

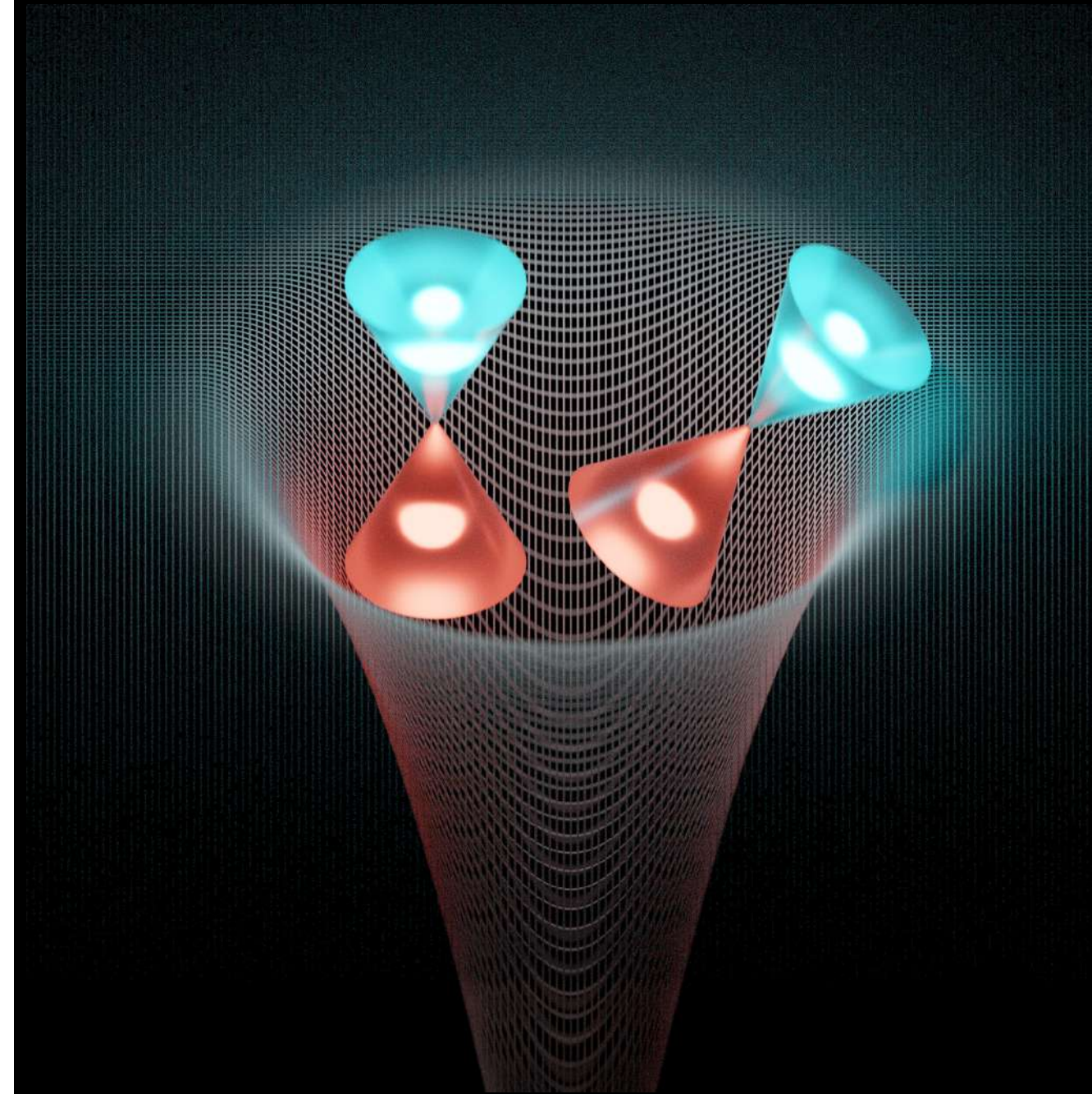
22.02.2022, UPC Barcelona

Quantum skyrmion lattices in Heisenberg ferromagnets

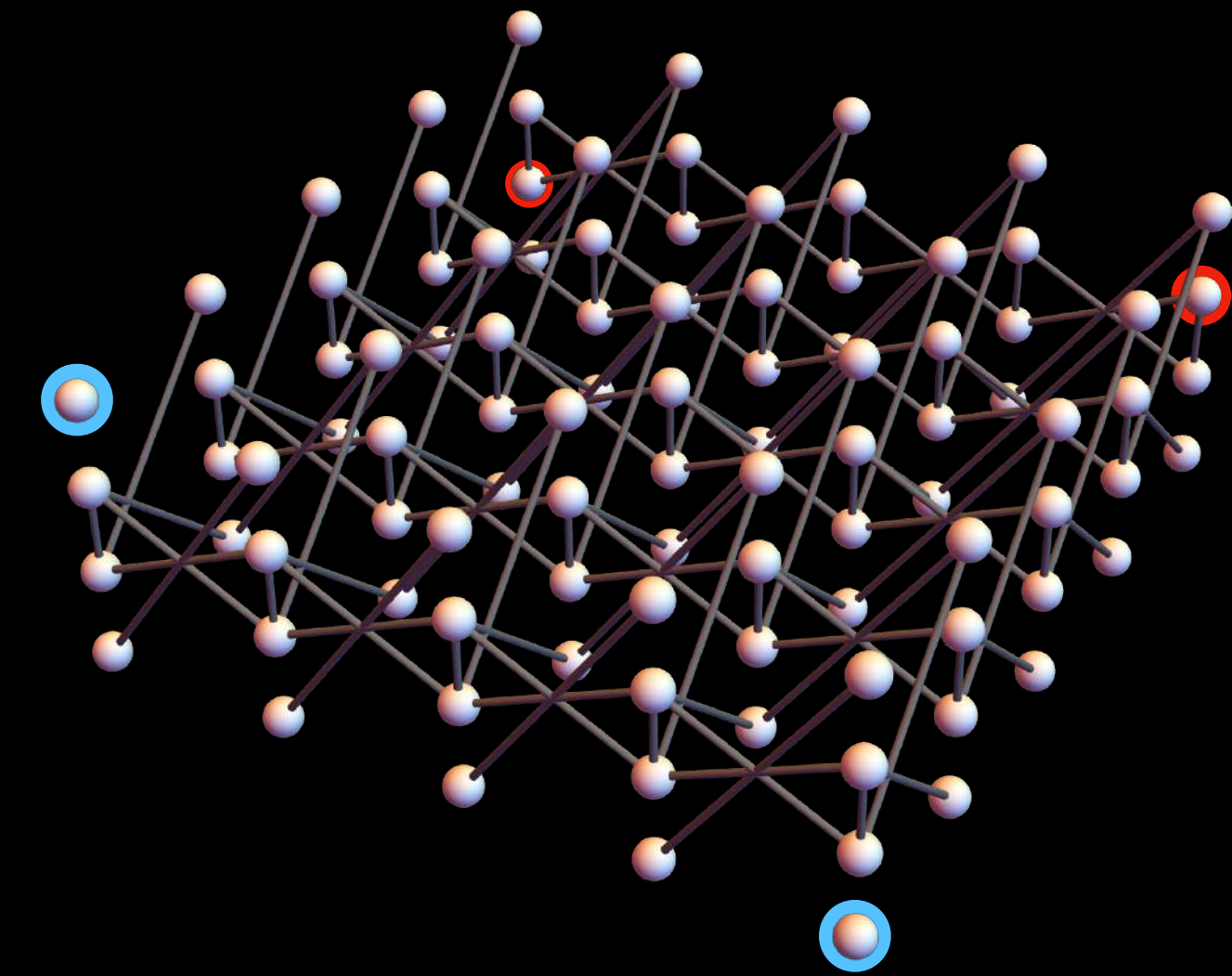
Andreas Haller, Department of Physics and Materials Science, University of Luxembourg

Outline

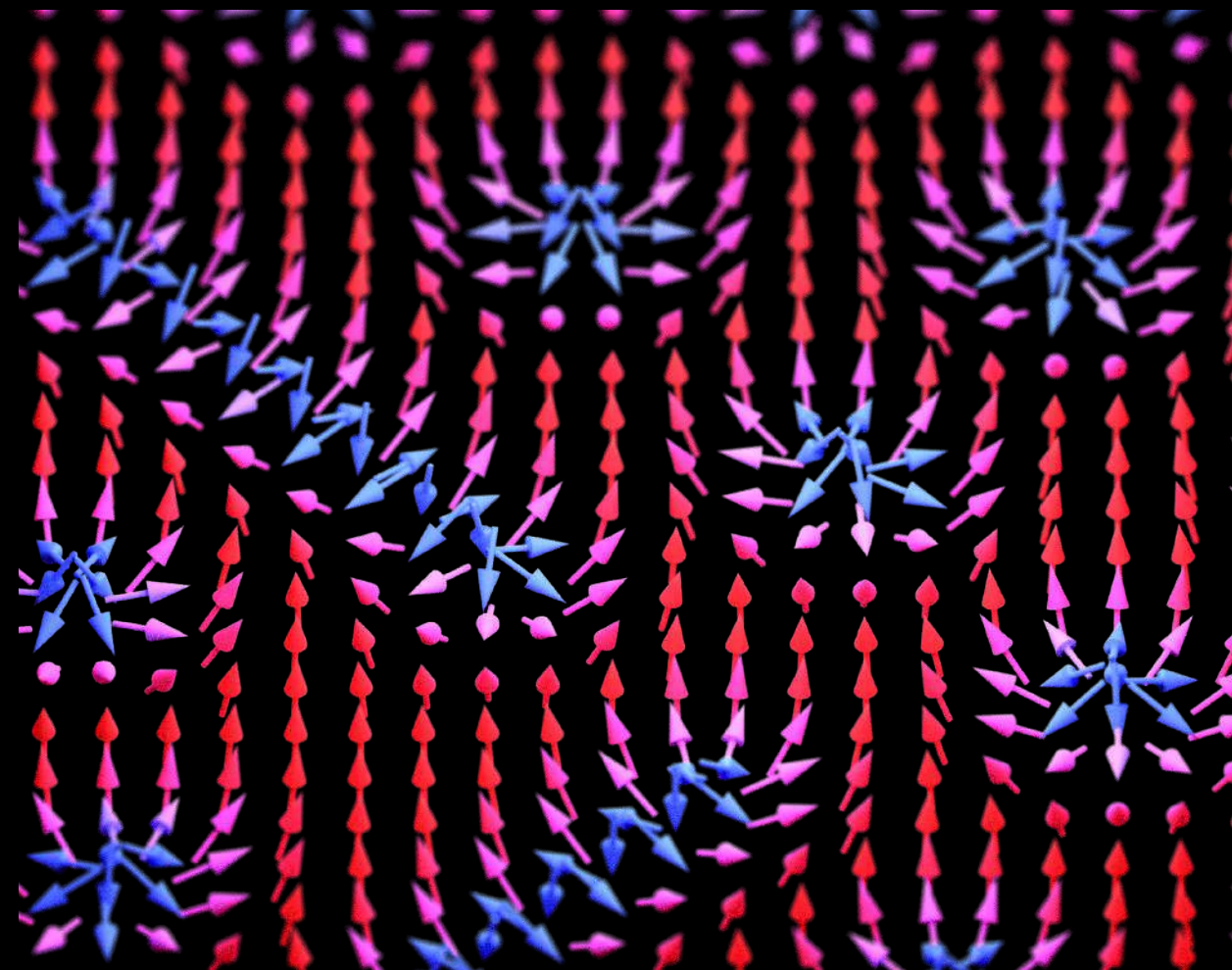
- Skyrmions
- Quantum skyrmion lattices
- Future directions
- Other research projects



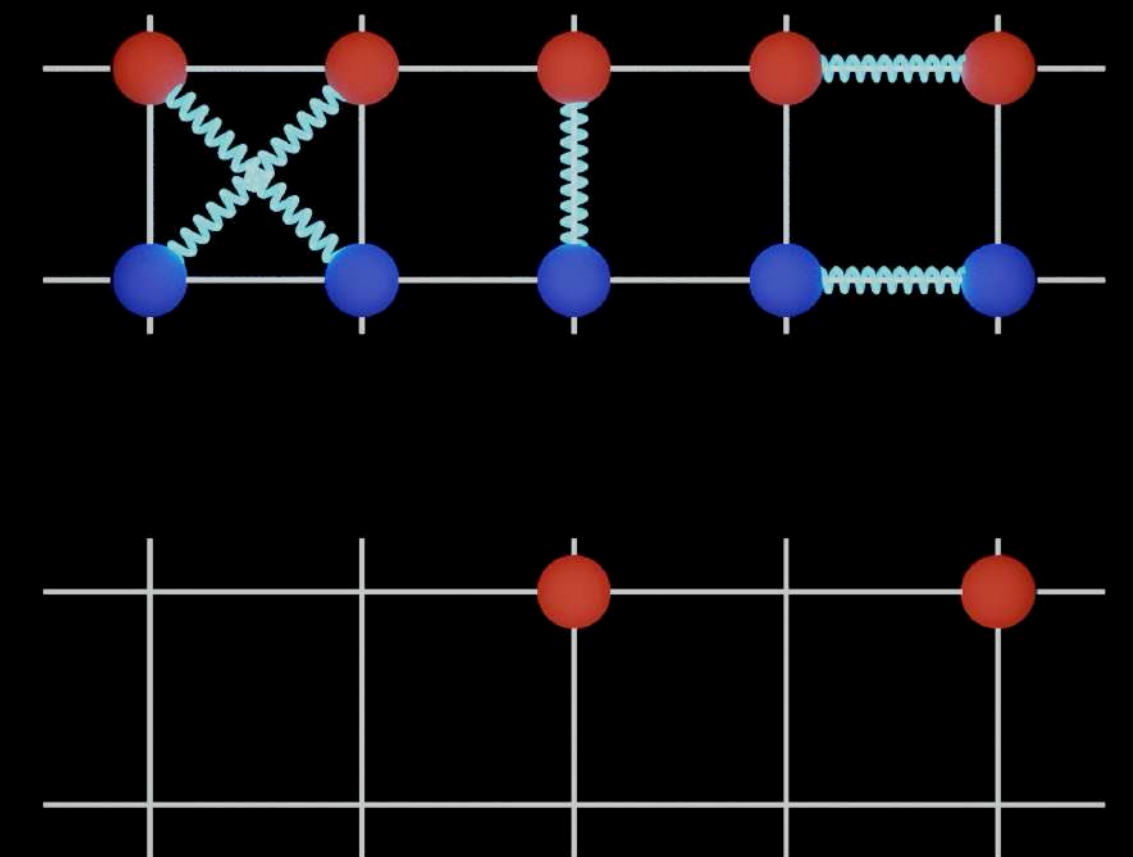
General relativity analogs
in tilted Weyl semimetals



Cookbook for
parafermion HOTIs



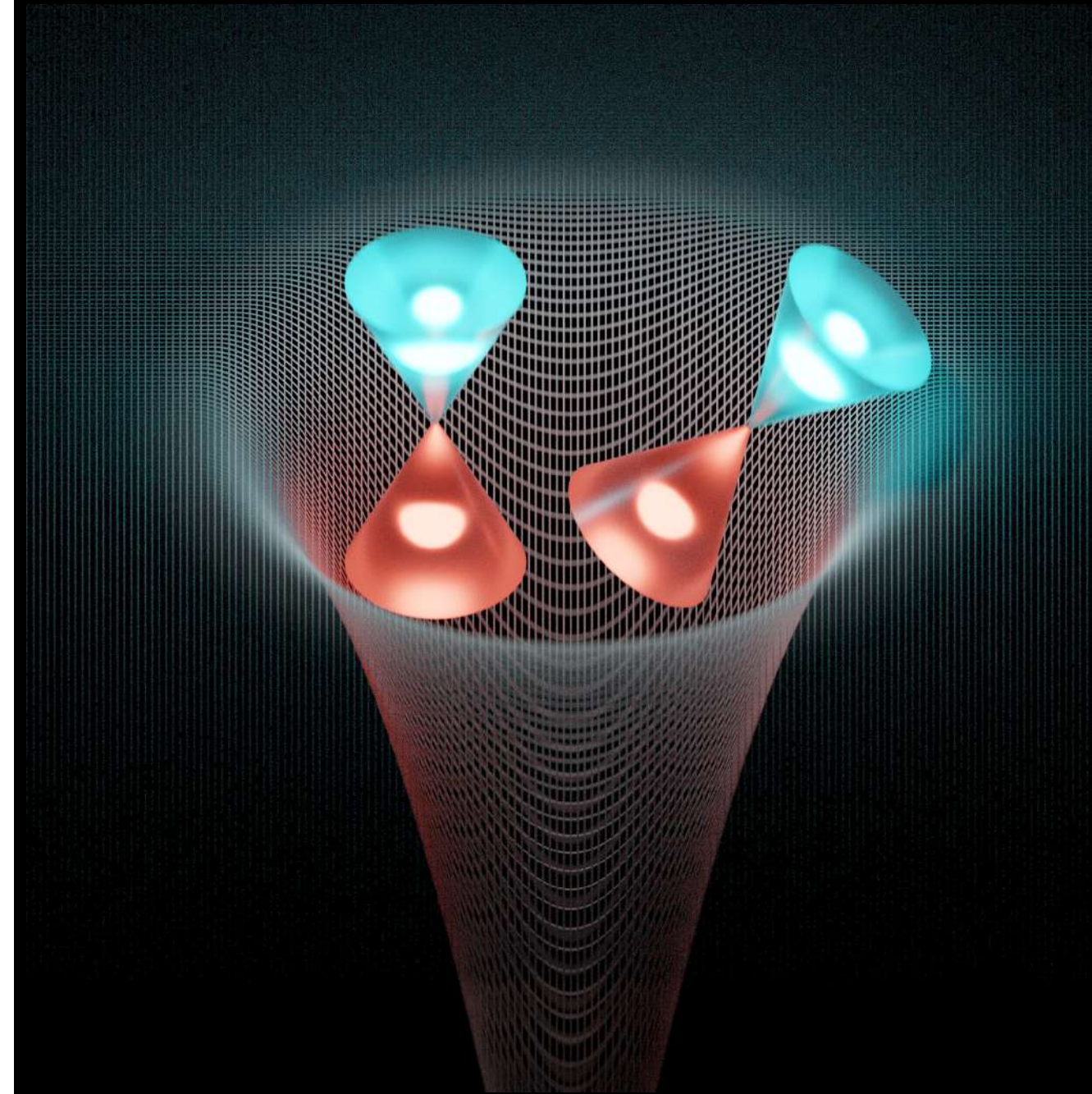
Quantum Spintronics



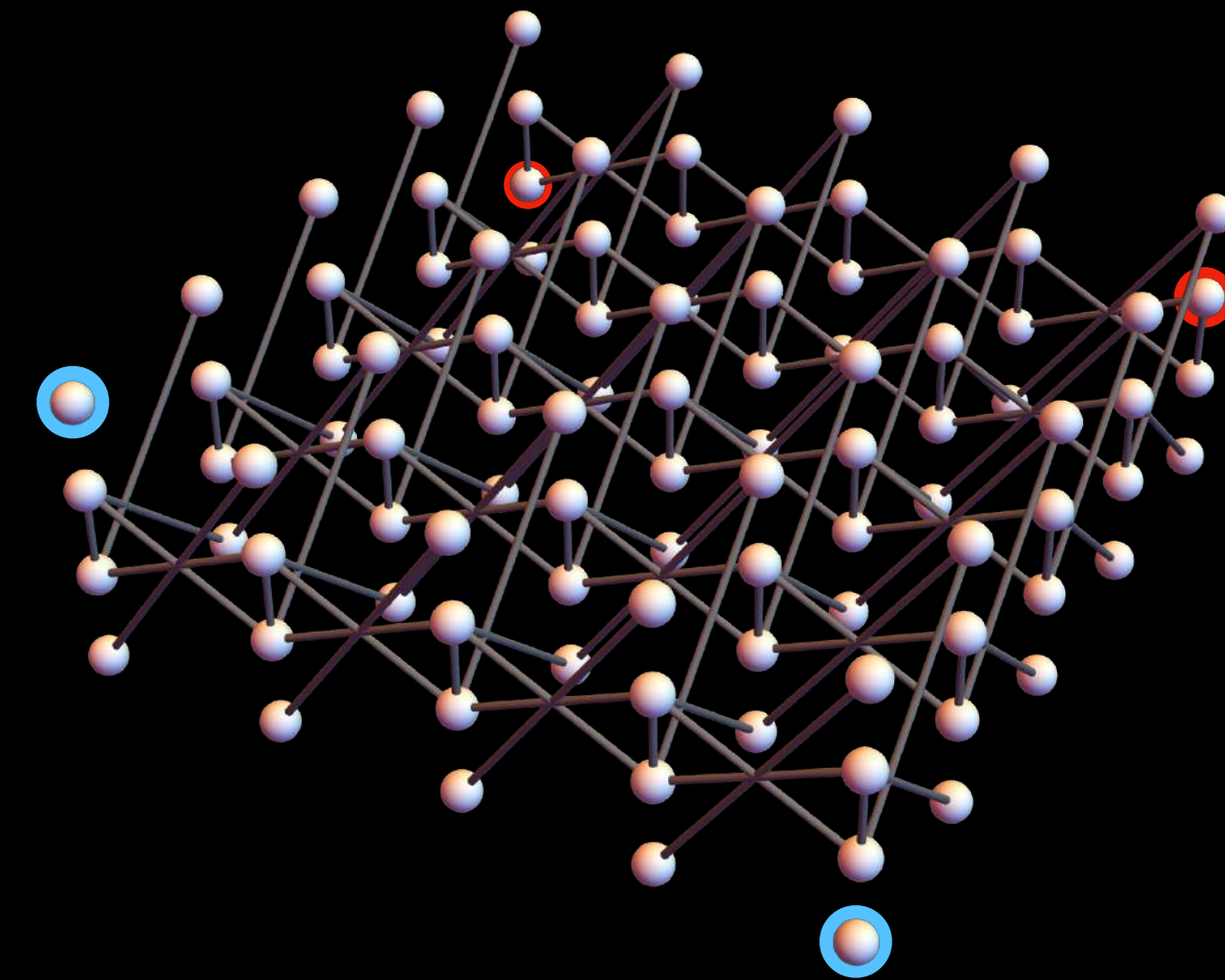
Interacting ladder systems

Outline

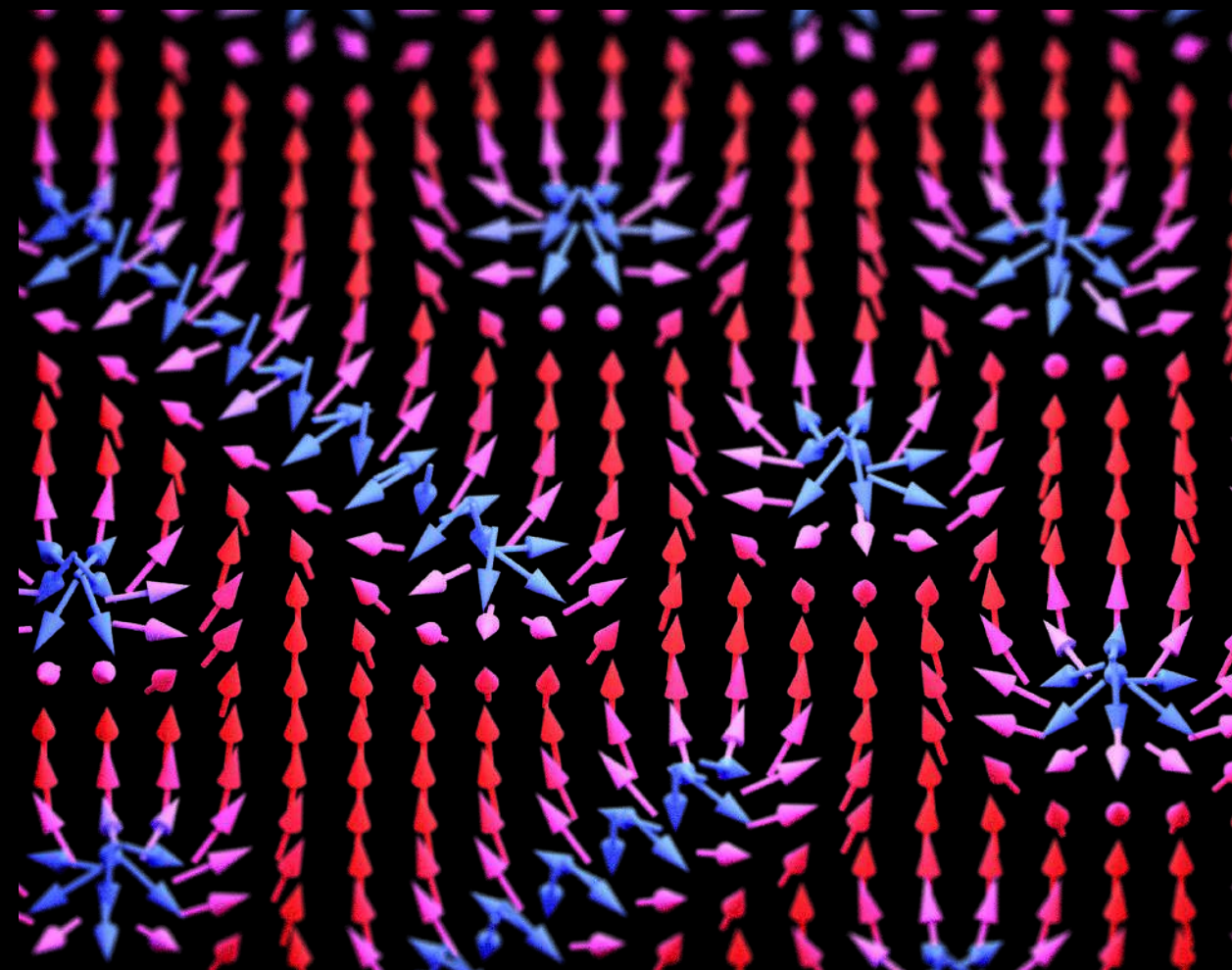
- Skyrmions
- Quantum skyrmion lattices
- Future directions
- Other research projects



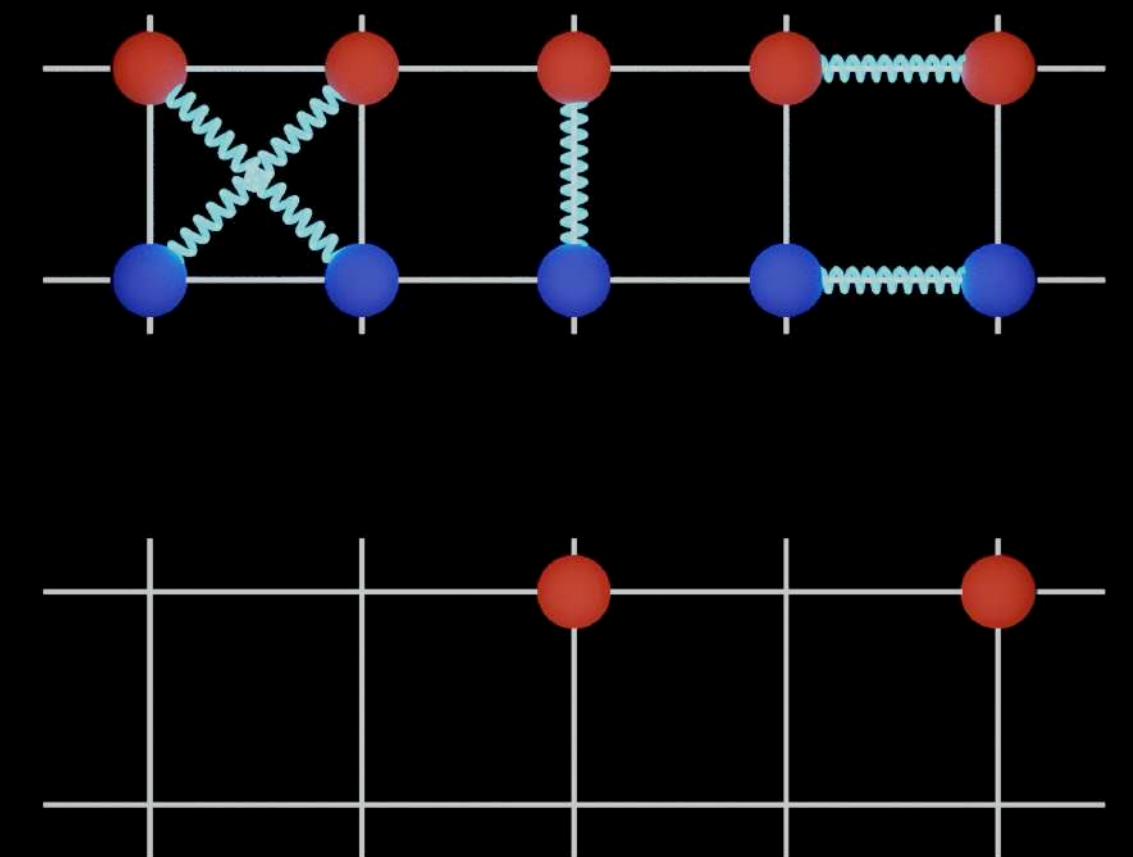
General relativity analogs
in tilted Weyl semimetals



Cookbook for
parafermion HOTIs



Quantum Spintronics

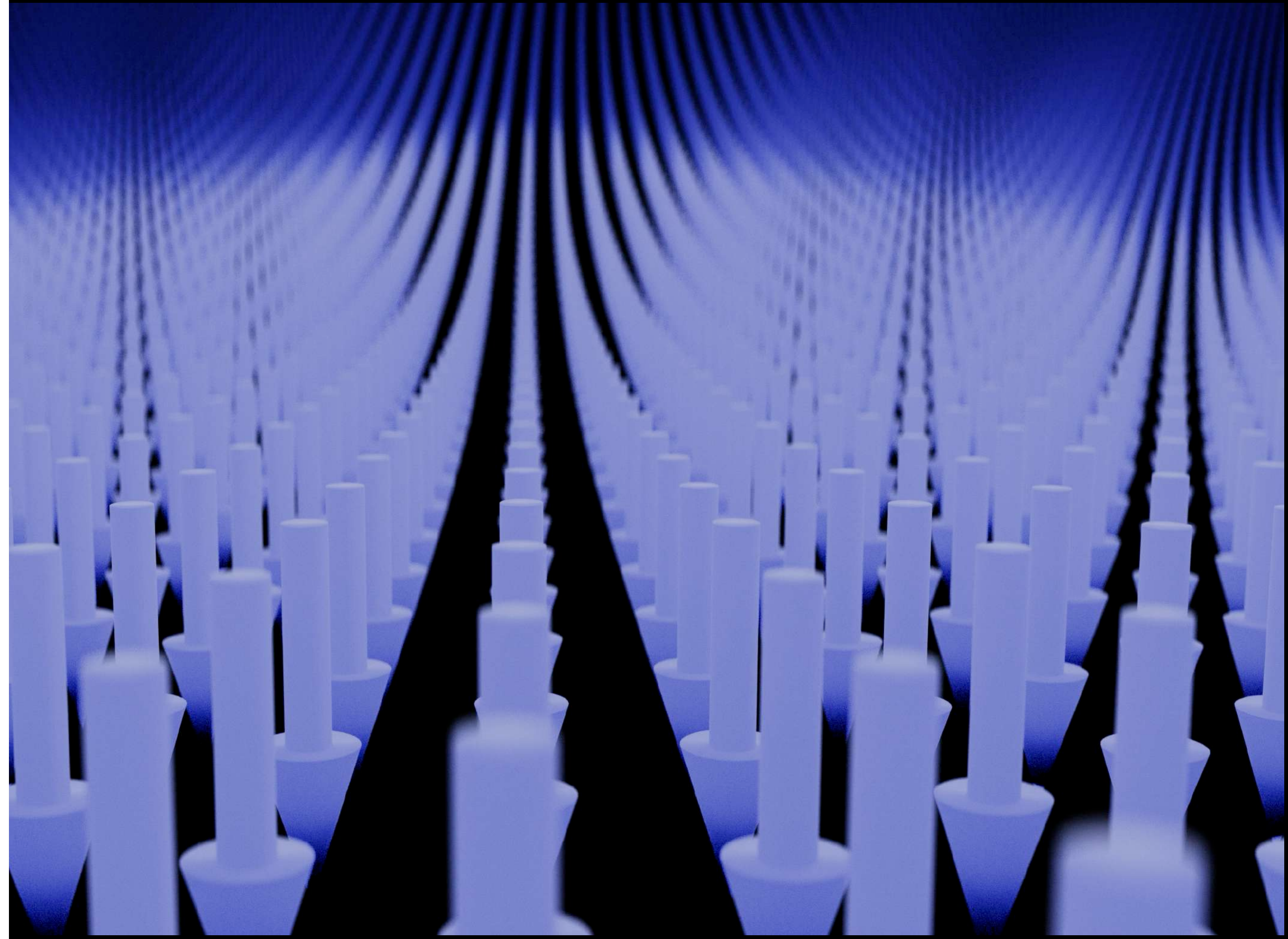


Interacting ladder systems

Skymions

Nature 442, 797 (2006), Science 323, 915 (2009)

Nature Reviews Materials 1, 1 (2016)

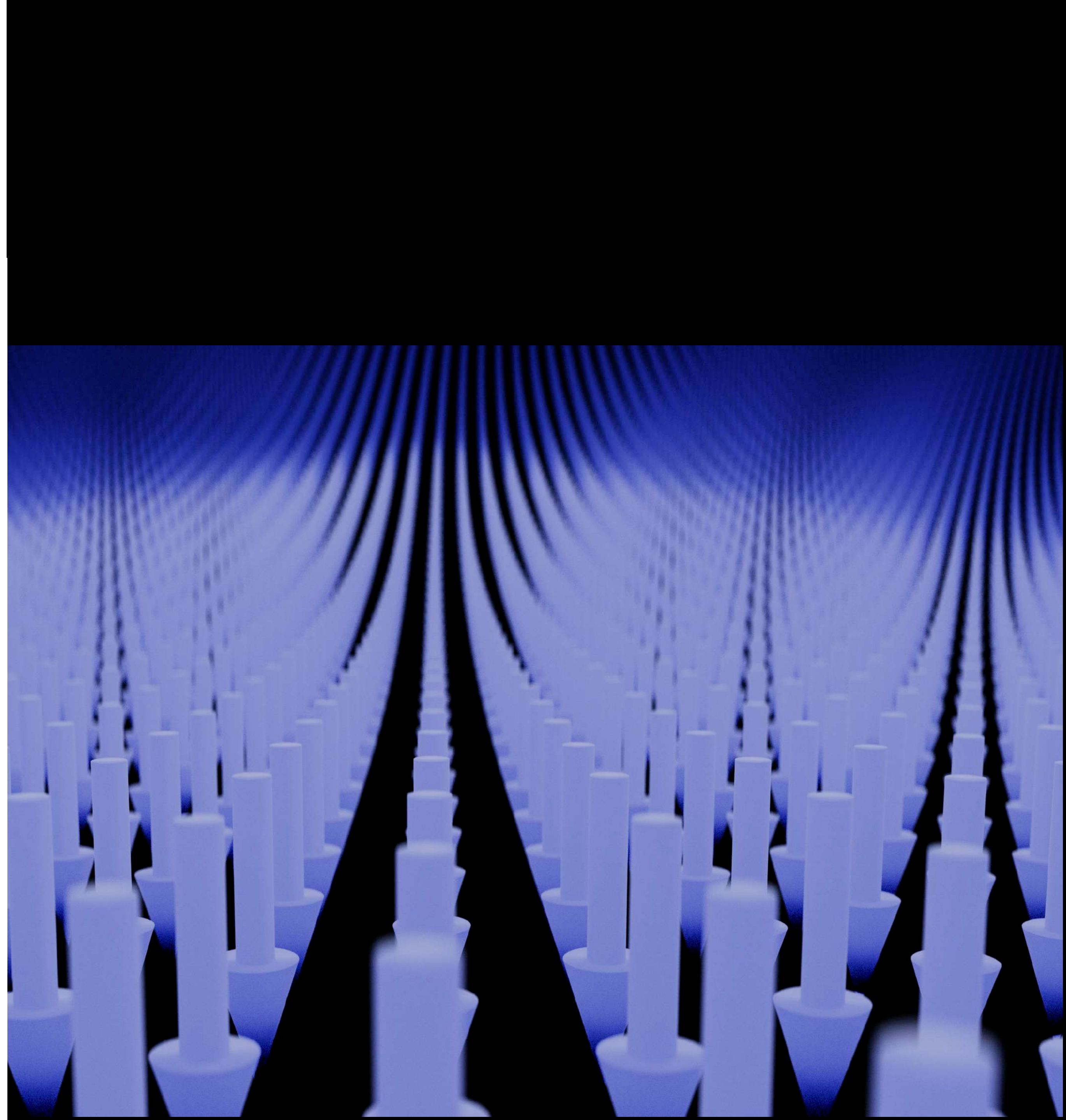


Skymions

Nature 442, 797 (2006), Science 323, 915 (2009)

Nature Reviews Materials 1, 1 (2016)

- quasiparticles in non-centrosymmetric materials

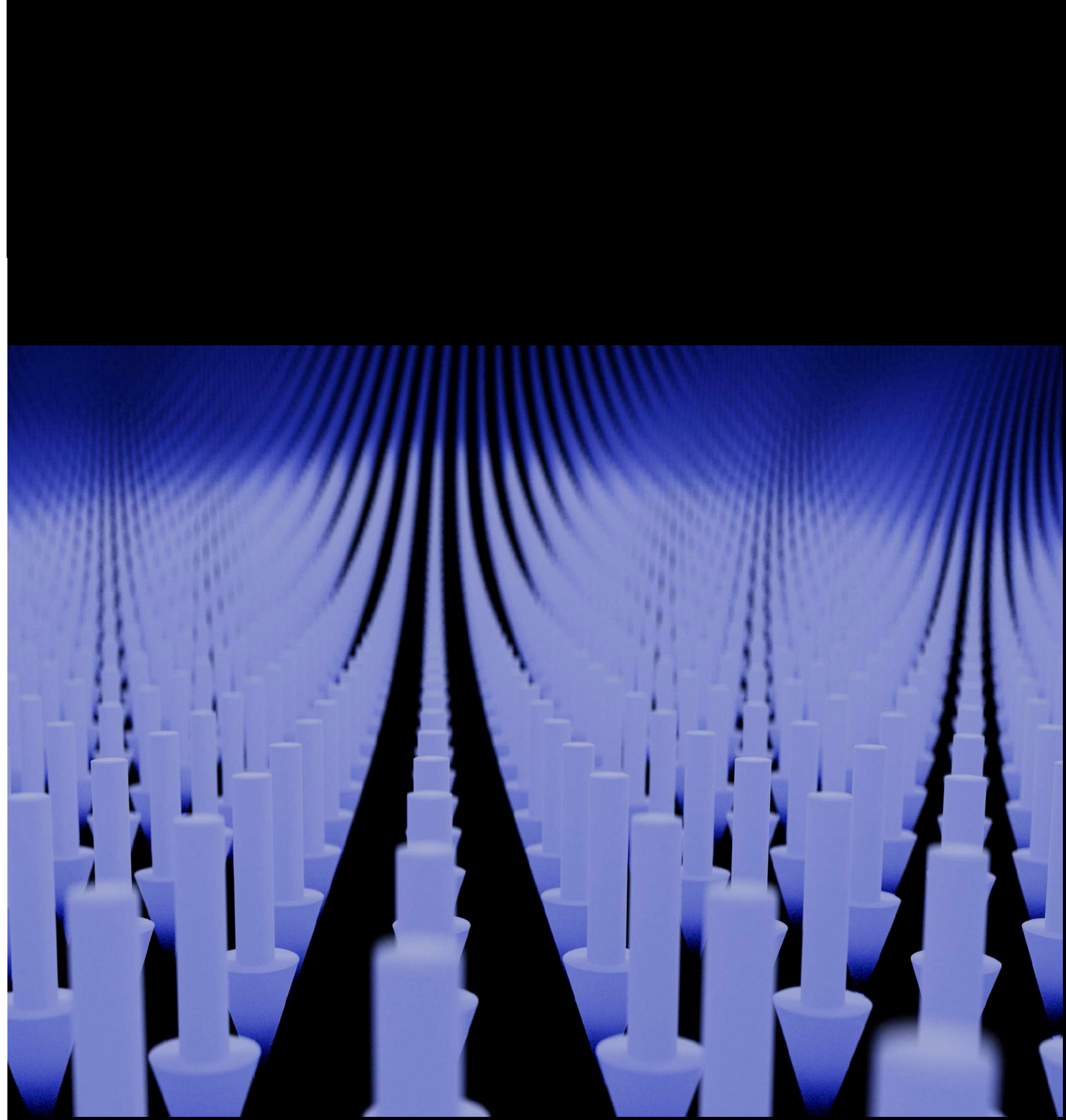


Skymions

Nature 442, 797 (2006), Science 323, 915 (2009)

Nature Reviews Materials 1, 1 (2016)

- quasiparticles in non-centrosymmetric materials
- Heisenberg model with Dzyaloshinskii-Moriya interaction

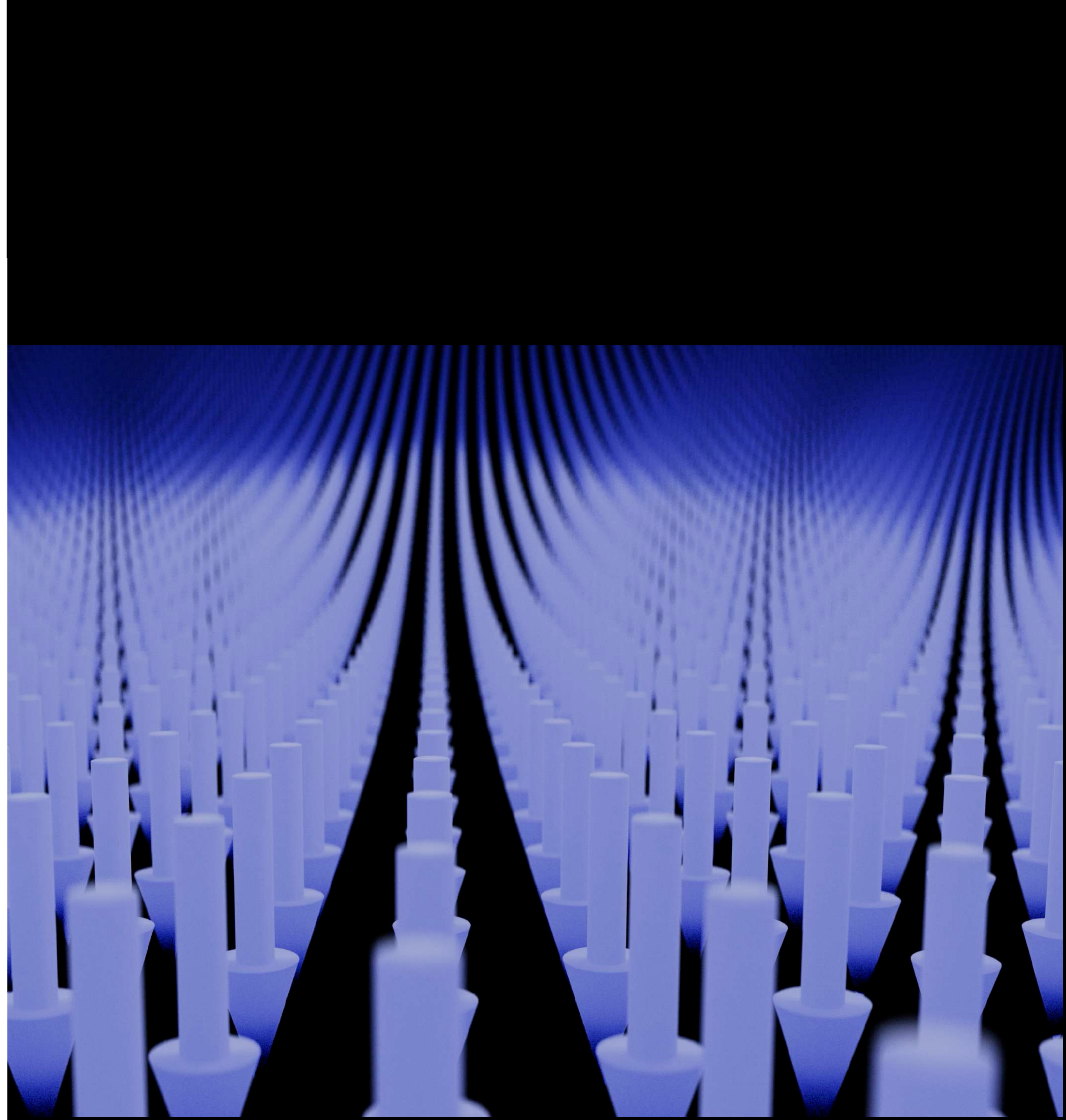


Skymions

Nature 442, 797 (2006), Science 323, 915 (2009)

Nature Reviews Materials 1, 1 (2016)

- quasiparticles in non-centrosymmetric materials
- Heisenberg model with Dzyaloshinskii-Moriya interaction
- real-space spin winding with quantized topological index

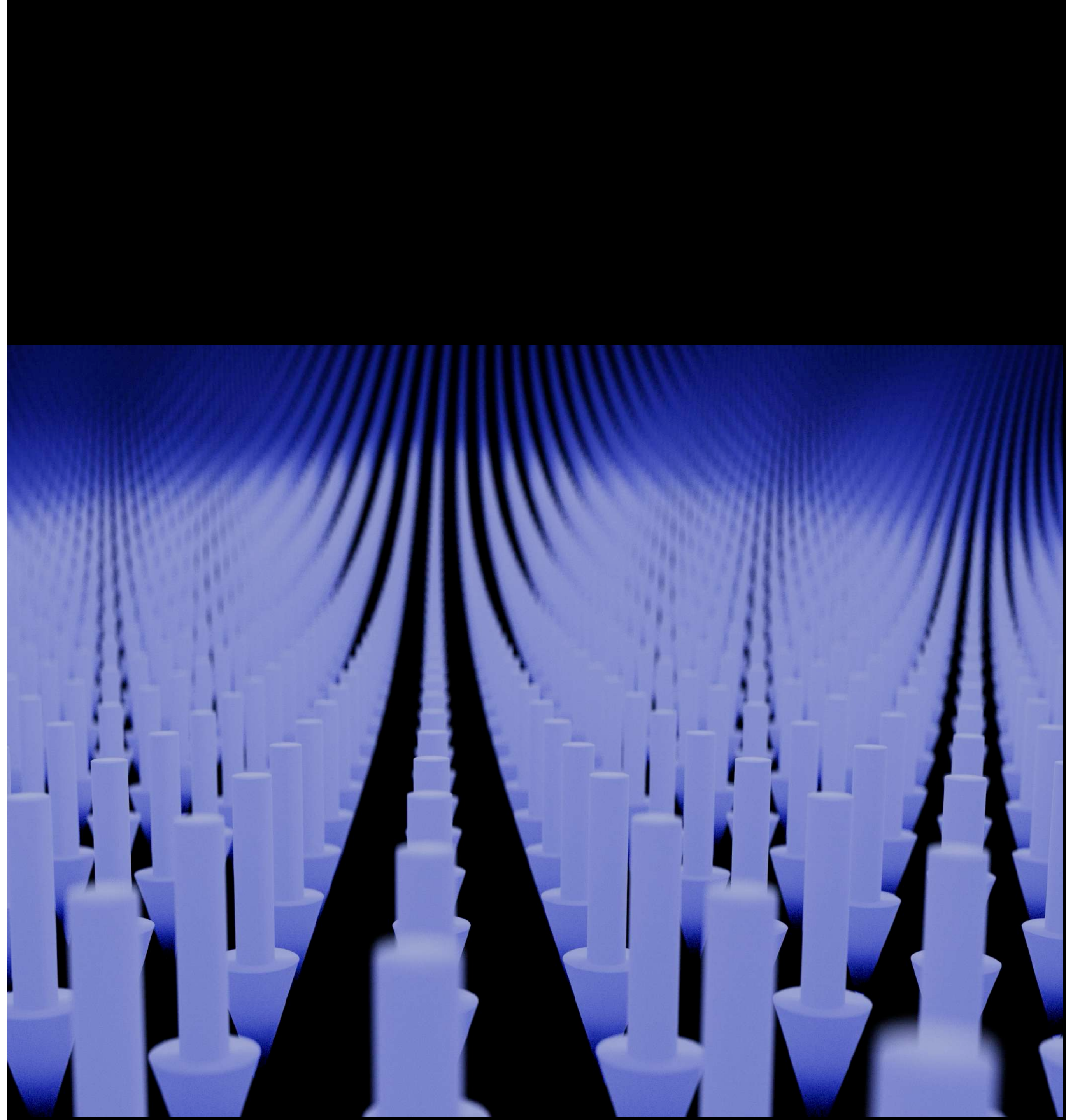


Skymions

Nature 442, 797 (2006), Science 323, 915 (2009)

Nature Reviews Materials 1, 1 (2016)

- quasiparticles in non-centrosymmetric materials
- Heisenberg model with Dzyaloshinskii-Moriya interaction
- real-space spin winding with quantized topological index
 - robust against disorder & temperature, long lifetimes, fast propagation, real-space manipulation by e.g. STM tips

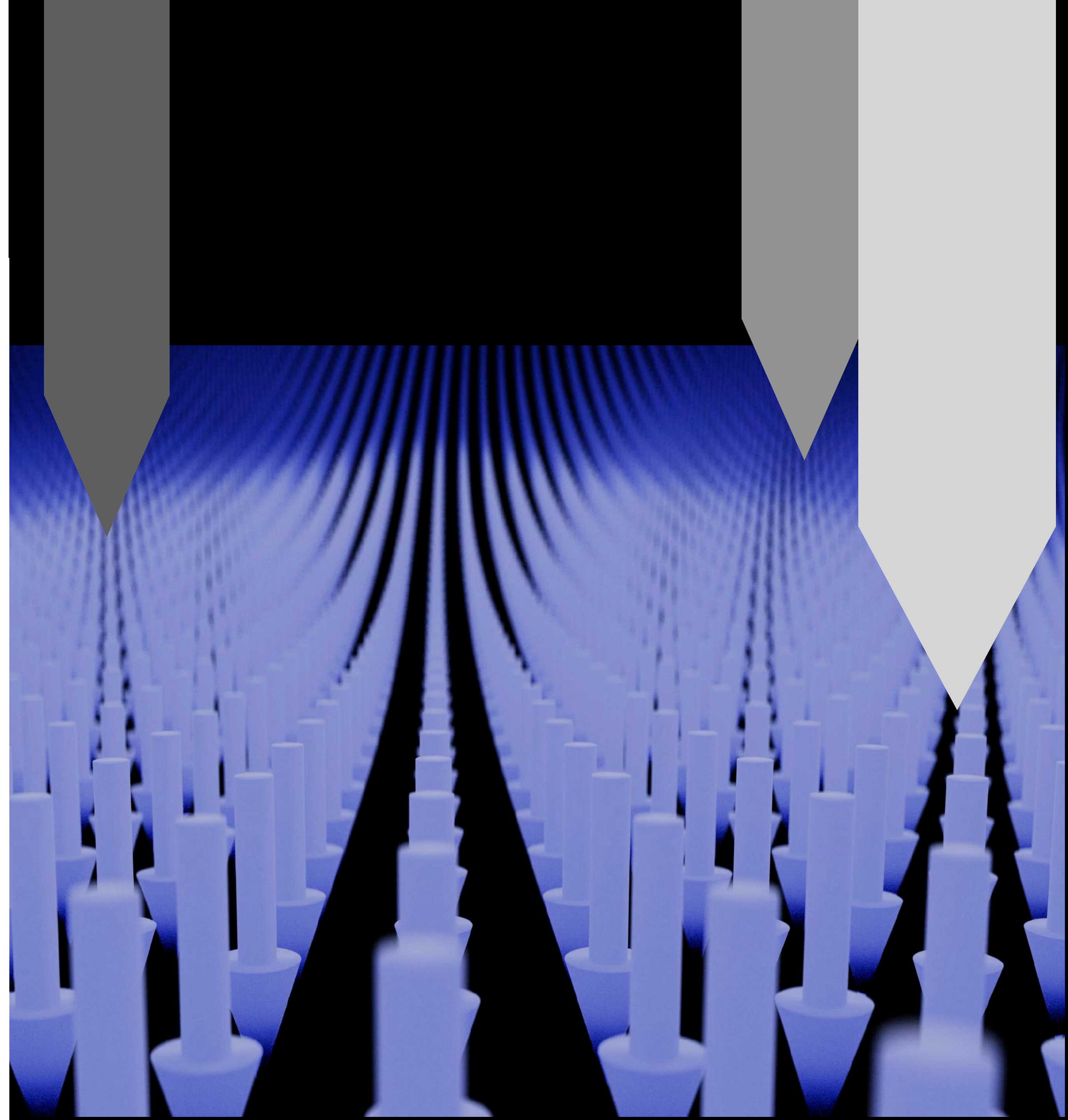


Skymions

Nature 442, 797 (2006), Science 323, 915 (2009)

Nature Reviews Materials 1, 1 (2016)

- quasiparticles in non-centrosymmetric materials
- Heisenberg model with Dzyaloshinskii-Moriya interaction
- real-space spin winding with quantized topological index
- robust against disorder & temperature, long lifetimes, fast propagation, real-space manipulation by e.g. STM tips



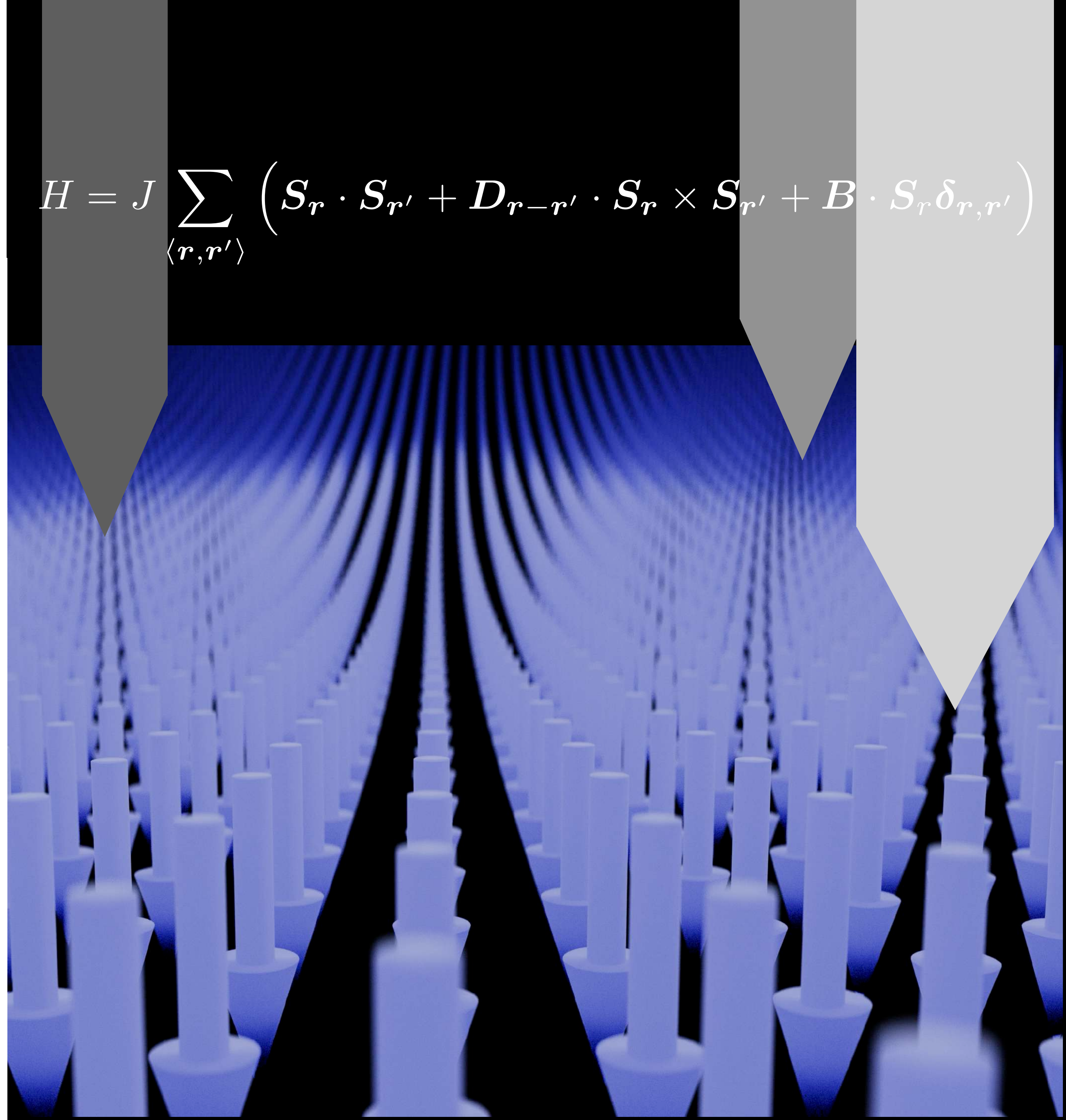
Skymions

Nature 442, 797 (2006), Science 323, 915 (2009)

Nature Reviews Materials 1, 1 (2016)

- quasiparticles in non-centrosymmetric materials
- Heisenberg model with Dzyaloshinskii-Moriya interaction
- real-space spin winding with quantized topological index
- robust against disorder & temperature, long lifetimes, fast propagation, real-space manipulation by e.g. STM tips

$$H = J \sum_{\langle r, r' \rangle} \left(\mathbf{S}_r \cdot \mathbf{S}_{r'} + \mathbf{D}_{r-r'} \cdot \mathbf{S}_r \times \mathbf{S}_{r'} + B \cdot \mathbf{S}_r \delta_{r,r'} \right)$$



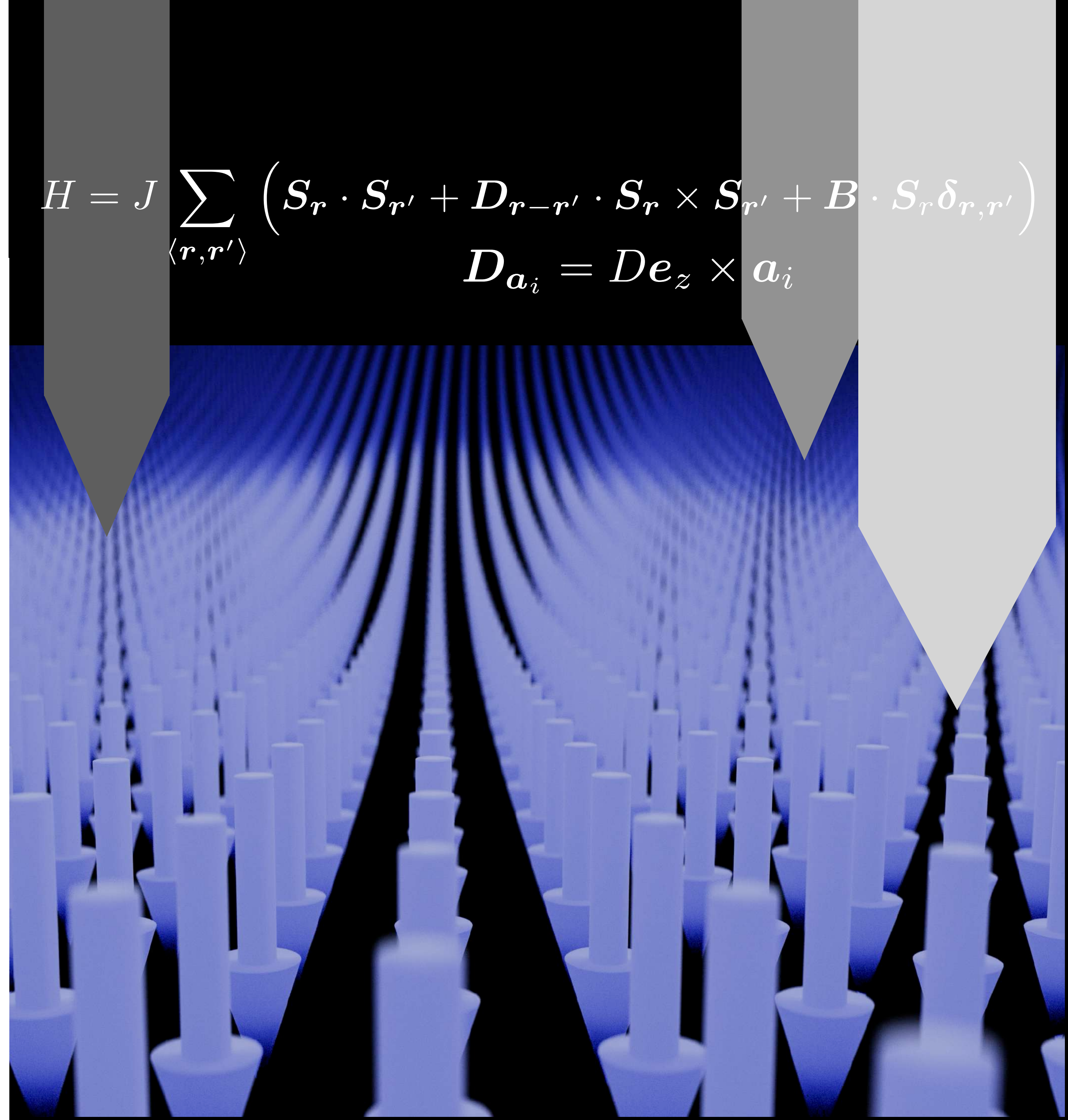
Skymions

Nature 442, 797 (2006), Science 323, 915 (2009)

Nature Reviews Materials 1, 1 (2016)

- quasiparticles in non-centrosymmetric materials
- Heisenberg model with Dzyaloshinskii-Moriya interaction
- real-space spin winding with quantized topological index
- robust against disorder & temperature, long lifetimes, fast propagation, real-space manipulation by e.g. STM tips

$$H = J \sum_{\langle r, r' \rangle} \left(\mathbf{S}_r \cdot \mathbf{S}_{r'} + \mathbf{D}_{r-r'} \cdot \mathbf{S}_r \times \mathbf{S}_{r'} + B \cdot \mathbf{S}_r \delta_{r,r'} \right)$$
$$\mathbf{D}_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$



Skyrmions

Nature 442, 797 (2006), Science 323, 915 (2009)

Nature Reviews Materials 1, 1 (2016)

- quasiparticles in non-centrosymmetric materials
- Heisenberg model with Dzyaloshinskii-Moriya interaction
- real-space spin winding with quantized topological index
- robust against disorder & temperature, long lifetimes, fast propagation, real-space manipulation by e.g. STM tips

$$H = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\mathbf{S}_{\mathbf{r}} \cdot \mathbf{S}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \mathbf{S}_{\mathbf{r}} \times \mathbf{S}_{\mathbf{r}'} + \mathbf{B} \cdot \mathbf{S}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$\mathbf{D}_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

Skyrmions

Nature 442, 797 (2006), Science 323, 915 (2009)

Nature Reviews Materials 1, 1 (2016)

- quasiparticles in non-centrosymmetric materials
- Heisenberg model with Dzyaloshinskii-Moriya interaction
- real-space spin winding with quantized topological index
- robust against disorder & temperature, long lifetimes, fast propagation, real-space manipulation by e.g. STM tips

$$J = -0.5D$$

$$H = J \sum_{\langle r, r' \rangle} \left(\mathbf{S}_r \cdot \mathbf{S}_{r'} + \mathbf{D}_{r-r'} \cdot \mathbf{S}_r \times \mathbf{S}_{r'} + \mathbf{B} \cdot \mathbf{S}_r \delta_{r, r'} \right)$$

$$\mathbf{D}_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

Skyrmions

Nature 442, 797 (2006), Science 323, 915 (2009)

Nature Reviews Materials 1, 1 (2016)

- quasiparticles in non-centrosymmetric materials
- Heisenberg model with Dzyaloshinskii-Moriya interaction
- real-space spin winding with quantized topological index
- robust against disorder & temperature, long lifetimes, fast propagation, real-space manipulation by e.g. STM tips

$$J = -0.5D \qquad B = -B_0 \mathbf{e}_z$$
$$H = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\mathbf{S}_{\mathbf{r}} \cdot \mathbf{S}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \mathbf{S}_{\mathbf{r}} \times \mathbf{S}_{\mathbf{r}'} + B \cdot \mathbf{S}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$\mathbf{D}_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

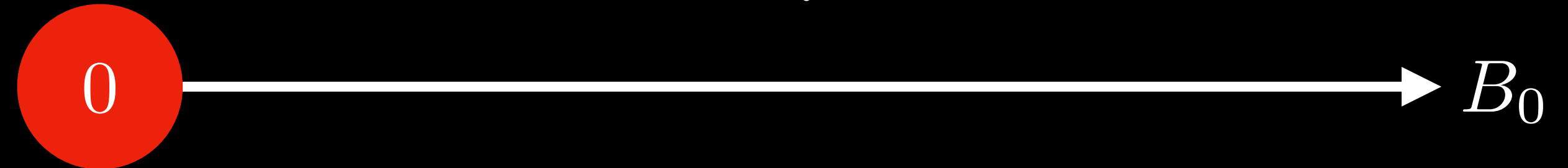
Skyrmions

Nature 442, 797 (2006), Science 323, 915 (2009)

Nature Reviews Materials 1, 1 (2016)

- quasiparticles in non-centrosymmetric materials
- Heisenberg model with Dzyaloshinskii-Moriya interaction
- real-space spin winding with quantized topological index
 - robust against disorder & temperature, long lifetimes, fast propagation, real-space manipulation by e.g. STM tips

$$J = -0.5D \quad B = -B_0 e_z$$
$$H = J \sum_{\langle r, r' \rangle} \left(S_r \cdot S_{r'} + D_{r-r'} \cdot S_r \times S_{r'} + B \cdot S_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$



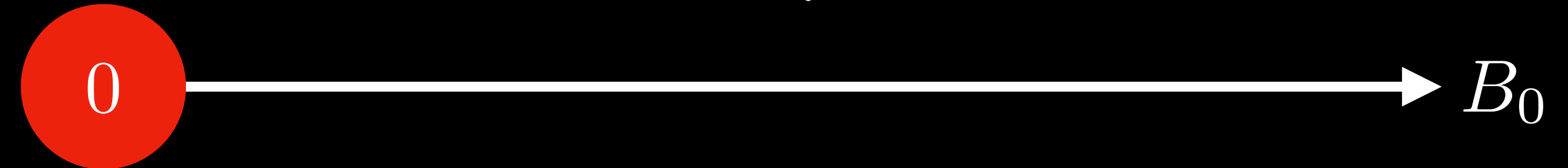
Skyrmions

Nature 442, 797 (2006), Science 323, 915 (2009)

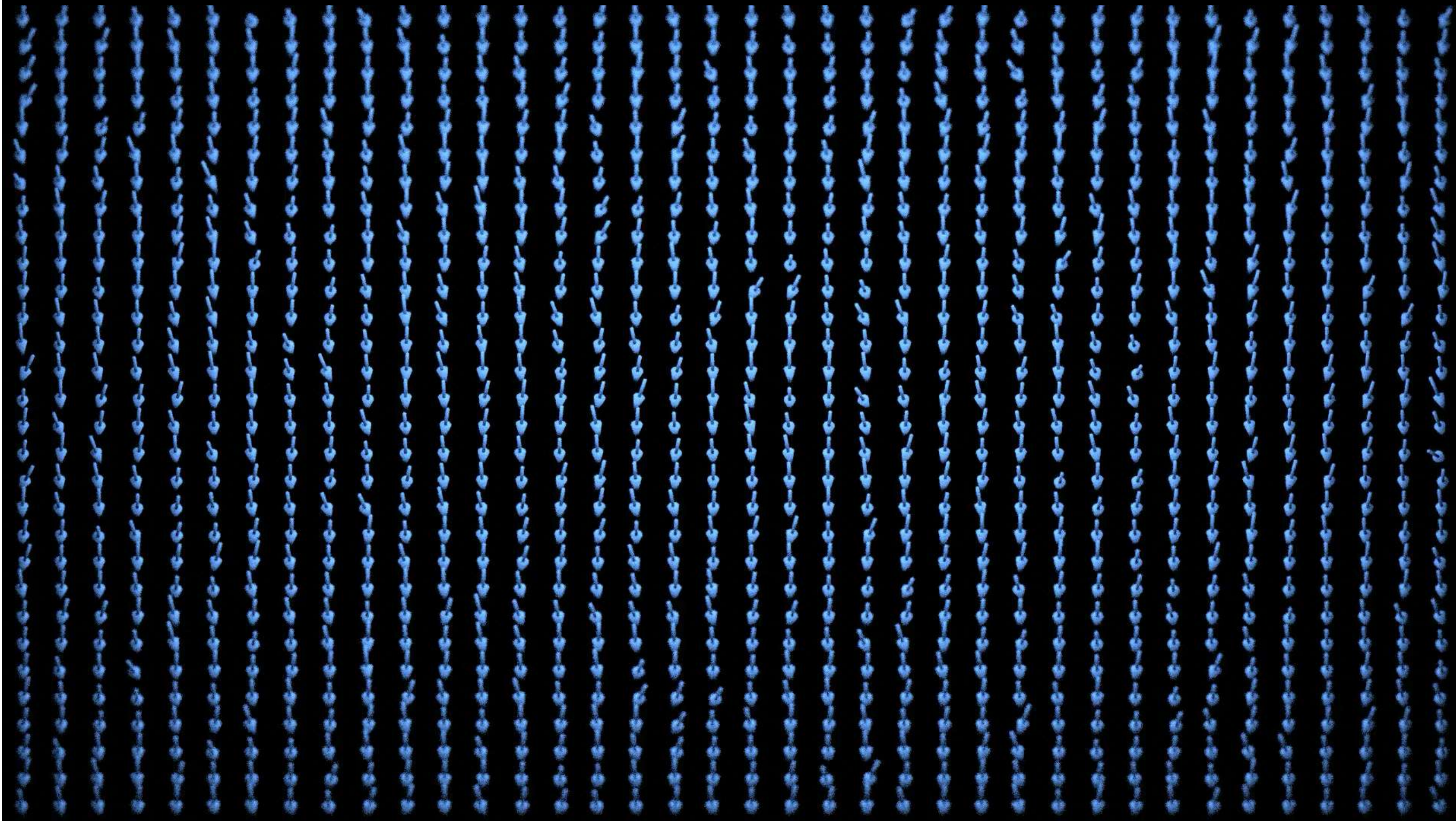
Nature Reviews Materials 1, 1 (2016)

- quasiparticles in non-centrosymmetric materials
- Heisenberg model with Dzyaloshinskii-Moriya interaction
- real-space spin winding with quantized topological index
 - robust against disorder & temperature, long lifetimes, fast propagation, real-space manipulation by e.g. STM tips

$$J = -0.5D \quad B = -B_0 e_z$$
$$H = J \sum_{\langle r, r' \rangle} \left(S_r \cdot S_{r'} + D_{r-r'} \cdot S_r \times S_{r'} + B \cdot S_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$



MC sampling with only local rotations



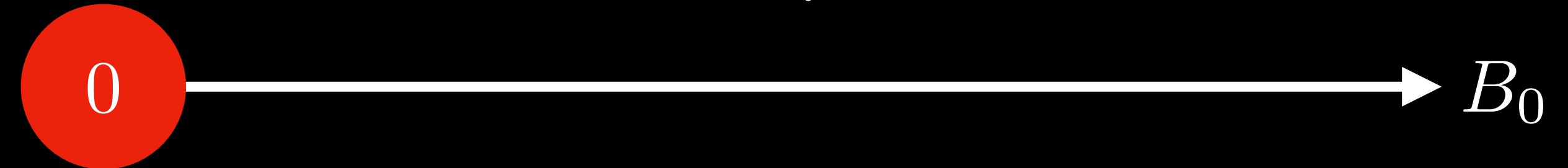
Skyrmions

Nature 442, 797 (2006), Science 323, 915 (2009)

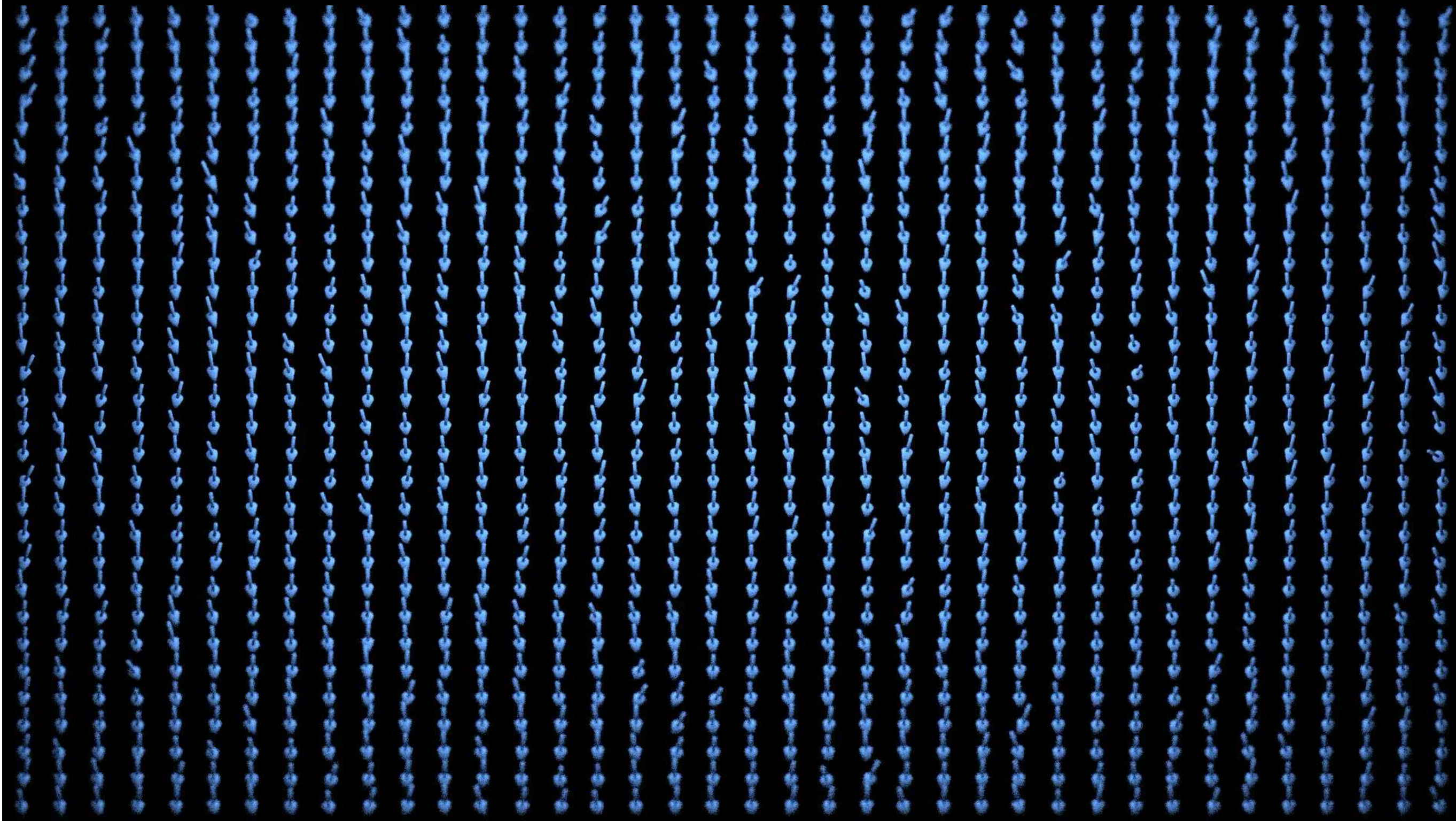
Nature Reviews Materials 1, 1 (2016)

- quasiparticles in non-centrosymmetric materials
- Heisenberg model with Dzyaloshinskii-Moriya interaction
- real-space spin winding with quantized topological index
 - robust against disorder & temperature, long lifetimes, fast propagation, real-space manipulation by e.g. STM tips

$$J = -0.5D \quad B = -B_0 e_z$$
$$H = J \sum_{\langle r, r' \rangle} \left(S_r \cdot S_{r'} + D_{r-r'} \cdot S_r \times S_{r'} + B \cdot S_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$



MC sampling with only local rotations

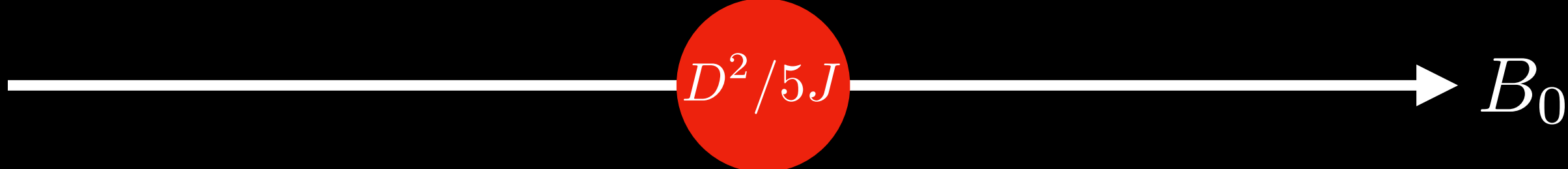


Skymions

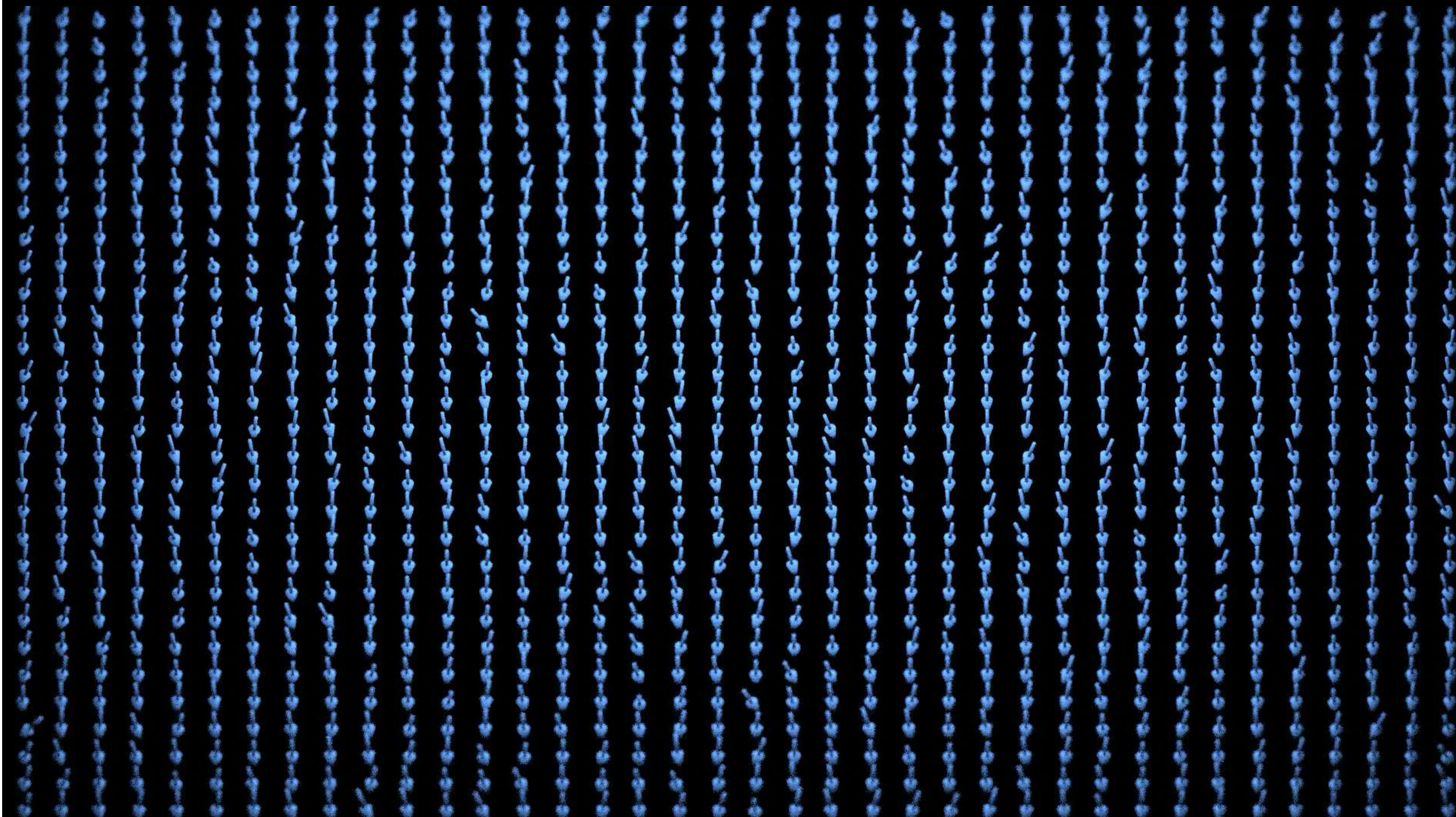
Nature 442, 797 (2006), Science 323, 915 (2009)

Nature Reviews Materials 1, 1 (2016)

- quasiparticles in non-centrosymmetric materials
- Heisenberg model with Dzyaloshinskii-Moriya interaction
- real-space spin winding with quantized topological index
 - robust against disorder & temperature, long lifetimes, fast propagation, real-space manipulation by e.g. STM tips

$$J = -0.5D \quad B = -B_0 \mathbf{e}_z$$
$$H = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\mathbf{S}_{\mathbf{r}} \cdot \mathbf{S}_{\mathbf{r}'} + D_{\mathbf{r}-\mathbf{r}'} \cdot \mathbf{S}_{\mathbf{r}} \times \mathbf{S}_{\mathbf{r}'} + B \cdot \mathbf{S}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$D_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$


MC sampling with only local rotations

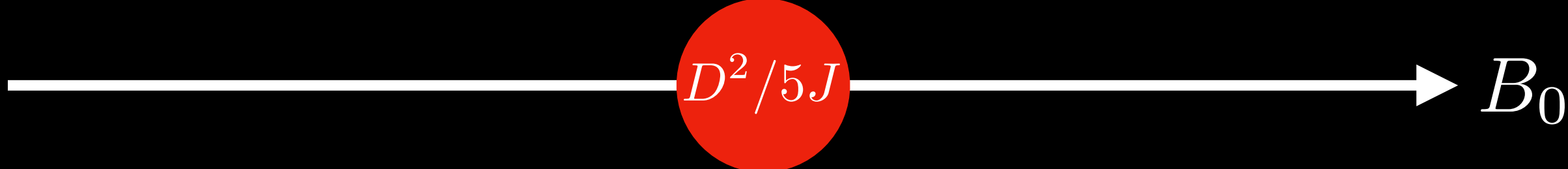


Skymions

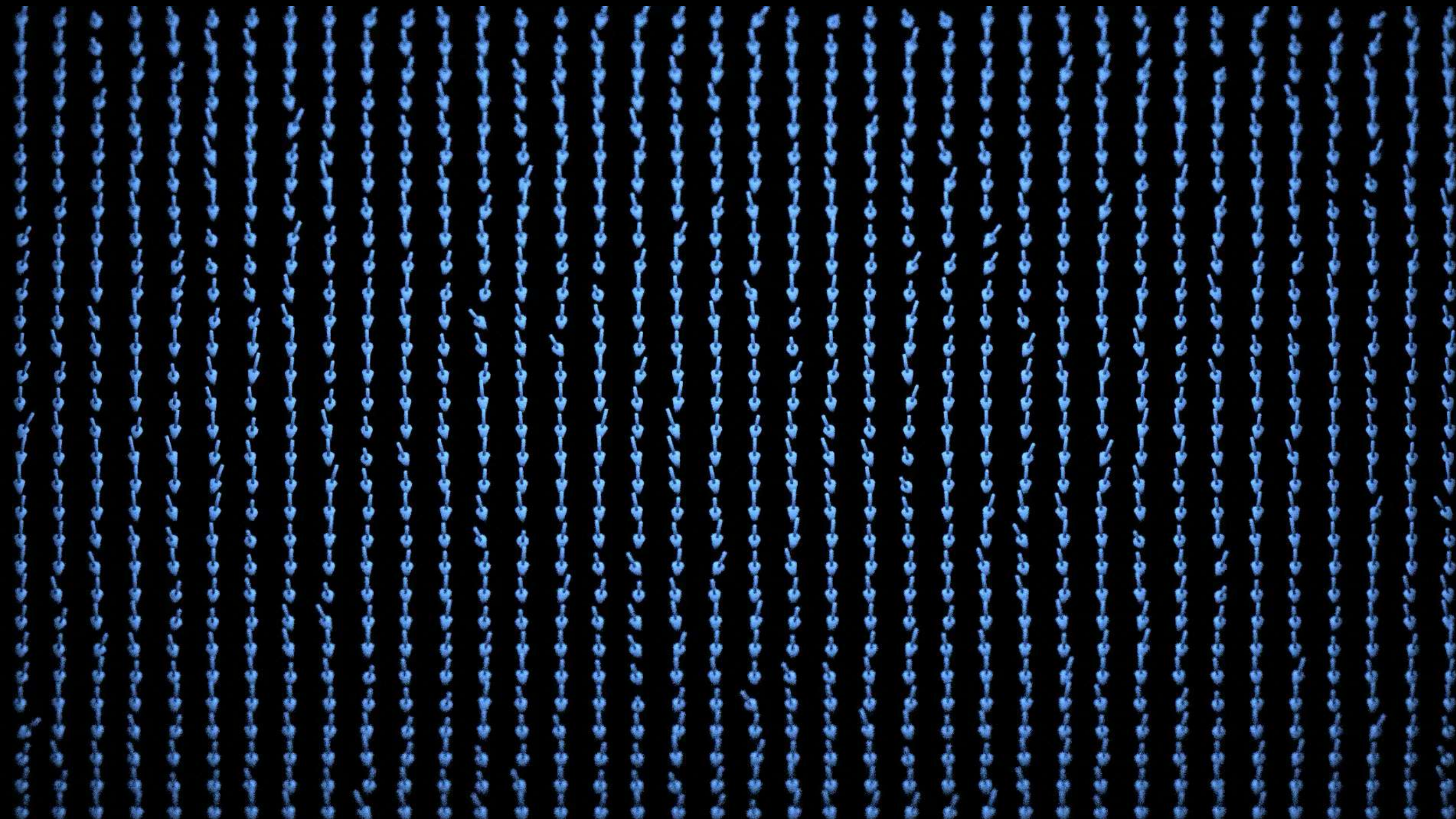
Nature 442, 797 (2006), Science 323, 915 (2009)

Nature Reviews Materials 1, 1 (2016)

- quasiparticles in non-centrosymmetric materials
- Heisenberg model with Dzyaloshinskii-Moriya interaction
- real-space spin winding with quantized topological index
 - robust against disorder & temperature, long lifetimes, fast propagation, real-space manipulation by e.g. STM tips

$$J = -0.5D \quad B = -B_0 e_z$$
$$H = J \sum_{\langle r, r' \rangle} \left(S_r \cdot S_{r'} + D_{r-r'} \cdot S_r \times S_{r'} + B \cdot S_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$


MC sampling with only local rotations



Skymions

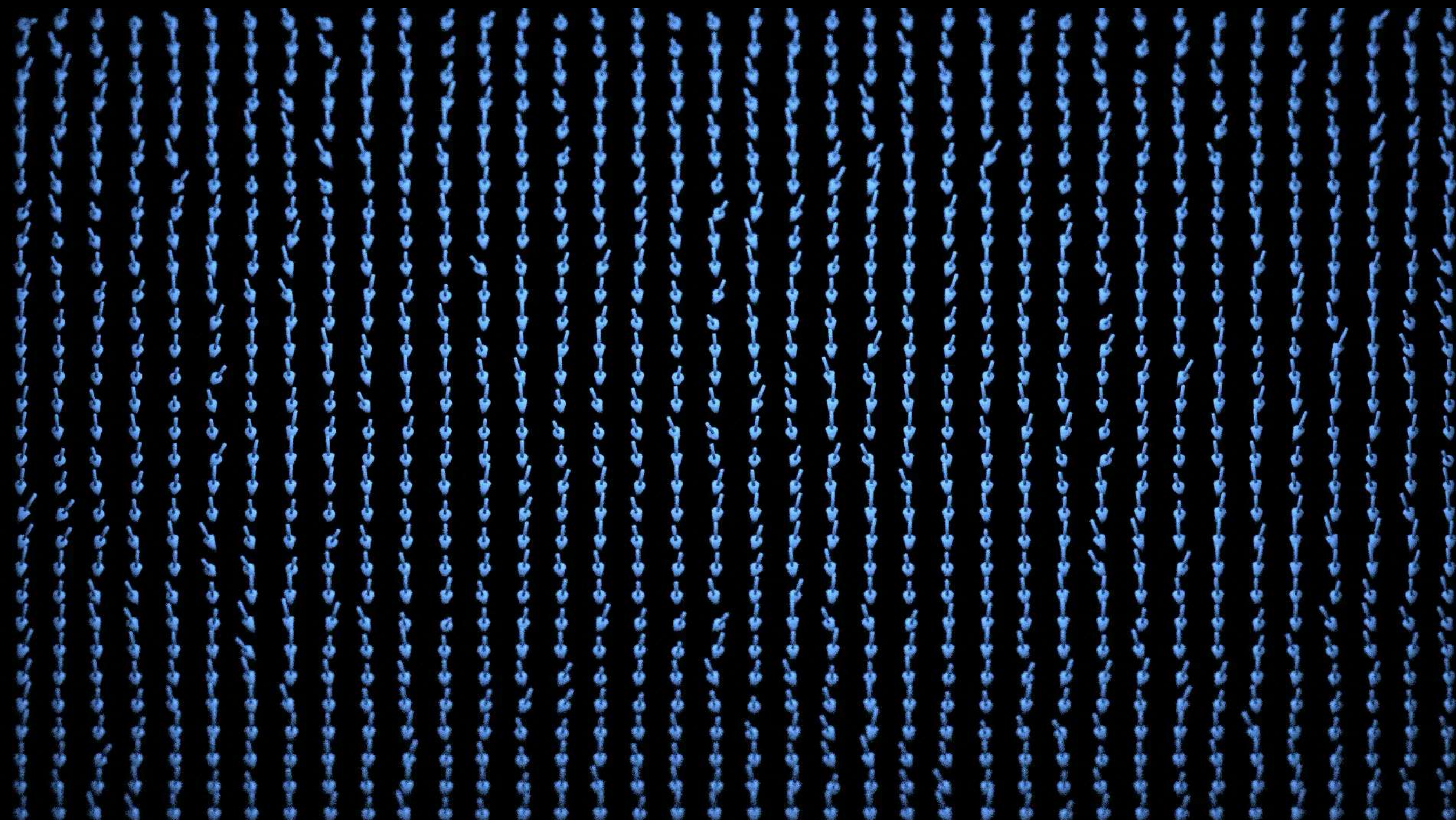
Nature 442, 797 (2006), Science 323, 915 (2009)

Nature Reviews Materials 1, 1 (2016)

- quasiparticles in non-centrosymmetric materials
- Heisenberg model with Dzyaloshinskii-Moriya interaction
- real-space spin winding with quantized topological index
 - robust against disorder & temperature, long lifetimes, fast propagation, real-space manipulation by e.g. STM tips

$$J = -0.5D \quad B = -B_0 e_z$$
$$H = J \sum_{\langle r, r' \rangle} \left(S_r \cdot S_{r'} + D_{r-r'} \cdot S_r \times S_{r'} + B \cdot S_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$


MC sampling with only local rotations



Skyrmions

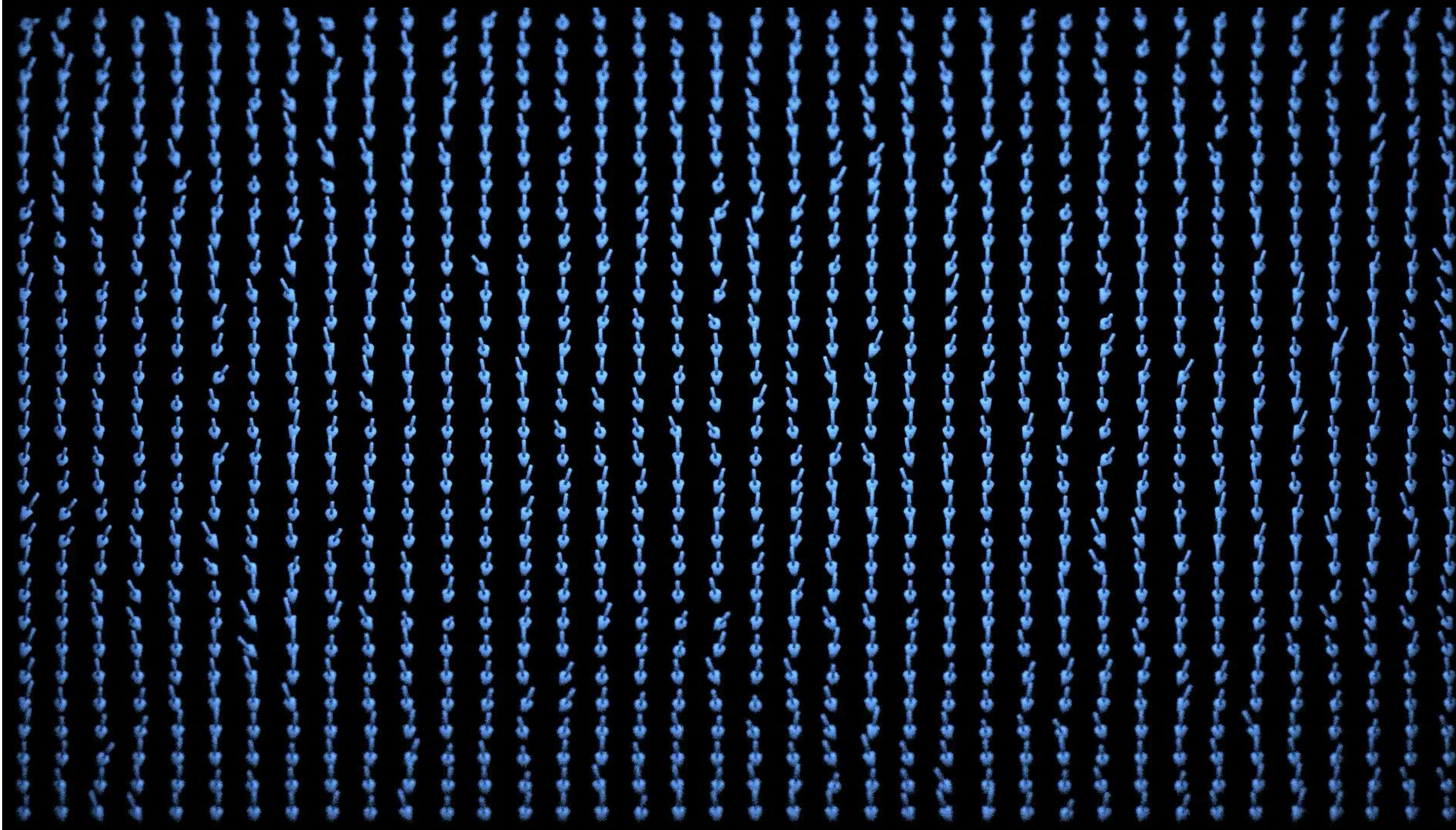
Nature 442, 797 (2006), Science 323, 915 (2009)

Nature Reviews Materials 1, 1 (2016)

- quasiparticles in non-centrosymmetric materials
- Heisenberg model with Dzyaloshinskii-Moriya interaction
- real-space spin winding with quantized topological index
 - robust against disorder & temperature, long lifetimes, fast propagation, real-space manipulation by e.g. STM tips

$$J = -0.5D \quad B = -B_0 \mathbf{e}_z$$
$$H = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\mathbf{S}_{\mathbf{r}} \cdot \mathbf{S}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \mathbf{S}_{\mathbf{r}} \times \mathbf{S}_{\mathbf{r}'} + B \cdot \mathbf{S}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$\mathbf{D}_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$


MC sampling with only local rotations



Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

$$\begin{aligned} J &= -0.5D & B &= -B_0 e_z \\ H &= J \sum_{\langle r, r' \rangle} \left(S_r \cdot S_{r'} + D_{r-r'} \cdot S_r \times S_{r'} + B \cdot S_r \delta_{r, r'} \right) \\ D_{a_i} &= D e_z \times a_i \end{aligned}$$

Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations

$$\begin{aligned} J &= -0.5D & B &= -B_0 e_z \\ H &= J \sum_{\langle r, r' \rangle} \left(S_r \cdot S_{r'} + D_{r-r'} \cdot S_r \times S_{r'} + B \cdot S_r \delta_{r, r'} \right) \\ D_{a_i} &= D e_z \times a_i \end{aligned}$$

Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations

$$J = -0.5D \quad B = -B_0 e_z$$
$$H = J \sum_{\langle r, r' \rangle} \left(S_r \cdot S_{r'} + D_{r-r'} \cdot S_r \times S_{r'} + B \cdot S_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

- Find the spin configuration Ω that minimizes H

Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations

$$J = -0.5D \quad B = -B_0 e_z$$
$$H = J \sum_{\langle r, r' \rangle} \left(S_r \cdot S_{r'} + D_{r-r'} \cdot S_r \times S_{r'} + B \cdot S_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

- Find the spin configuration Ω that minimizes H
- Quantize the Hamiltonian $H \rightarrow \hat{H}$

Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

- Find the spin configuration Ω that minimizes H
- Quantize the Hamiltonian $H \rightarrow \hat{H}$

Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

- Find the spin configuration Ω that minimizes H
- Quantize the Hamiltonian $H \rightarrow \hat{H}$
- Do a Holstein-Primakoff transformation
 $\hat{S}_i \Omega_i = S - \hat{n}_i \quad S^+ = \sqrt{2S - \hat{n}_i} \hat{a}_i$

Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

- Find the spin configuration Ω that minimizes H
- Quantize the Hamiltonian $H \rightarrow \hat{H}$
- Do a Holstein-Primakoff transformation
- Approximate Hamiltonian depends on Ω

$$\hat{S}_i \Omega_i = S - \hat{n}_i \quad S^+ = \sqrt{2S - \hat{n}_i} \hat{a}_i$$
$$\hat{H}_{SW} = E_0 + \sum_{i,j} \begin{pmatrix} \hat{a}_i^\dagger & \hat{a}_j \end{pmatrix} \begin{pmatrix} t_{ij} & \tau_{ij} \\ \tau_{ij}^* & t_{ij}^* \end{pmatrix} \begin{pmatrix} \hat{a}_i \\ \hat{a}_j^\dagger \end{pmatrix}$$

Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

- Find the spin configuration Ω that minimizes H
- Quantize the Hamiltonian $H \rightarrow \hat{H}$
- Do a Holstein-Primakoff transformation
 $\hat{S}_i \Omega_i = S - \hat{n}_i \quad S^+ = \sqrt{2S - \hat{n}_i} \hat{a}_i$

- Approximate Hamiltonian depends on Ω
$$\hat{H}_{SW} = E_0 + \sum_{i,j} \begin{pmatrix} \hat{a}_i^\dagger & \hat{a}_j \end{pmatrix} \begin{pmatrix} t_{ij} & \tau_{ij} \\ \tau_{ij}^* & t_{ij}^* \end{pmatrix} \begin{pmatrix} \hat{a}_i \\ \hat{a}_j^\dagger \end{pmatrix}$$

- Bogoliubov transformation for bosons
$$\hat{H}_{SW} = E_0 + \sum_{\nu=1}^N \hbar \omega_\nu \left(\alpha_\nu^\dagger \alpha_\nu + \frac{1}{2} \right)$$

Quantum Skyrmions

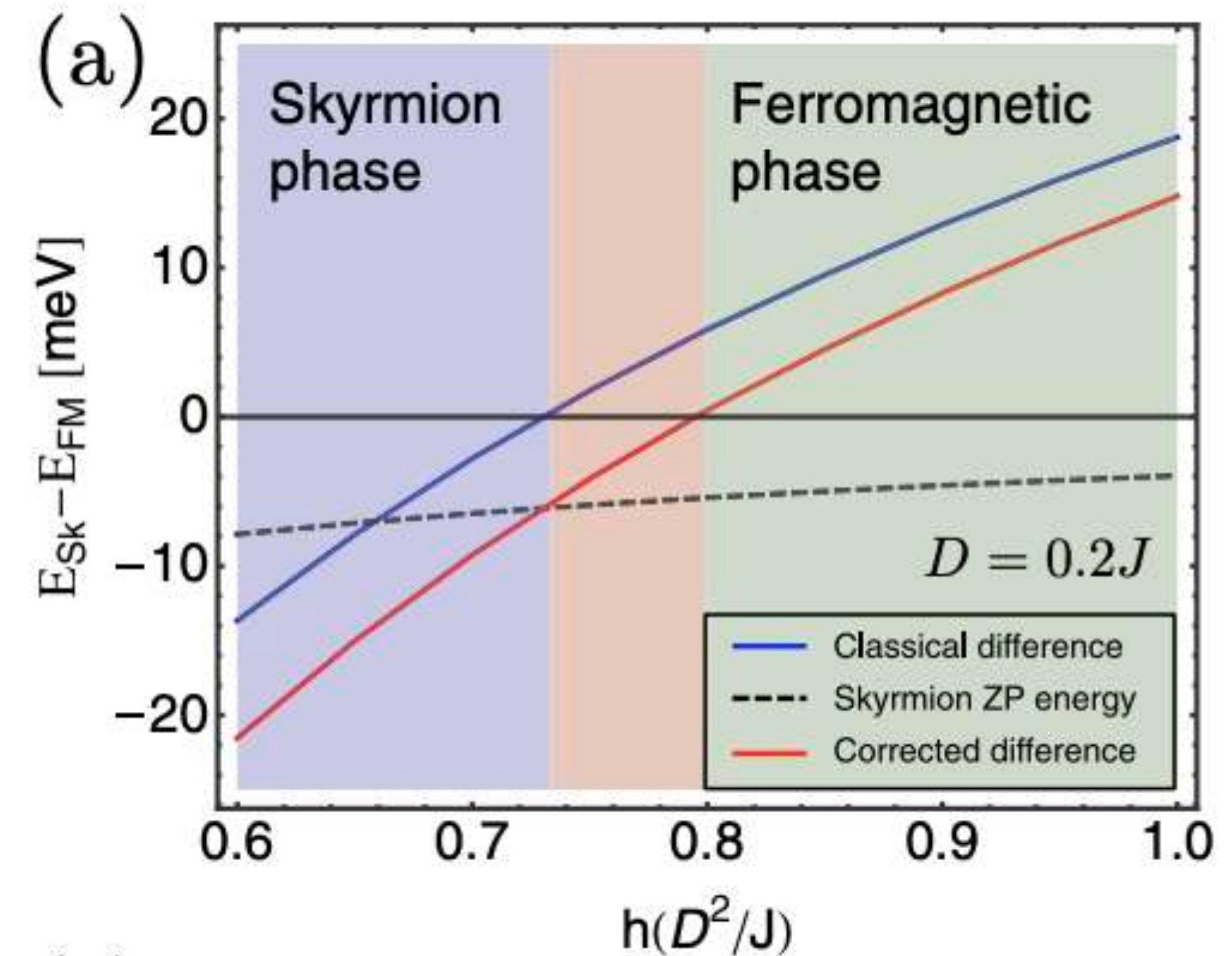
PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations

$$J = -0.5D \quad B = -B_0 e_z$$

$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r,r'} \right)$$

$$D_{a_i} = D e_z \times a_i$$



- Bogoliubov transformation for bosons

$$\hat{H}_{SW} = E_0 + \sum_{\nu=1}^N \hbar \omega_{\nu} \left(\alpha_{\nu}^{\dagger} \alpha_{\nu} + \frac{1}{2} \right)$$

Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations
- Lohani et al. (PRX): frustrated AF lattices can host quasi-classical skyrmions

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations
- Lohani et al. (PRX): frustrated AF lattices can host quasi-classical skyrmions

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

$$\bullet \quad |\{\phi_i\}, \{\theta_i\}\rangle = \prod_i e^{-i\phi_i \hat{S}_i^z} e^{i\theta_i \hat{S}_i^y} e^{+i\phi_i \hat{S}_i^z} |FM\rangle$$

Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations
- Lohani et al. (PRX): frustrated AF lattices can host quasi-classical skyrmions

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

- $|\{\phi_i\}, \{\theta_i\}\rangle = \prod_i e^{-i\phi_i \hat{S}_i^z} e^{i\theta_i \hat{S}_i^y} e^{+i\phi_i \hat{S}_i^z} |FM\rangle$
- Skyrmion states are superpositions of classical configurations

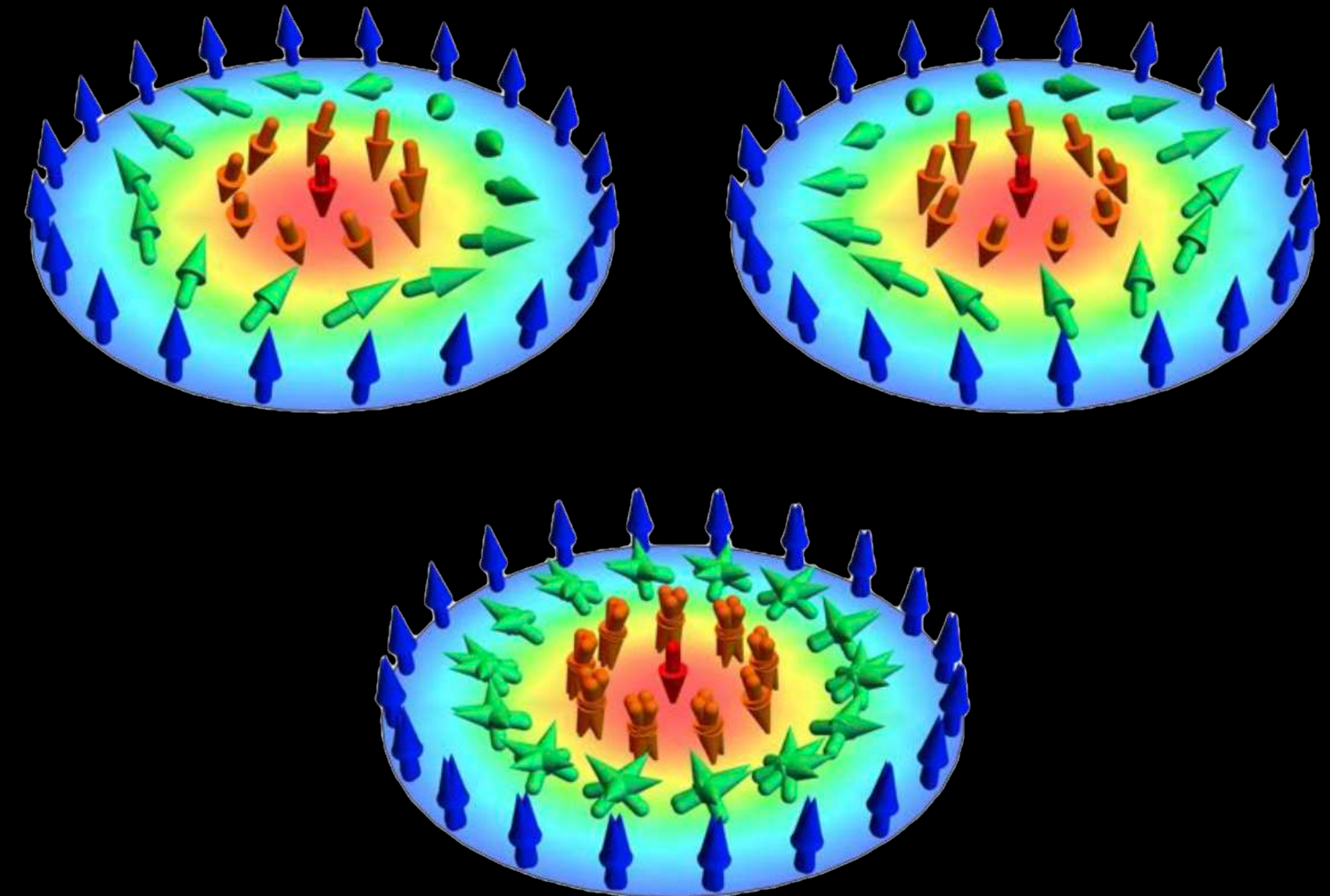
Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations
- Lohani et al. (PRX): frustrated AF lattices can host quasi-classical skyrmions

$$J = -0.5D \quad B = -B_0 \mathbf{e}_z$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{B} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$\mathbf{D}_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

- $|\{\phi_i\}, \{\theta_i\}\rangle = \prod_i e^{-i\phi_i \hat{S}_i^z} e^{i\theta_i \hat{S}_i^y} e^{+i\phi_i \hat{S}_i^z} |FM\rangle$
- Skyrmion states are superpositions of classical configurations



Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations
- Lohani et al. (PRX): frustrated AF lattices can host quasi-classical skyrmions

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations
- Lohani et al. (PRX): frustrated AF lattices can host quasi-classical skyrmions
- Siegl et al. (arXiv preprint): 3x3 quantum system coupled to classical ferromagnet gives quantum analog of skyrmions

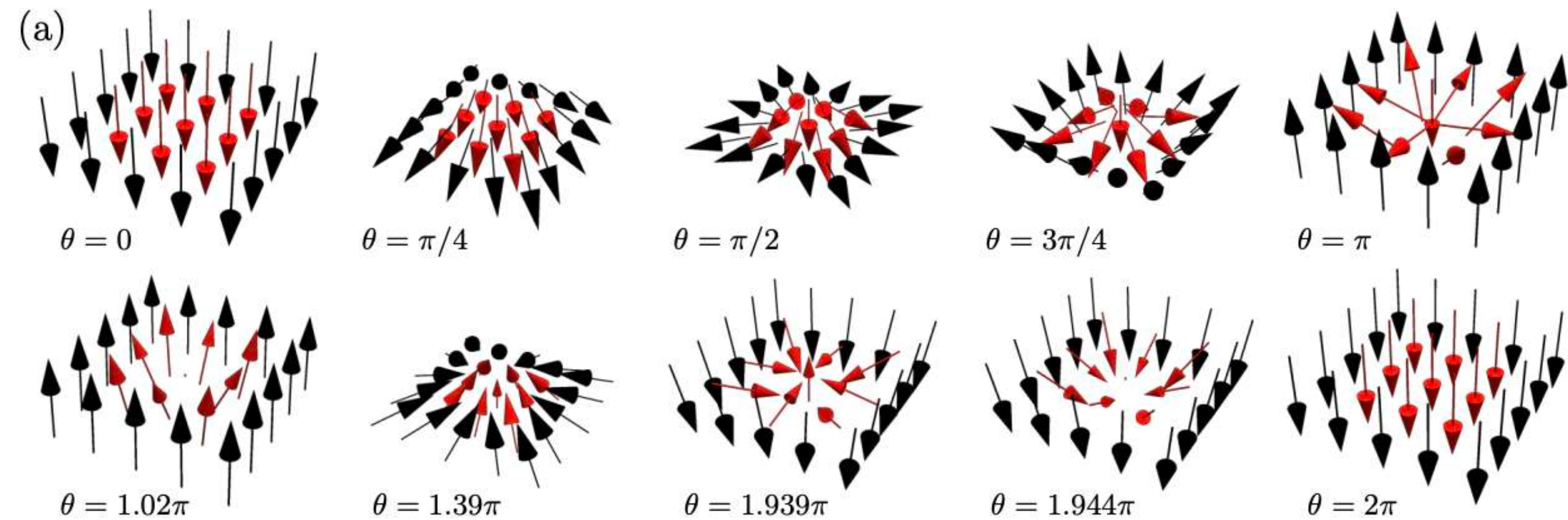
$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations
- Lohani et al. (PRX): frustrated AF lattices can host quasi-classical skyrmions
- Siegl et al. (arXiv preprint): 3x3 quantum system coupled to classical ferromagnet gives quantum analog of skyrmions

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$



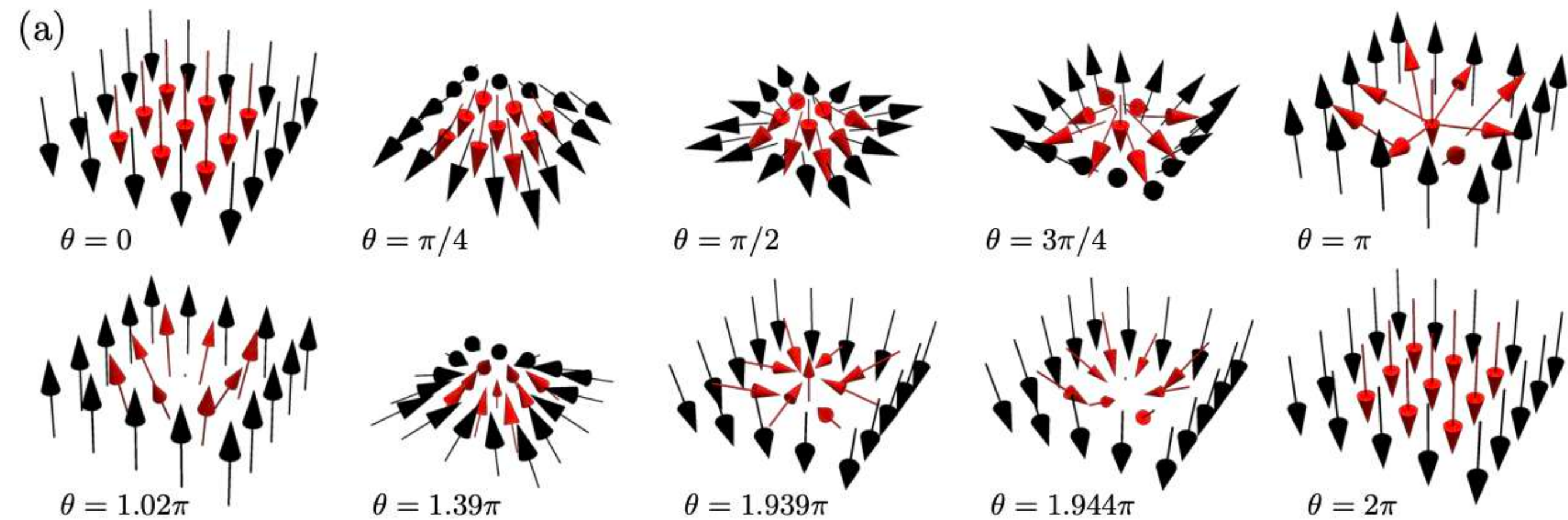
- quantum system (red) coupled to classical magnetic environment (black)

Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations
- Lohani et al. (PRX): frustrated AF lattices can host quasi-classical skyrmions
- Siegl et al. (arXiv preprint): 3x3 quantum system coupled to classical ferromagnet gives quantum analog of skyrmions

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r,r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$



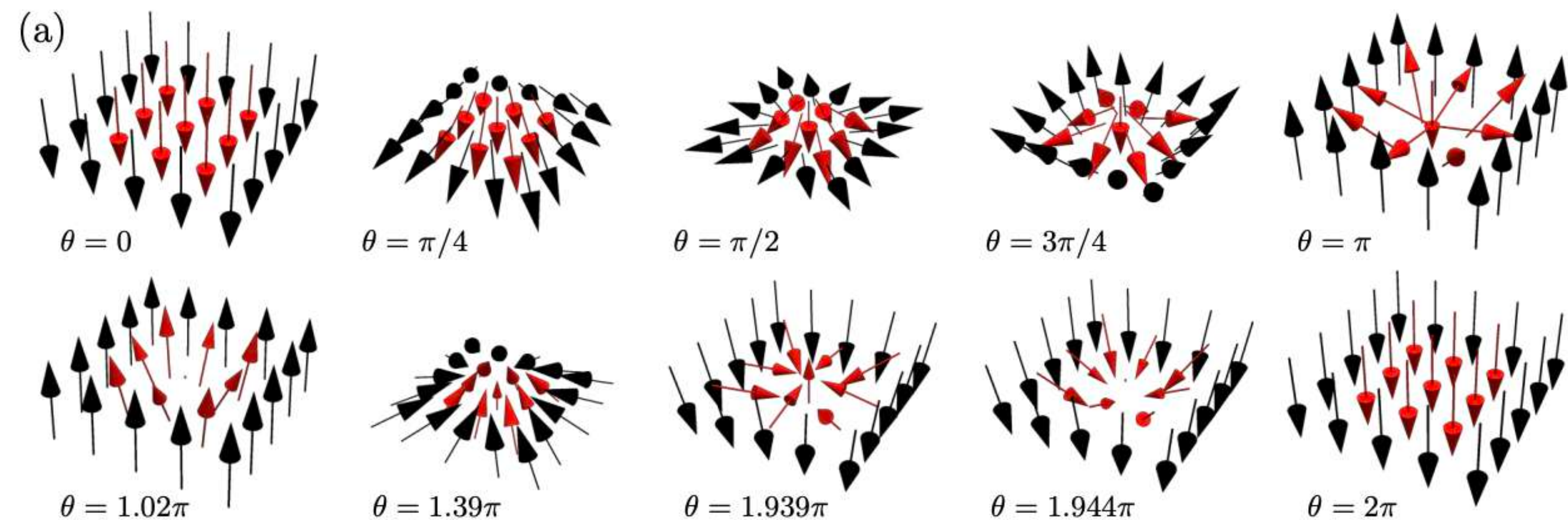
- quantum system (red) coupled to classical magnetic environment (black)
- adiabatic twisting of the environment results in a tiny quantum skyrmion

Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations
- Lohani et al. (PRX): frustrated AF lattices can host quasi-classical skyrmions
- Siegl et al. (arXiv preprint): 3x3 quantum system coupled to classical ferromagnet gives quantum analog of skyrmions

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$



- quantum system (red) coupled to classical magnetic environment (black)
- adiabatic twisting of the environment results in a tiny quantum skyrmion
- controlled creation and annihilation of quantum skyrmions

Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations
- Lohani et al. (PRX): frustrated AF lattices can host quasi-classical skyrmions
- Siegl et al. (arXiv preprint): 3x3 quantum system coupled to classical ferromagnet gives quantum analog of skyrmions

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

Quantum Skyrmions

PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations
- Lohani et al. (PRX): frustrated AF lattices can host quasi-classical skyrmions
- Siegl et al. (arXiv preprint): 3x3 quantum system coupled to classical ferromagnet gives quantum analog of skyrmions
- Sotnikov et al. (PRB): ED study of quantum analog of classical Hamiltonian, but system sizes too small to observe real-space structure

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

Quantum Skyrmions

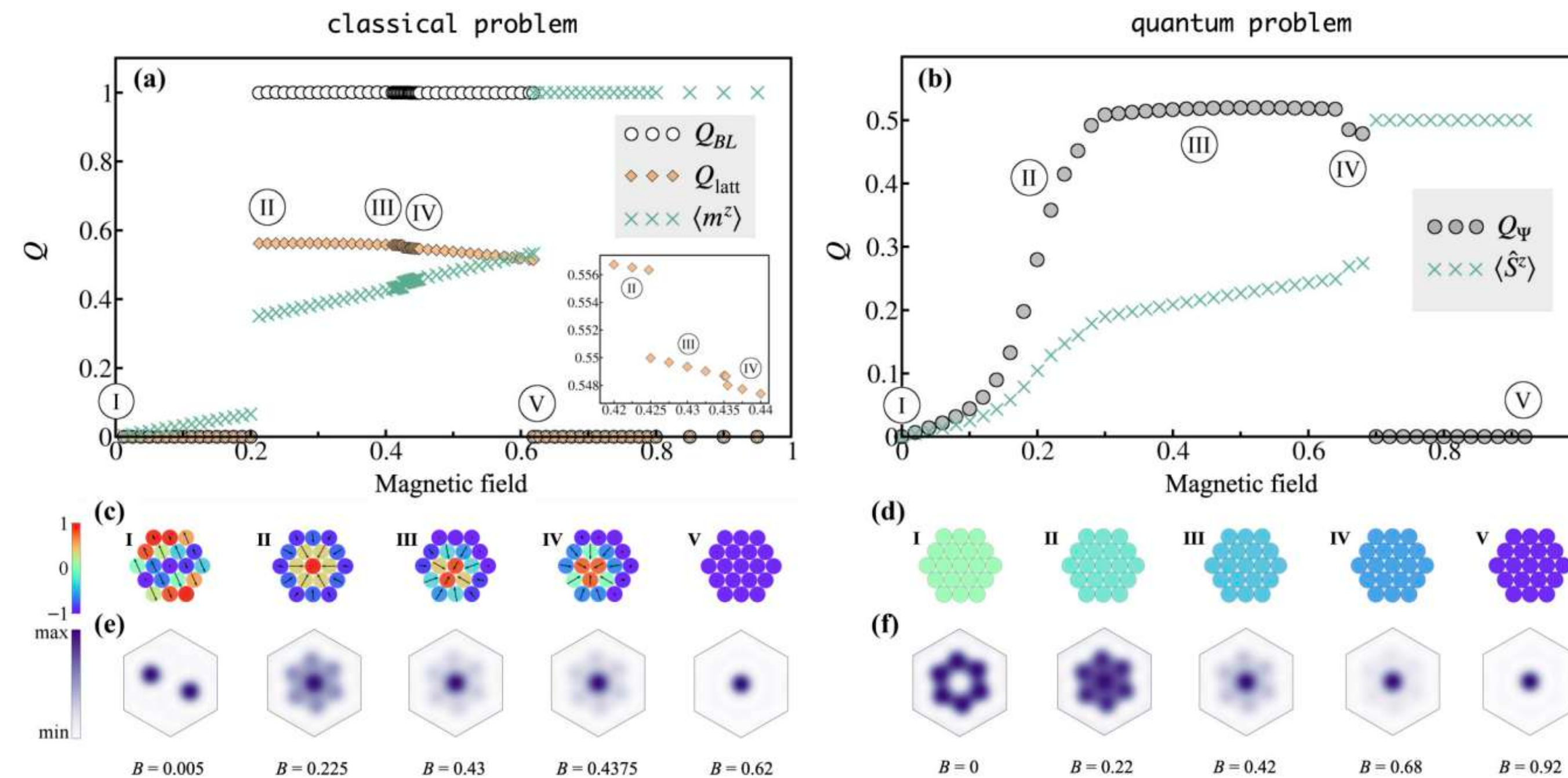
PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations
- Lohani et al. (PRX): frustrated AF lattices can host quasi-classical skyrmions
- Siegl et al. (arXiv preprint): 3x3 quantum system coupled to classical ferromagnet gives quantum analog of skyrmions
- Sotnikov et al. (PRB): ED study of quantum analog of classical Hamiltonian, but system sizes too small to observe real-space structure

$$J = -0.5D \quad B = -B_0 e_z$$

$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r,r'} \right)$$

$$D_{a_i} = D e_z \times a_i$$



Quantum Skyrmions

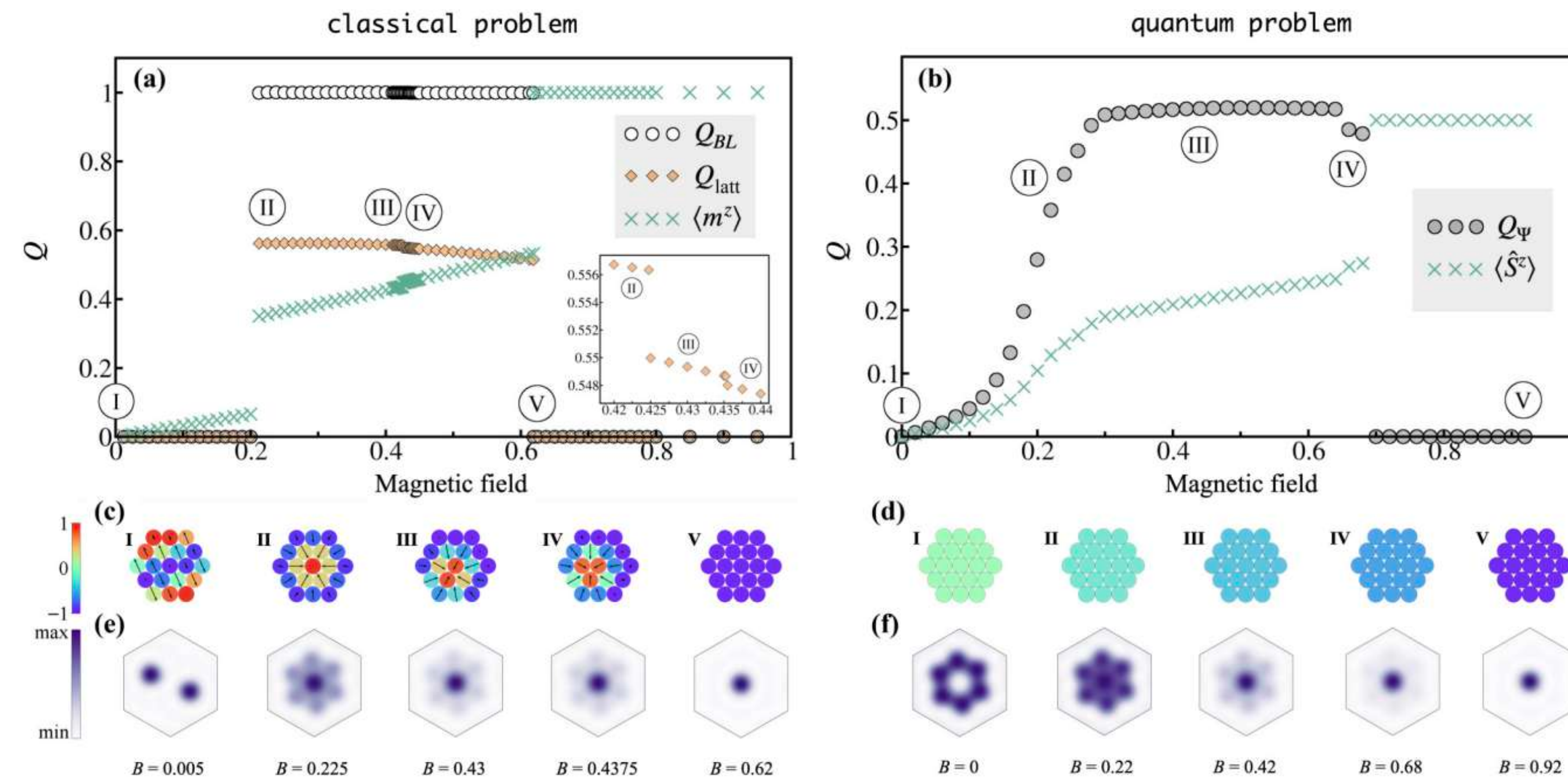
PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations
- Lohani et al. (PRX): frustrated AF lattices can host quasi-classical skyrmions
- Siegl et al. (arXiv preprint): 3x3 quantum system coupled to classical ferromagnet gives quantum analog of skyrmions
- Sotnikov et al. (PRB): ED study of quantum analog of classical Hamiltonian, but system sizes too small to observe real-space structure

$$J = -0.5D \quad B = -B_0 e_z$$

$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r,r'} \right)$$

$$D_{a_i} = D e_z \times a_i$$



- Skyrmion states are superpositions of classical configurations, skyrmion structure revealed in spin structure factor

Quantum Skyrmions

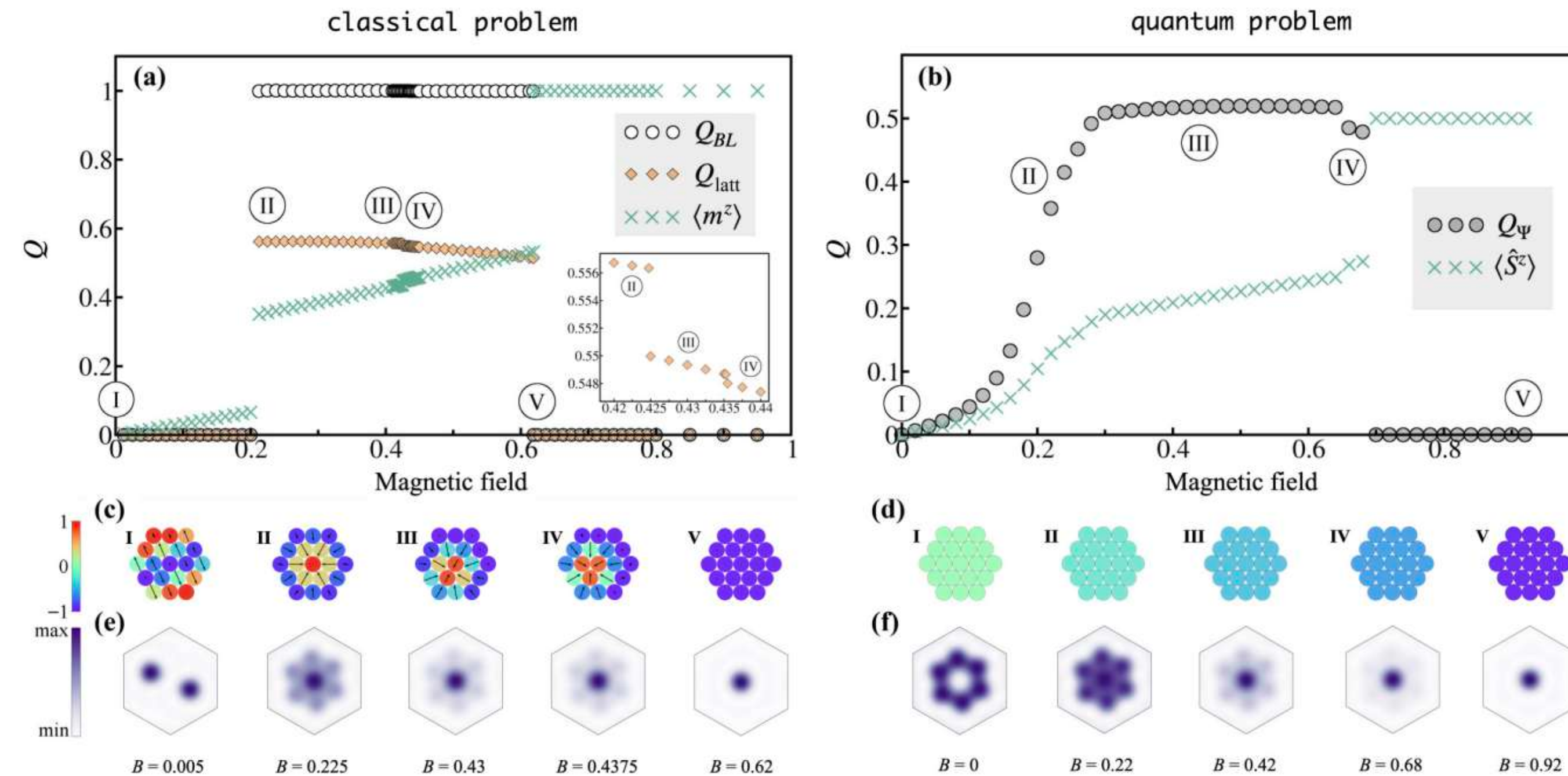
PRB 92, 245436 (2015), PRX 9, 041063 (2019),
PRB 103, L060404 (2021), arXiv:2110.00348(2021)

- Molina et al.: skyrmions are stabilized by quantum fluctuations
- Lohani et al. (PRX): frustrated AF lattices can host quasi-classical skyrmions
- Siegl et al. (arXiv preprint): 3x3 quantum system coupled to classical ferromagnet gives quantum analog of skyrmions
- Sotnikov et al. (PRB): ED study of quantum analog of classical Hamiltonian, but system sizes too small to observe real-space structure

$$J = -0.5D \quad B = -B_0 e_z$$

$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r,r'} \right)$$

$$D_{a_i} = D e_z \times a_i$$



- Skyrmion states are superpositions of classical configurations, skyrmion structure revealed in spin structure factor
- No clean topological classification possible, because invariant quantization breaks down for quantum skyrmions

Quantum Skyrmions

arXiv:2112.12475*

$$\begin{aligned} J &= -0.5D & B &= -B_0 \mathbf{e}_z \\ \hat{H} &= J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{B} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right) \\ \mathbf{D}_{\mathbf{a}_i} &= D \mathbf{e}_z \times \mathbf{a}_i \end{aligned}$$

Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)

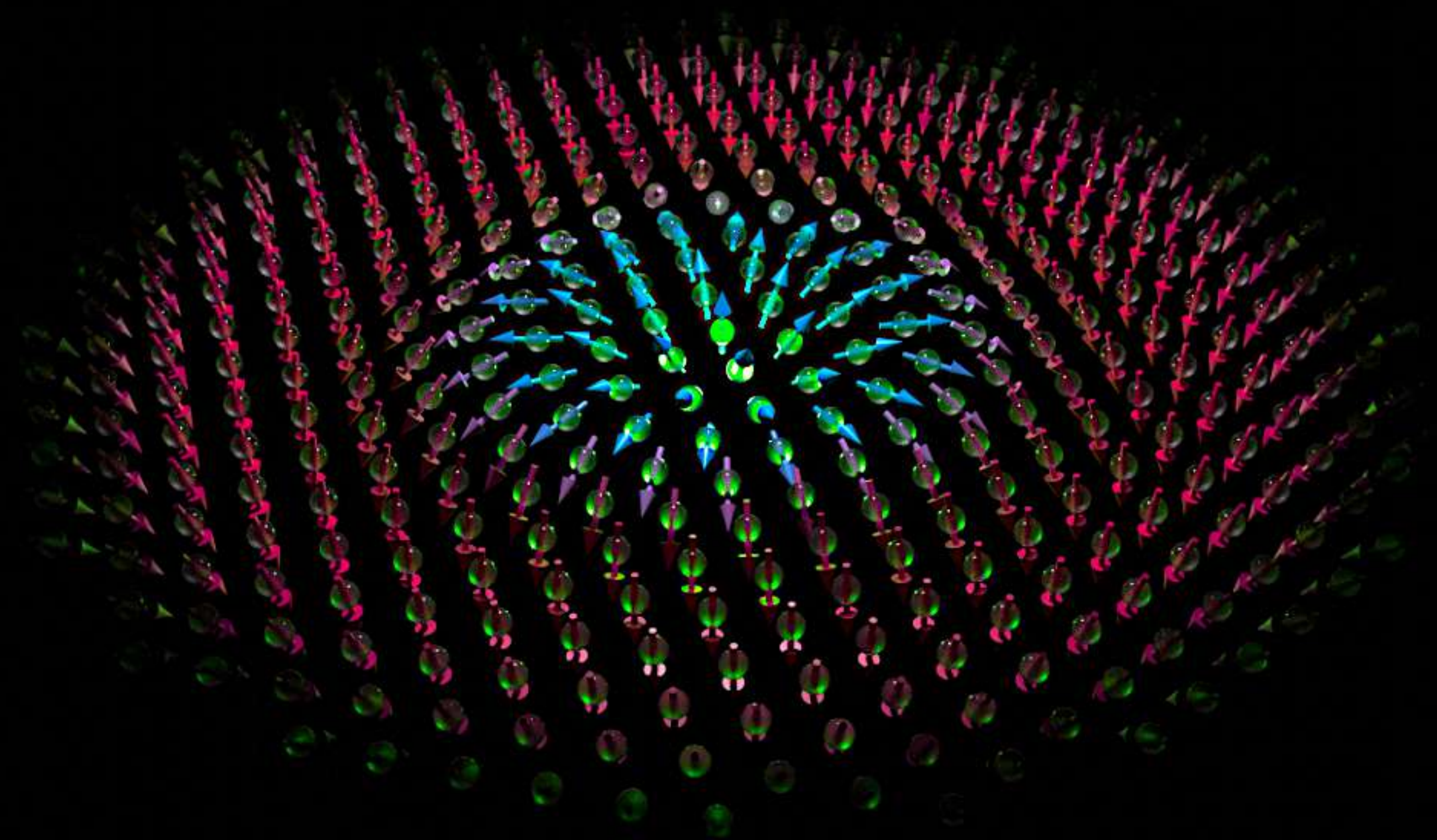
$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)

$$J = -0.5D \quad B = -B_0 \mathbf{e}_z$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{B} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$\mathbf{D}_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

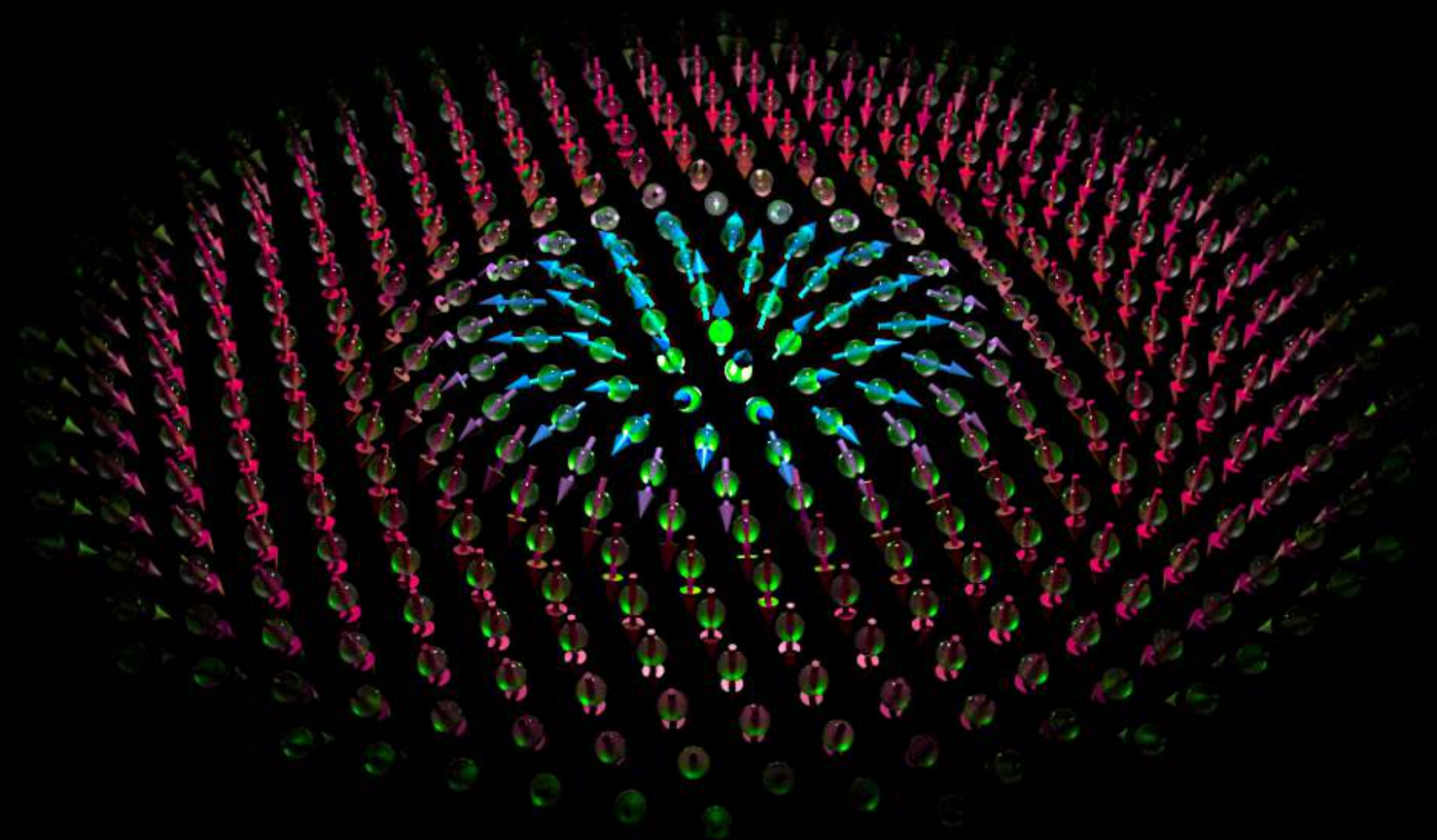


Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- **Quantum skyrmions are *entangled*** superpositions of classical states

$$J = -0.5D \quad B = -B_0 \mathbf{e}_z$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{B} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$\mathbf{D}_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

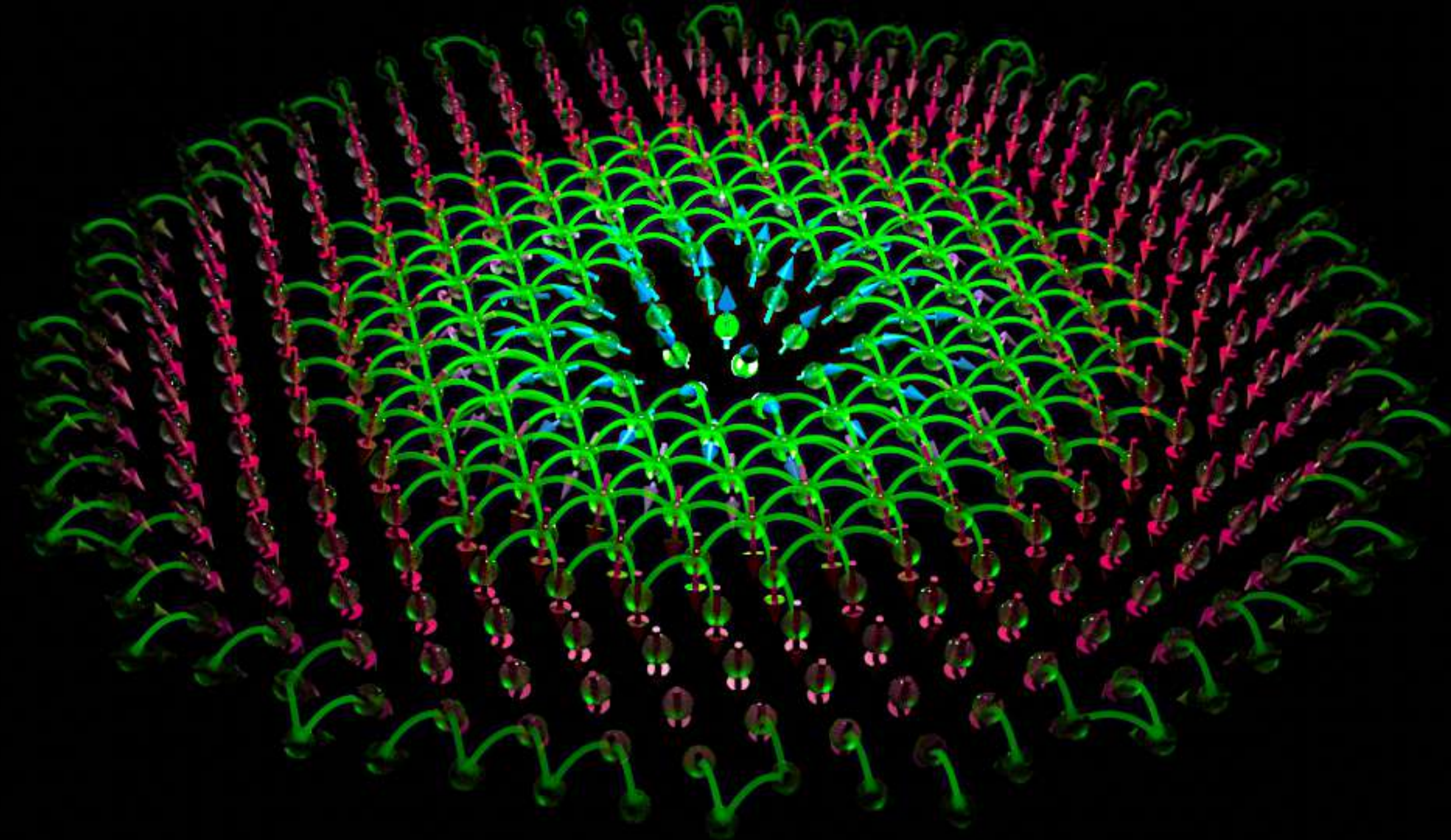


Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- **Quantum skyrmions are *entangled*** superpositions of classical states

$$J = -0.5D \quad B = -B_0 \mathbf{e}_z$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{B} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$\mathbf{D}_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

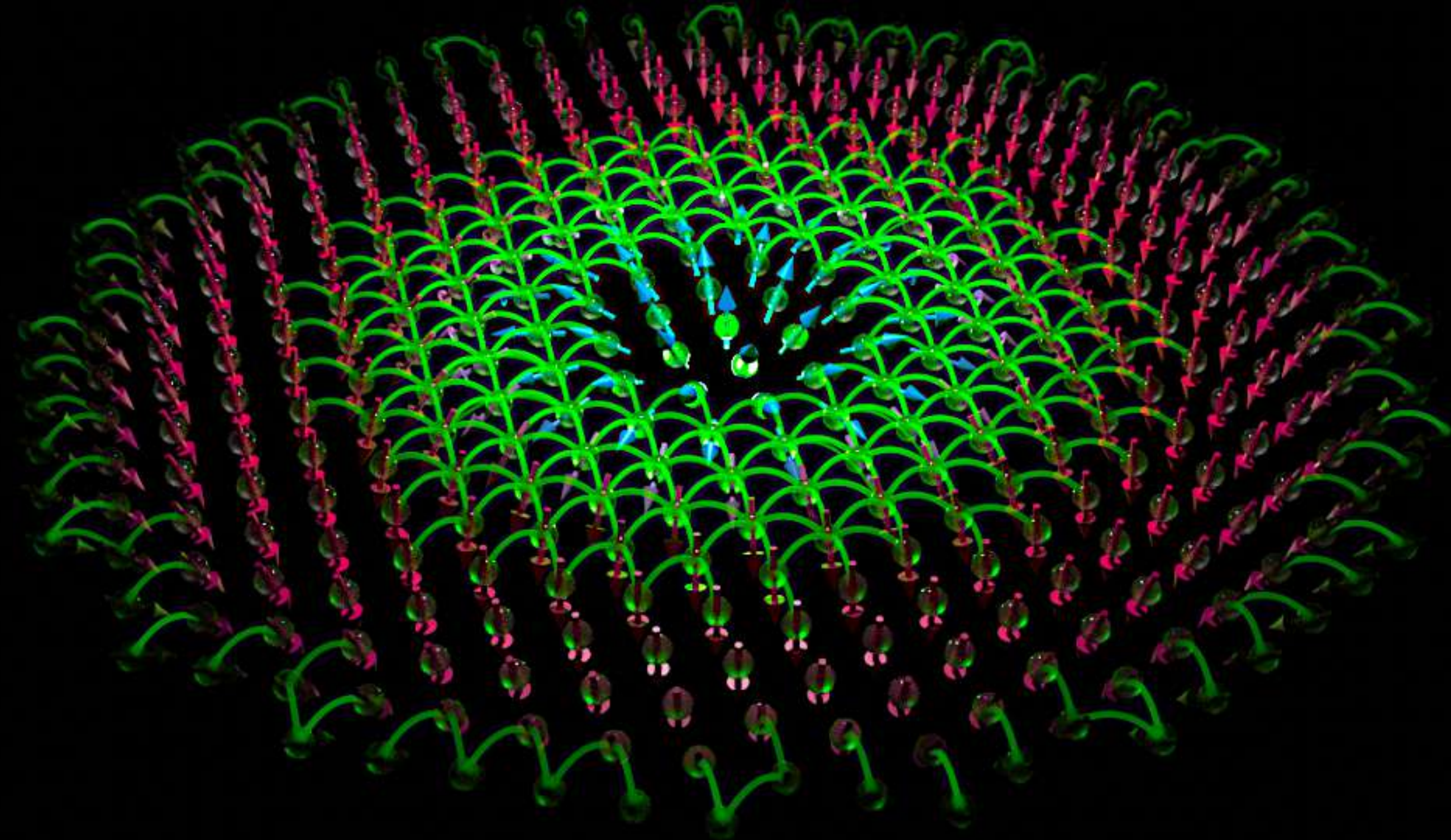


Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- **Quantum skyrmions are *entangled*** superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section

$$J = -0.5D \quad B = -B_0 \mathbf{e}_z$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{B} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$\mathbf{D}_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

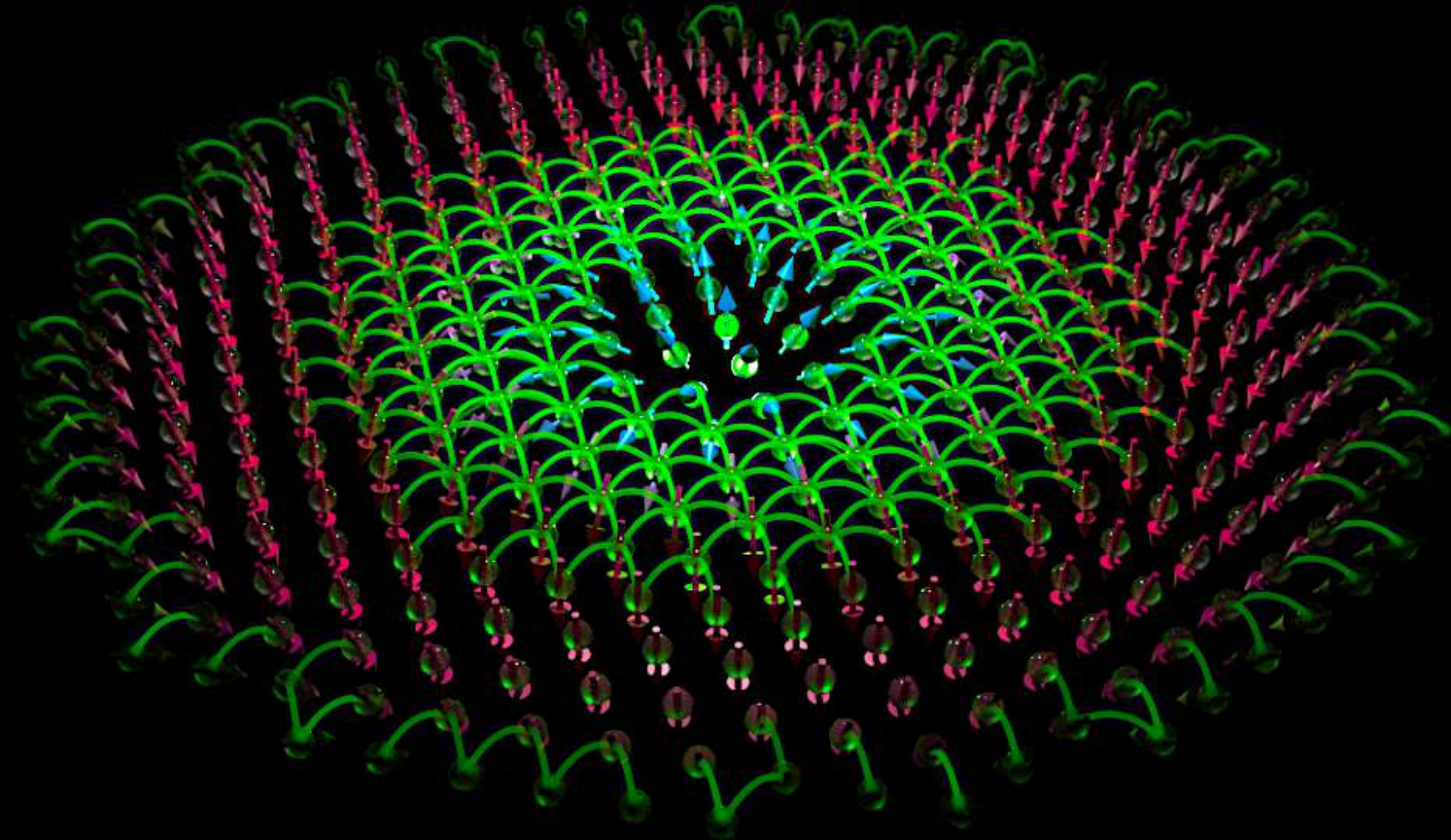


Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- **Quantum skyrmions are *entangled*** superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- **Quantum skyrmion lattices** in the thermodynamic limit

$$J = -0.5D \quad B = -B_0 \mathbf{e}_z$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{B} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$\mathbf{D}_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$



Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- *Quantum skyrmions are entangled* superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- *Quantum skyrmion lattices* in the thermodynamic limit

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

- matrix product states (MPS)

$$|\psi\rangle = \sum_{i_1, i_2, \dots, i_L} C(i_1, i_2, \dots, i_L) |i_1, i_2, \dots, i_L\rangle$$

Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- *Quantum skyrmions are entangled* superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- *Quantum skyrmion lattices* in the thermodynamic limit

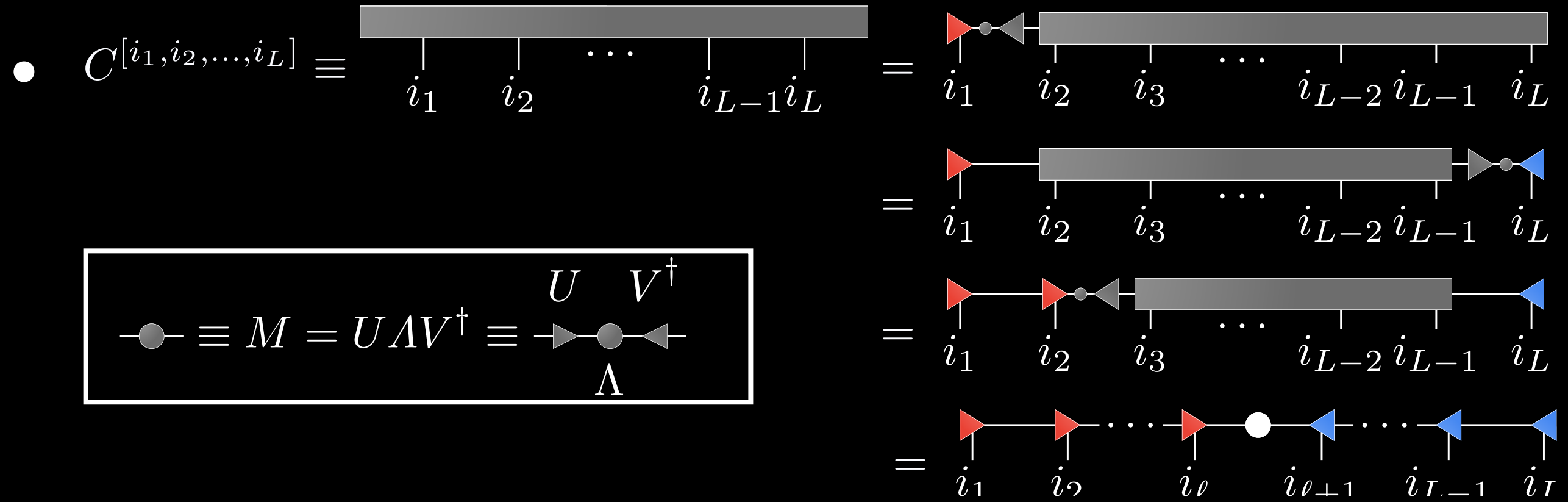
$$J = -0.5D \qquad B = -B_0 e_z$$

$$\hat{H} = J \sum_{\langle r,r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r,r'} \right)$$

$$D_{a_i} = D e_z \times a_i$$

- matrix product states (MPS)

$$|\psi\rangle = \sum_{i_1,i_2,\dots,i_L} C(i_1,i_2,\dots,i_L) |i_1,i_2,\dots,i_L\rangle$$



arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- *Quantum skyrmions are entangled* superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- *Quantum skyrmion lattices* in the thermodynamic limit

$$\begin{aligned} J &= -0.5D & \mathbf{B} &= -B_0 \mathbf{e}_z \\ \hat{H} &= J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{B} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right) \\ \mathbf{D}_{\mathbf{a}_i} &= D \mathbf{e}_z \times \mathbf{a}_i \end{aligned}$$

- matrix product states (MPS)

- $|\psi\rangle = \sum_{i_1, i_2, \dots, i_L} C(i_1, i_2, \dots, i_L) |i_1, i_2, \dots, i_L\rangle$

$\bullet \quad C^{[i_1, i_2, \dots, i_L]} \equiv$

- every quantum state has an MPS form

Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- *Quantum skyrmions are entangled* superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- *Quantum skyrmion lattices* in the thermodynamic limit

$$J = -0.5D \quad B = -B_0 \mathbf{e}_z$$

$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{B} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$

$$\mathbf{D}_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

- matrix product states (MPS)

$$|\psi\rangle = \sum_{i_1, i_2, \dots, i_L} C(i_1, i_2, \dots, i_L) |i_1, i_2, \dots, i_L\rangle$$

$$|\psi\rangle = \sum_{\mathbf{i}} C^{[\mathbf{i}]} |\mathbf{i}\rangle = \sum_k s_k(\ell) |\psi_{A(\ell), k}\rangle |\psi_{B(\ell), k}\rangle$$

$$|\psi_{A(\ell), k}\rangle = U_{k_1}^{[i_1]} U_{k_1, k_2}^{[i_2]} U_{k_2, k_3}^{[i_3]} \dots U_{k_{\ell-1}, k}^{[i_{\ell}]} |i_1, i_2, i_3, \dots, i_{\ell}\rangle = \prod_{k=1}^{\ell} U^{[i_k]} |i_1, \dots, i_{\ell}\rangle$$

$$|\psi_{B(\ell), k}\rangle = V_{k, k_{\ell+1}}^{\dagger [i_{\ell+1}]} V_{k_{\ell+1}, k_{\ell+2}}^{\dagger [i_{\ell+2}]} \dots V_{k_L, k}^{\dagger [i_L]} |i_{\ell+1}, \dots, i_L\rangle = \prod_{k=\ell+1}^L V^{\dagger [i_k]} |i_{\ell+1}, \dots, i_L\rangle$$

Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- *Quantum skyrmions are entangled* superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- *Quantum skyrmion lattices* in the thermodynamic limit

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

- matrix product states (MPS)
- $|\psi\rangle = \sum_{i_1, i_2, \dots, i_L} C(i_1, i_2, \dots, i_L) |i_1, i_2, \dots, i_L\rangle$
- $|\psi\rangle = \sum_i C^{[i]} |i\rangle = \sum_k s_k(\ell) |\psi_{A(\ell), k}\rangle |\psi_{B(\ell), k}\rangle$
- MPS form gives access to the eigenvalues of the reduced density matrix
 - measure bipartite entanglement for free

$$S_\alpha = \frac{1}{1-\alpha} \log_2 \text{tr} \hat{\rho}(\ell)^\alpha = \frac{1}{1-\alpha} \log_2 \sum_k s_k(\ell)^\alpha$$

Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- *Quantum skyrmions are entangled* superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- *Quantum skyrmion lattices* in the thermodynamic limit

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

- matrix product states (MPS)
- $|\psi\rangle = \sum_{i_1, i_2, \dots, i_L} C(i_1, i_2, \dots, i_L) |i_1, i_2, \dots, i_L\rangle$
- $|\psi\rangle = \sum_i C^{[i]} |i\rangle = \sum_{k=1}^{\infty} s_k(\ell) |\psi_{A(\ell), k}\rangle |\psi_{B(\ell), k}\rangle$
- canonical truncation: keep the largest contributions to reduced density matrix

Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- *Quantum skyrmions are entangled* superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- *Quantum skyrmion lattices* in the thermodynamic limit

$$J = -0.5D \quad B = -B_0 \mathbf{e}_z$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{B} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$\mathbf{D}_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

- matrix product states (MPS)
- $|\psi\rangle = \sum_{i_1, i_2, \dots, i_L} C(i_1, i_2, \dots, i_L) |i_1, i_2, \dots, i_L\rangle$
- $|\psi\rangle = \sum_{\mathbf{i}} C^{[\mathbf{i}]} |\mathbf{i}\rangle = \sum_{k=1}^{\chi} s_k(\ell) |\psi_{A(\ell), k}\rangle |\psi_{B(\ell), k}\rangle$
- canonical truncation: keep the largest contributions to reduced density matrix

Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- *Quantum skyrmions are entangled* superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- *Quantum skyrmion lattices* in the thermodynamic limit

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

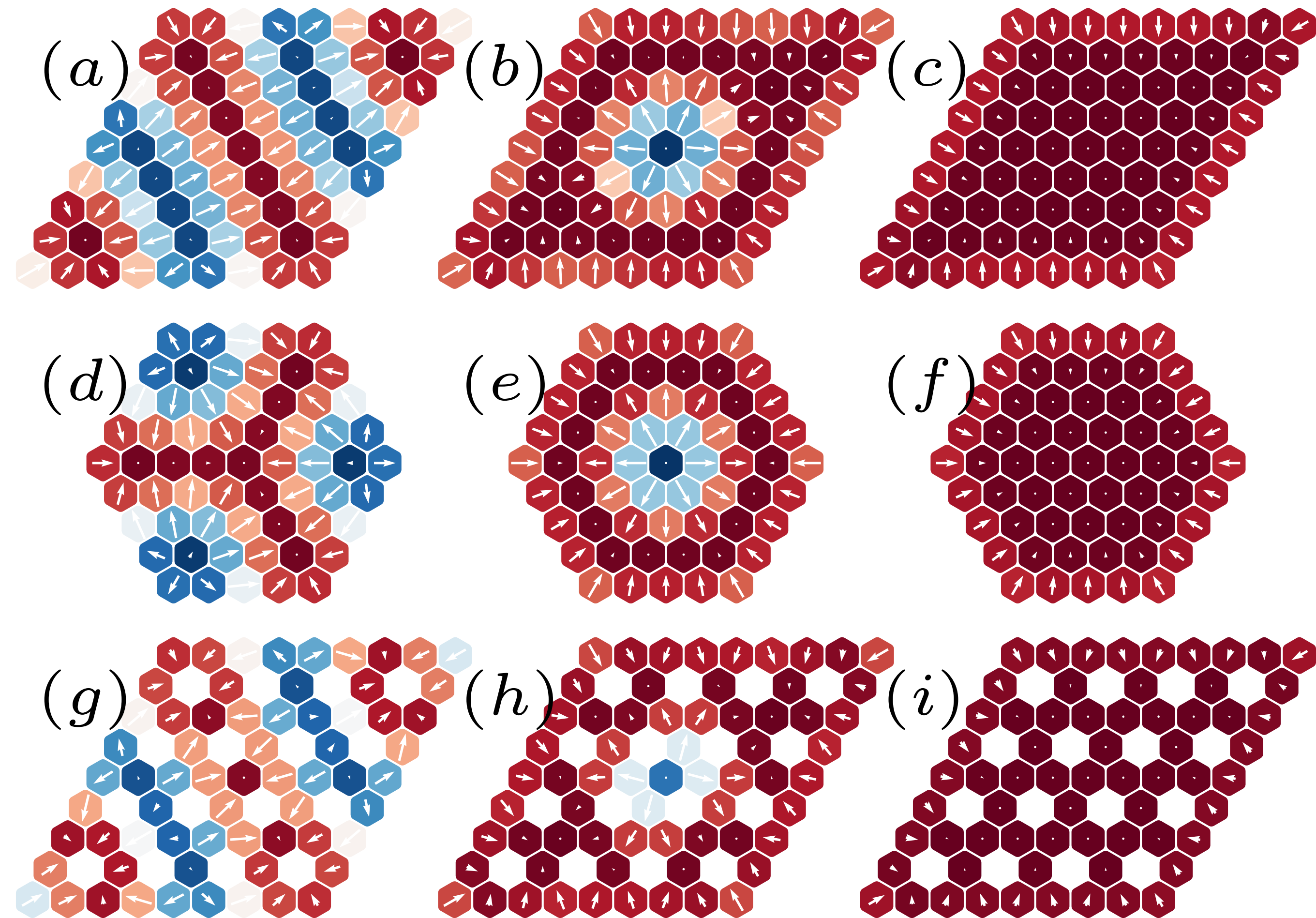
- matrix product states (MPS)
- $|\psi\rangle = \sum_{i_1, i_2, \dots, i_L} C(i_1, i_2, \dots, i_L) |i_1, i_2, \dots, i_L\rangle$
- $|\psi\rangle = \sum_i C^{[i]} |i\rangle = \sum_{k=1}^{\chi} s_k(\ell) |\psi_{A(\ell), k}\rangle |\psi_{B(\ell), k}\rangle$
- canonical truncation: keep the largest contributions to reduced density matrix
- effective variational energy minimization (density matrix renormalization group)

Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- *Quantum skyrmions are entangled* superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- *Quantum skyrmion lattices* in the thermodynamic limit

$$J = -0.5D \quad B = -B_0 \mathbf{e}_z$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{B} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$\mathbf{D}_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

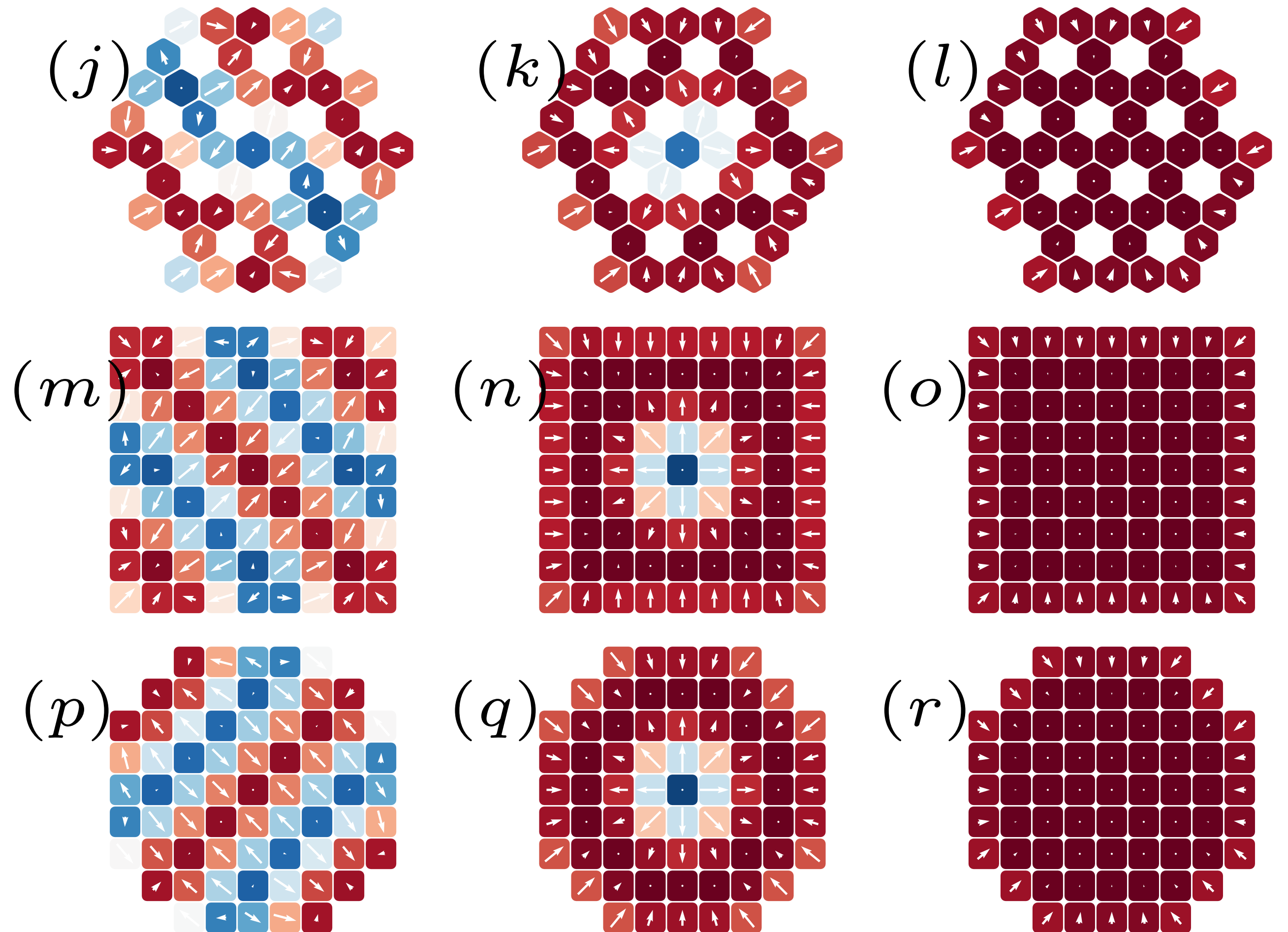


Quantum Skyrmions

arXiv:2112.12475*

$$J = -0.5D \quad B = -B_0 \mathbf{e}_z$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + B \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$\mathbf{D}_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- *Quantum skyrmions are entangled* superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- *Quantum skyrmion lattices* in the thermodynamic limit

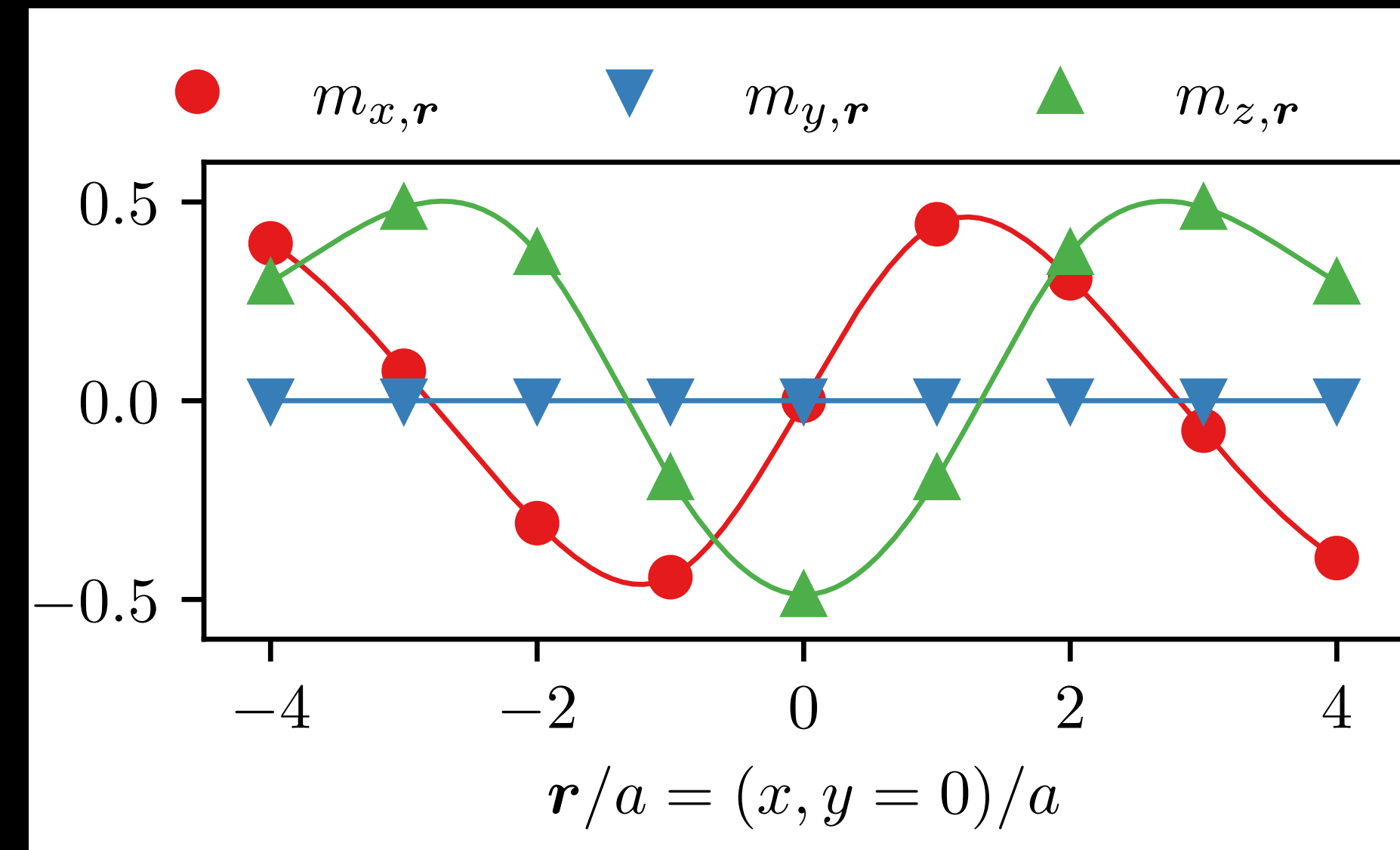


Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- *Quantum skyrmions are entangled* superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- *Quantum skyrmion lattices* in the thermodynamic limit

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{\mathbf{S}}_r \cdot \hat{\mathbf{S}}_{r'} + D_{r-r'} \cdot \hat{\mathbf{S}}_r \times \hat{\mathbf{S}}_{r'} + B \cdot \hat{\mathbf{S}}_r \delta_{r,r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$



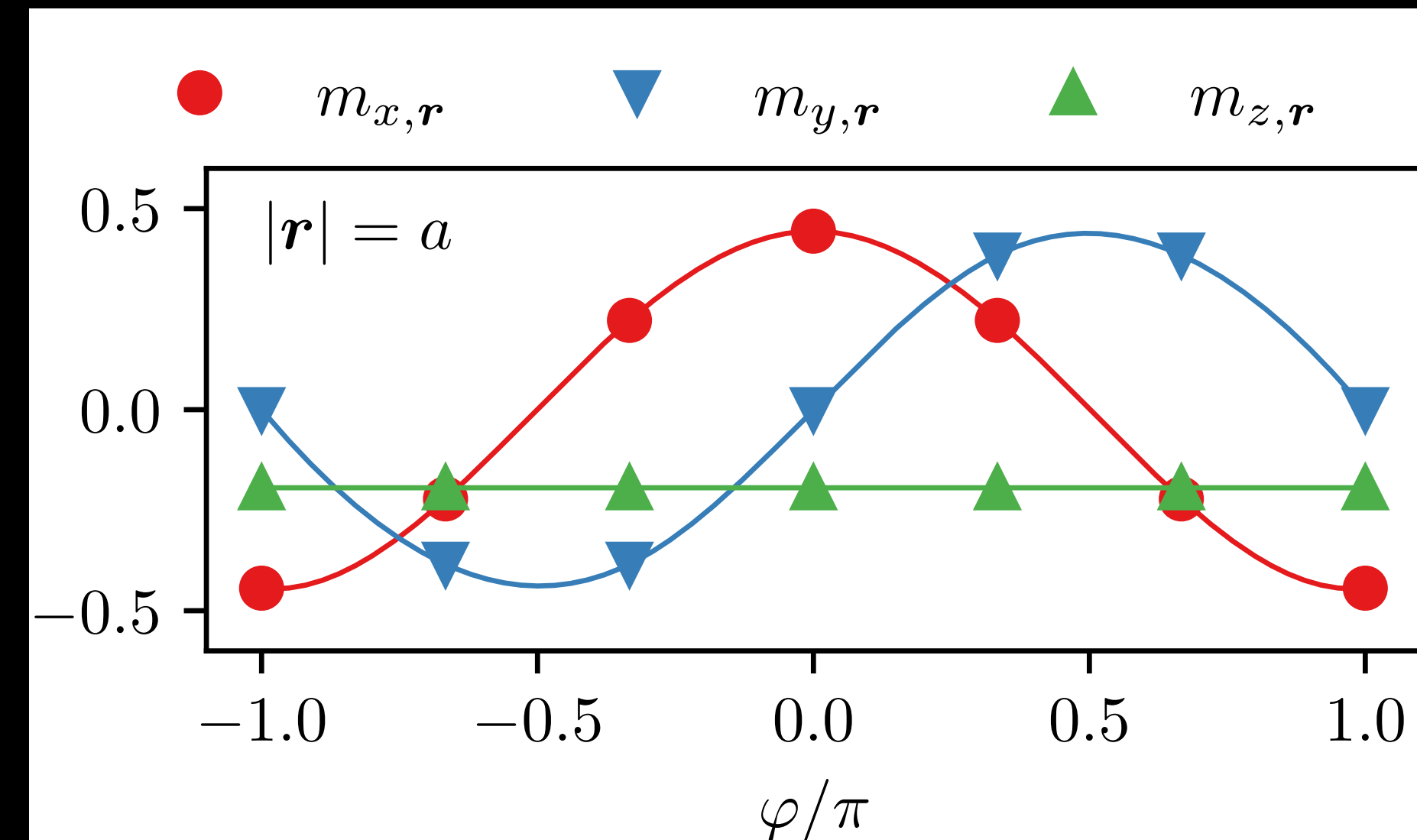
- radial spin winding of 3 lattice sites

Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- *Quantum skyrmions are entangled* superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- *Quantum skyrmion lattices* in the thermodynamic limit

$$J = -0.5D \quad B = -B_0 \mathbf{e}_z$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{B} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$\mathbf{D}_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$



- radial spin winding of 3 lattice sites
- angular spin rotation according to lattice symmetries (here, rotation by $\pi/3$)
- Skyrmion type: Hedgehog family

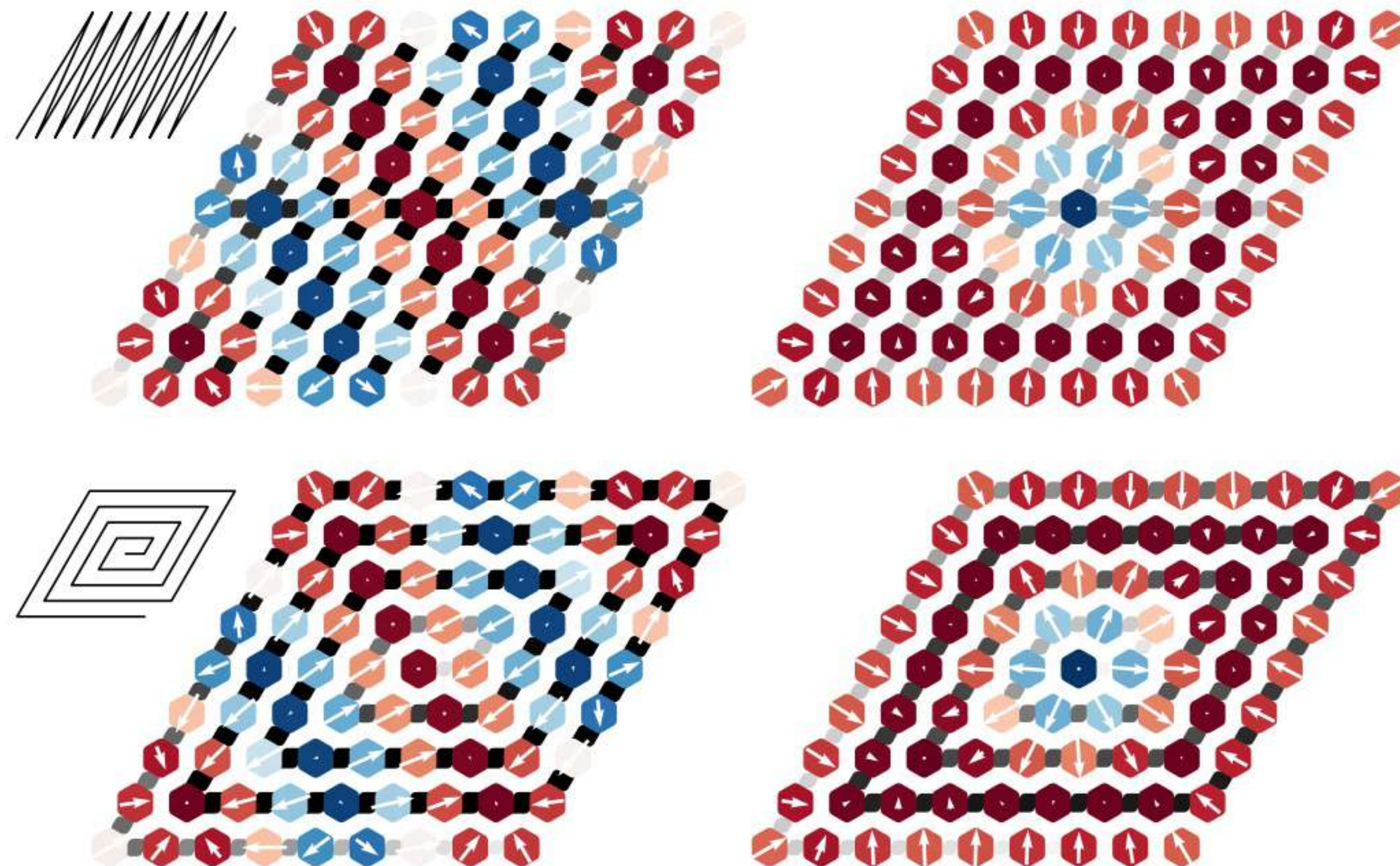
Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- **Quantum skyrmions are entangled** superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- **Quantum skyrmion lattices** in the thermodynamic limit

$$J = -0.5D \quad B = -B_0 \mathbf{e}_z$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + D_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + B \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$D_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

- use MPS form to confirm bipartite entanglement (von Neumann entropy)



Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- **Quantum skyrmions are entangled** superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- **Quantum skyrmion lattices** in the thermodynamic limit

$$J = -0.5D \quad B = -B_0 \mathbf{e}_z$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + B \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$\mathbf{D}_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

- use MPS form to confirm bipartite entanglement (von Neumann entropy)
- confirm entanglement shared between two spins (Osterloh et al., Nature 416)
 - $C_{\mathbf{r}, \mathbf{r}'} = \max(0, \lambda_1 - \lambda_2 - \lambda_3 - \lambda_4)$
 - $R_{\mathbf{r}, \mathbf{r}'} = \rho_{\mathbf{r}, \mathbf{r}'} \tilde{\rho}_{\mathbf{r}, \mathbf{r}'}$ $\rho_{\mathbf{r}, \mathbf{r}'} = \text{tr}_{\{\mathbf{r}, \mathbf{r}'\}^C} \hat{\rho}$
 - eigenvalues $\lambda_1 \geq \lambda_2 \geq \lambda_3 \geq \lambda_4$

Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- **Quantum skyrmions are entangled** superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- **Quantum skyrmion lattices** in the thermodynamic limit

$$J = -0.5D \quad B = -B_0 \mathbf{e}_z$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + D_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + B \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$D_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

- use MPS form to confirm bipartite entanglement (von Neumann entropy)
- confirm entanglement shared between two spins (Osterloh et al., Nature 416)
 - $C_{\mathbf{r}, \mathbf{r}'} = \max(0, \lambda_1 - \lambda_2 - \lambda_3 - \lambda_4)$
 - $R_{\mathbf{r}, \mathbf{r}'} = \rho_{\mathbf{r}, \mathbf{r}'} \tilde{\rho}_{\mathbf{r}, \mathbf{r}'} \quad \rho_{\mathbf{r}, \mathbf{r}'} = \text{tr}_{\{\mathbf{r}, \mathbf{r}'\}^C} \hat{\rho}$
 - eigenvalues $\lambda_1 \geq \lambda_2 \geq \lambda_3 \geq \lambda_4$
- side remark:
$$\rho_{\mathbf{r}, \mathbf{r}'} = \sum_{\alpha, \beta} \langle S_{\mathbf{r}}^{\alpha} S_{\mathbf{r}'}^{\beta} \rangle \sigma_{\alpha} \otimes \sigma_{\beta}$$

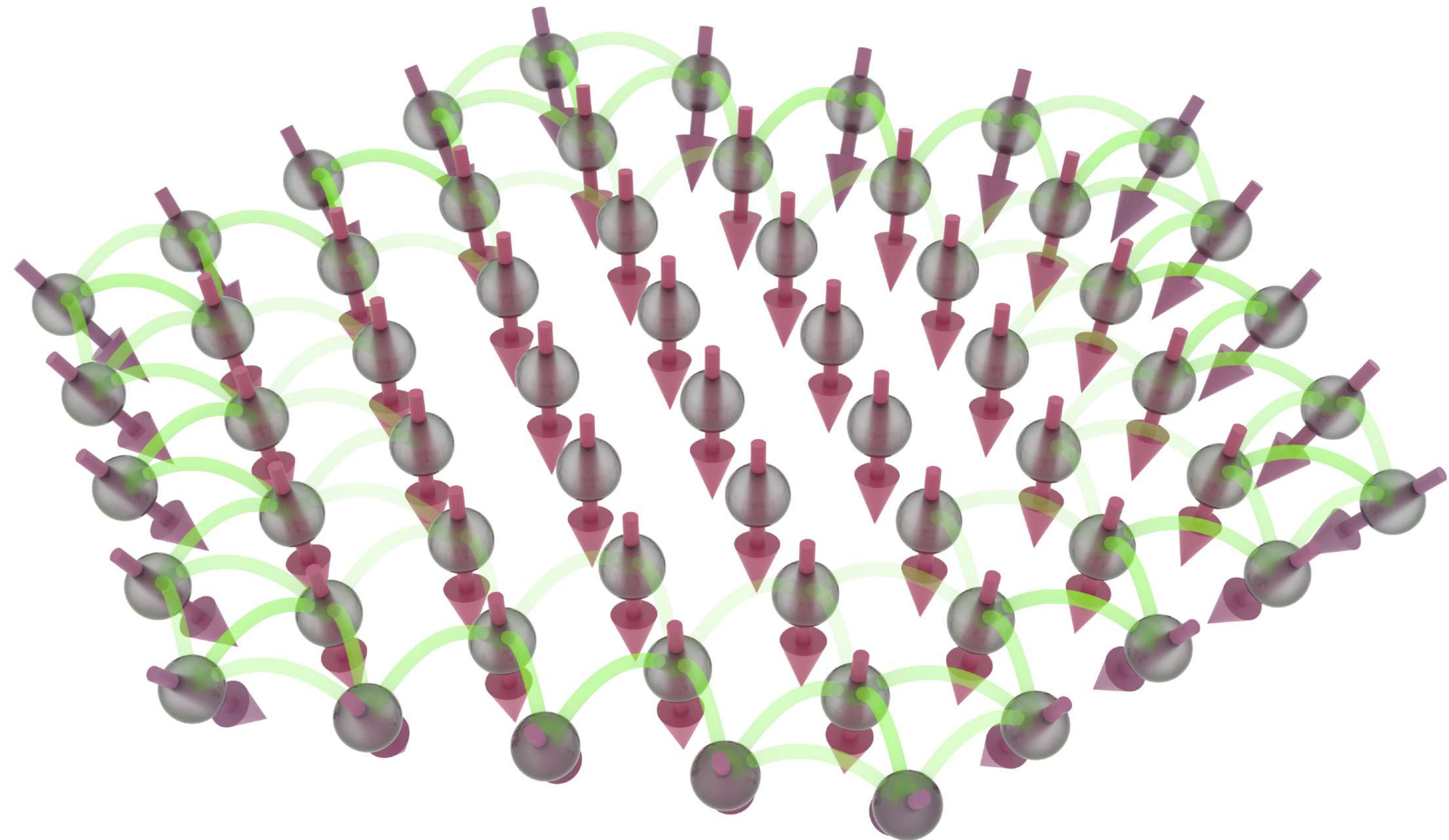
Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- **Quantum skyrmions are entangled** superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- **Quantum skyrmion lattices** in the thermodynamic limit

$$\begin{aligned} B_0 &= 1 \\ J &= -0.5D \\ B &= -B_0 \mathbf{e}_z \\ \hat{H} &= J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + D_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + B \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right) \\ D_{\mathbf{a}_i} &= D \mathbf{e}_z \times \mathbf{a}_i \end{aligned}$$

- confirm entanglement shared between two spins (Osterloh et al., Nature 416)



Quantum Skyrmions

arXiv:2112.12475*

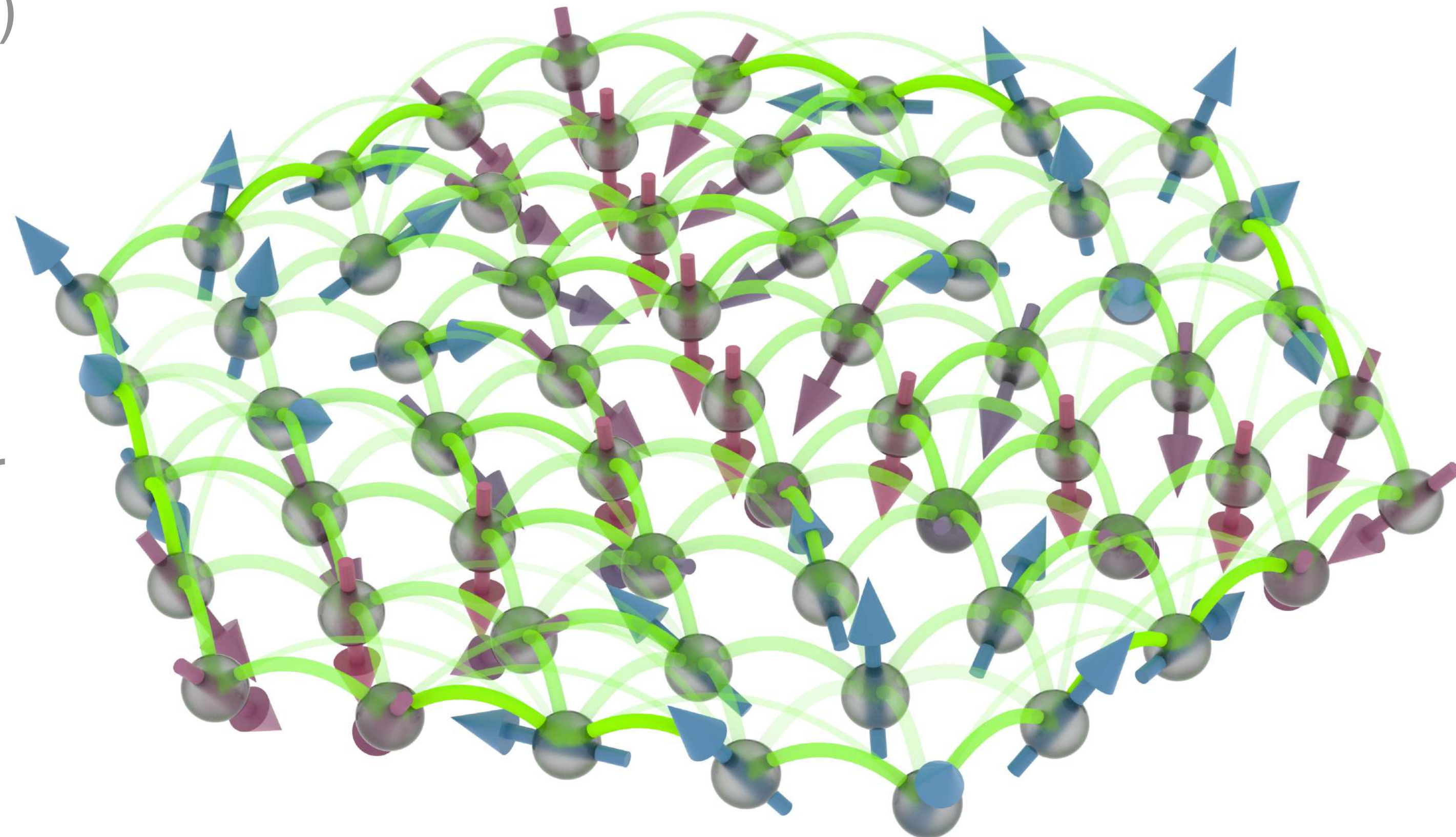
- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- **Quantum skyrmions are entangled** superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- **Quantum skyrmion lattices** in the thermodynamic limit

$$B_0 = 1/8$$

$$B = -B_0 \mathbf{e}_z$$

$$J = -0.5D$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + D_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + B \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$D_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

- confirm entanglement shared between two spins (Osterloh et al., Nature 416)



Quantum Skyrmions

arXiv:2112.12475*

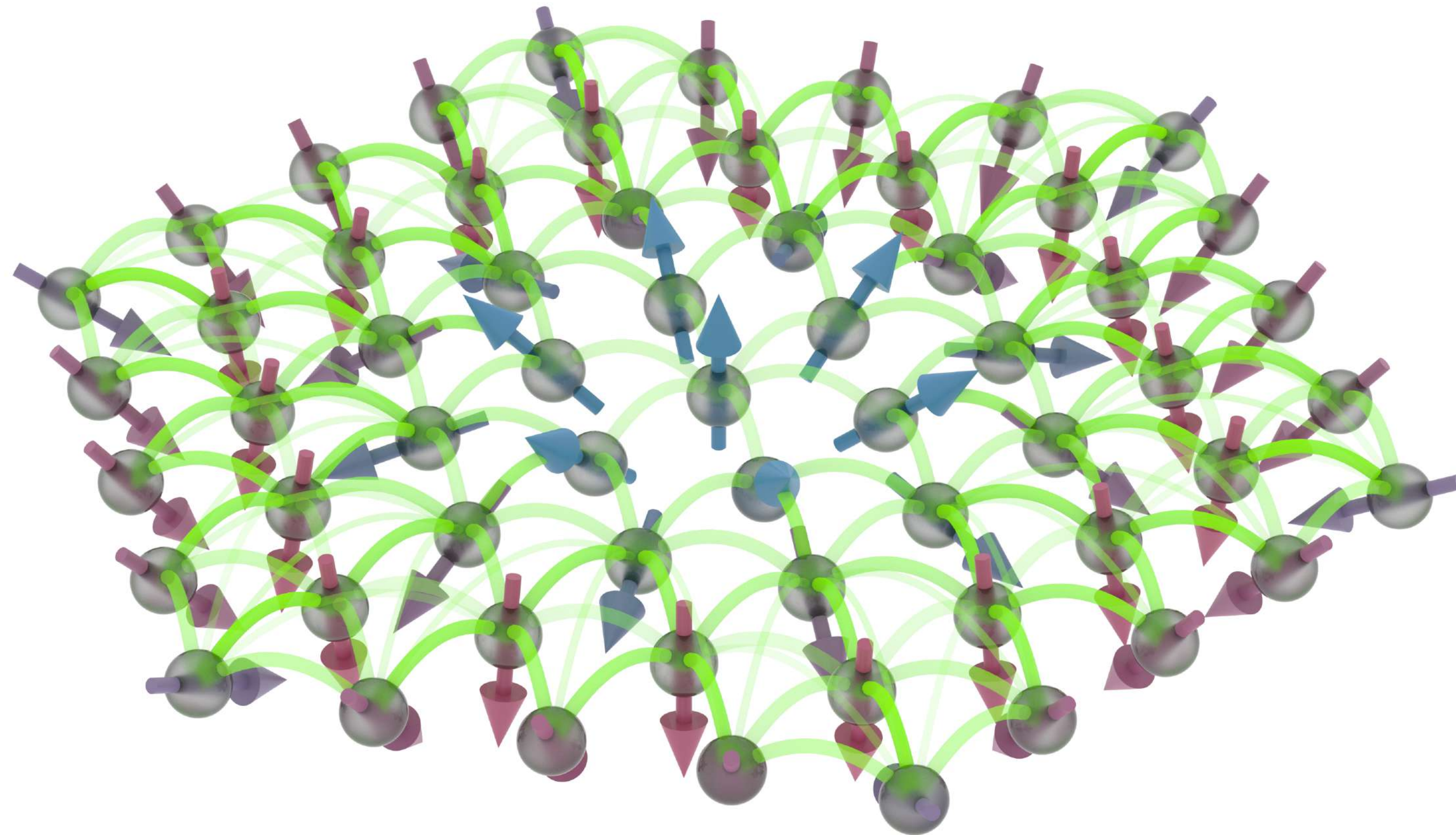
- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- **Quantum skyrmions are entangled** superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- **Quantum skyrmion lattices** in the thermodynamic limit

$$B_0 = 1/2$$

$$B = -B_0 \mathbf{e}_z$$

$$J = -0.5D$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + D_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + B \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$D_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

- confirm entanglement shared between two spins (Osterloh et al., Nature 416)



Quantum Skyrmions

arXiv:2112.12475*

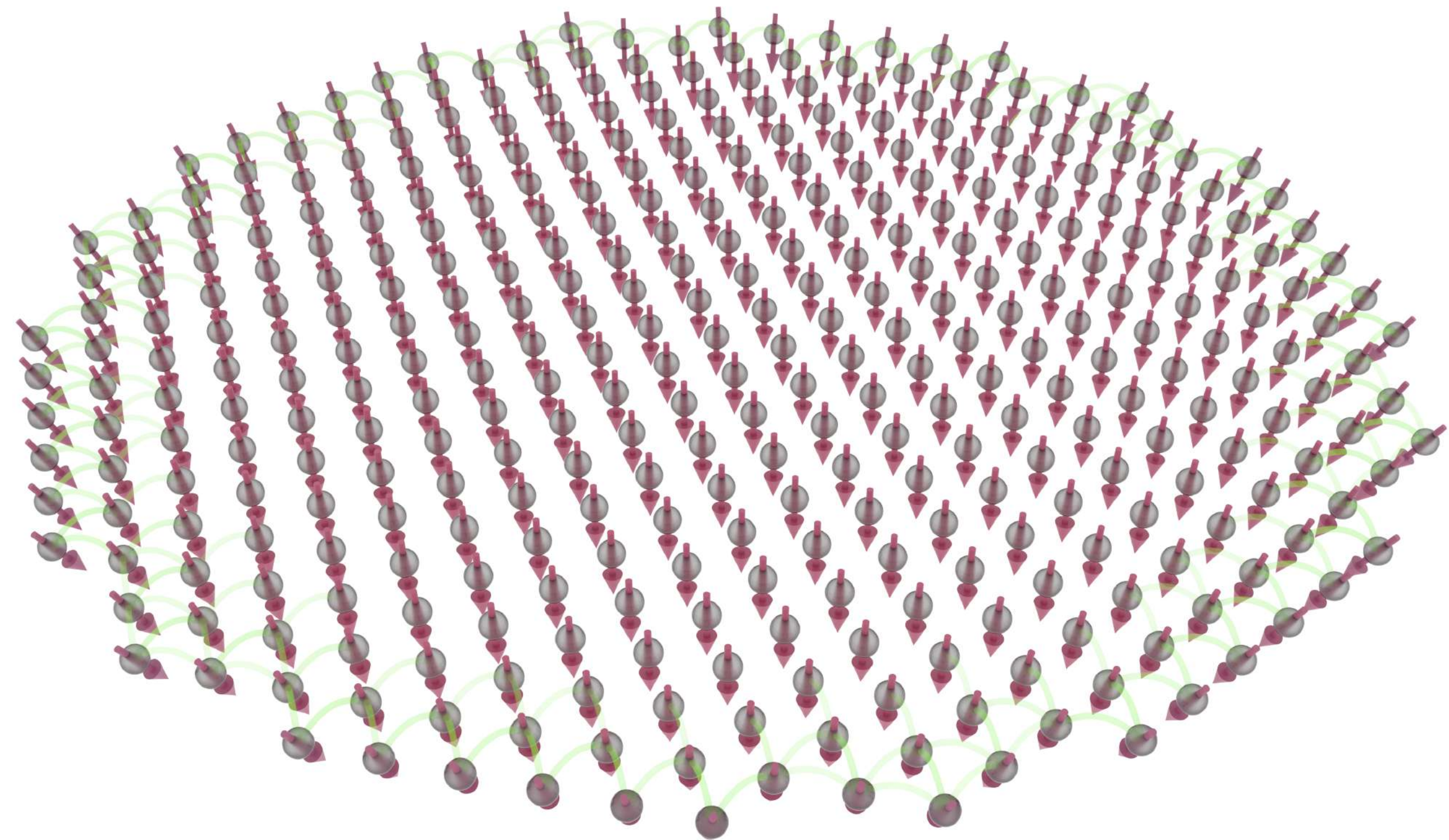
- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- **Quantum skyrmions are entangled** superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- **Quantum skyrmion lattices** in the thermodynamic limit

$$B_0 = 1$$

$$\mathbf{B} = -B_0 \mathbf{e}_z$$

$$J = -2.0D$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + D_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{B} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$D_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

- confirm entanglement shared between two spins (Osterloh et al., Nature 416)
 - field-polarized ground state



Quantum Skyrmions

arXiv:2112.12475*

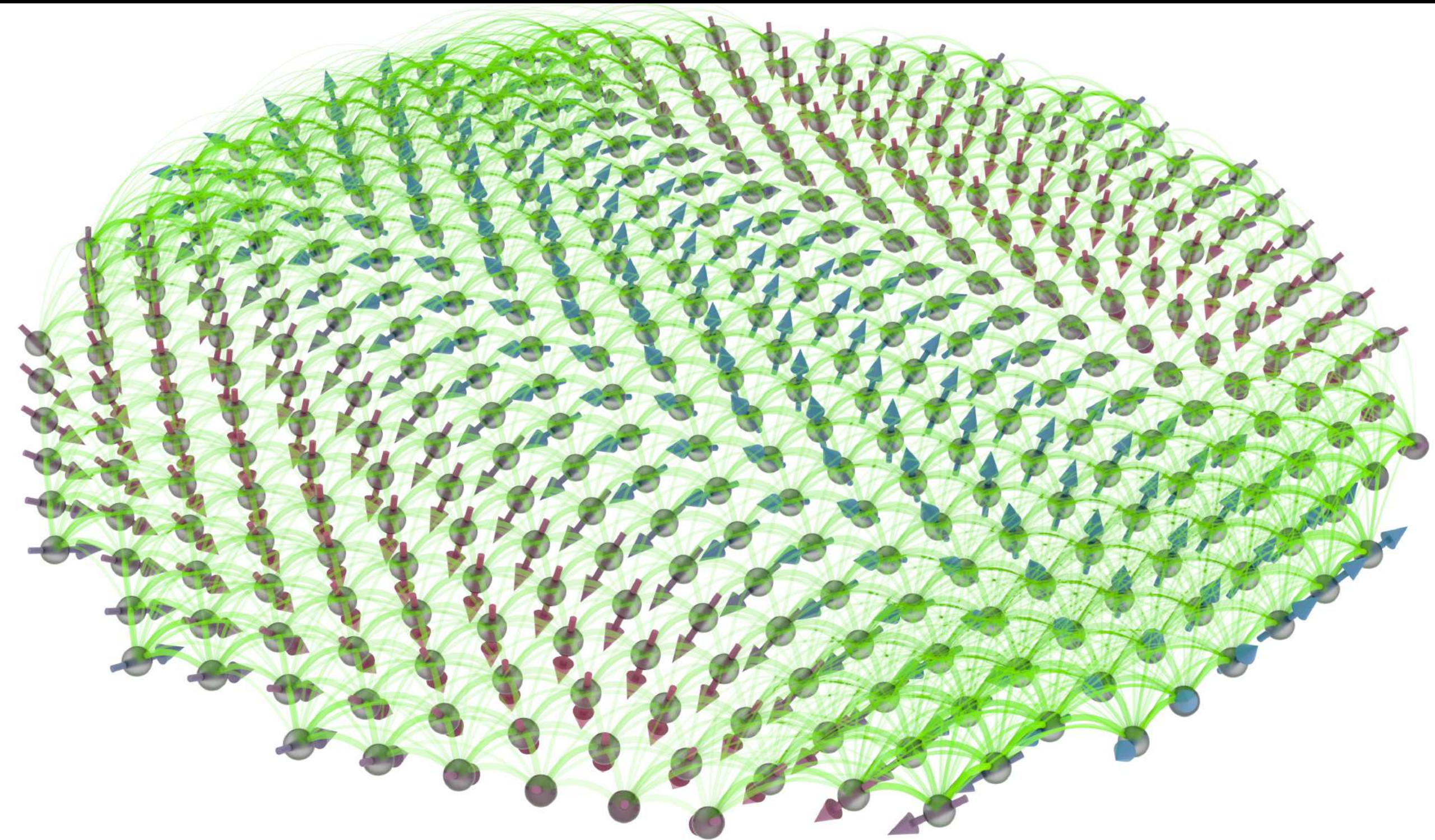
- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- **Quantum skyrmions are entangled** superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- **Quantum skyrmion lattices** in the thermodynamic limit

$$B_0 = 1/8$$

$$B = -B_0 \mathbf{e}_z$$

$$J = -2.0D$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + D_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + B \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$D_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

- confirm entanglement shared between two spins (Osterloh et al., Nature 416)
- density wave ground state



Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- **Quantum skyrmions are entangled** superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- **Quantum skyrmion lattices** in the thermodynamic limit

$$B_0 = 1/2$$

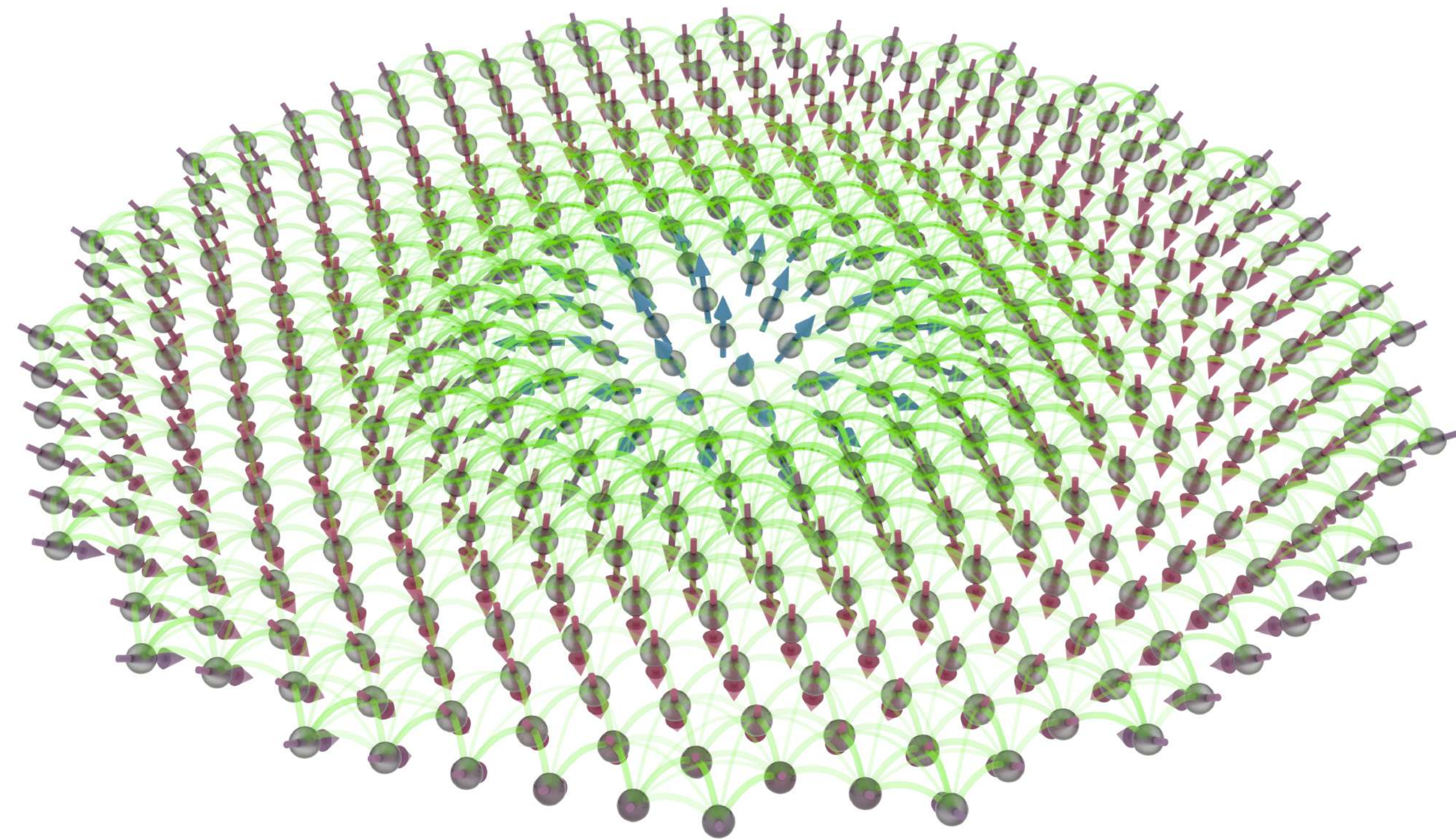
$$\mathbf{B} = -B_0 \mathbf{e}_z$$

$$J = -2.0D$$

$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{B} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$

$$\mathbf{D}_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

- confirm entanglement shared between two spins (Osterloh et al., Nature 416)
 - quantum skyrmion ground state



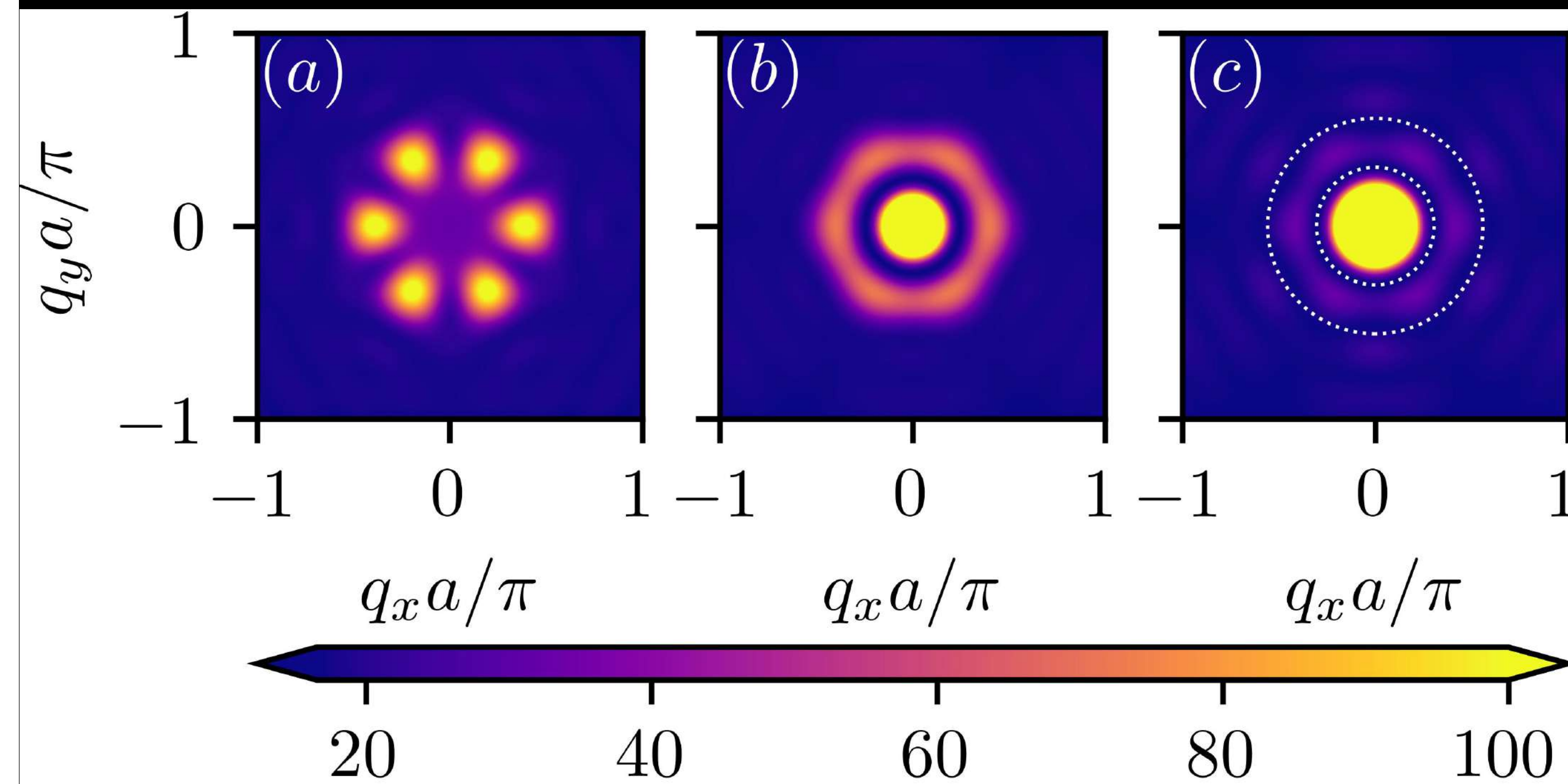
Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- **Quantum skyrmions are *entangled*** superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- **Quantum skyrmion lattices** in the thermodynamic limit

$$J = -0.5D \quad B = -B_0 \mathbf{e}_z$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + D_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + B \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$D_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

- Elastic magnetic differential neutron scattering cross section (extracted from spin structure factors)



Quantum Skyrmions

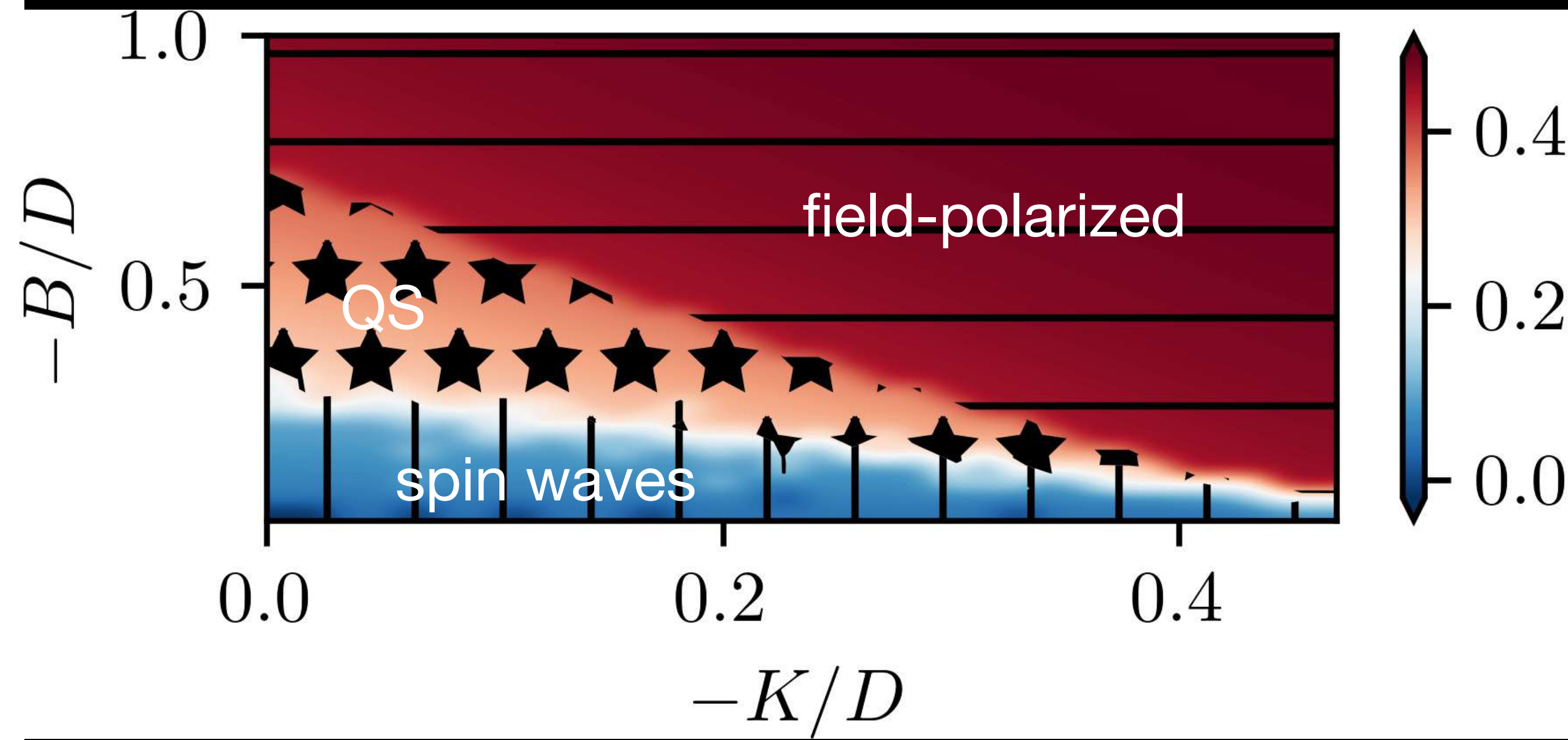
arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- *Quantum skyrmions are entangled* superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- *Quantum skyrmion lattices* in the thermodynamic limit

$$J = -0.5D \quad B = -B_0 e_z$$
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$
$$D_{a_i} = D e_z \times a_i$$

- Add anisotropy to study meltdown of the skyrmion phase

$$\hat{H}_{\text{pert.}} = -K \sum_{\langle r, r' \rangle} \hat{S}_r^z \hat{S}_{r'}^z$$



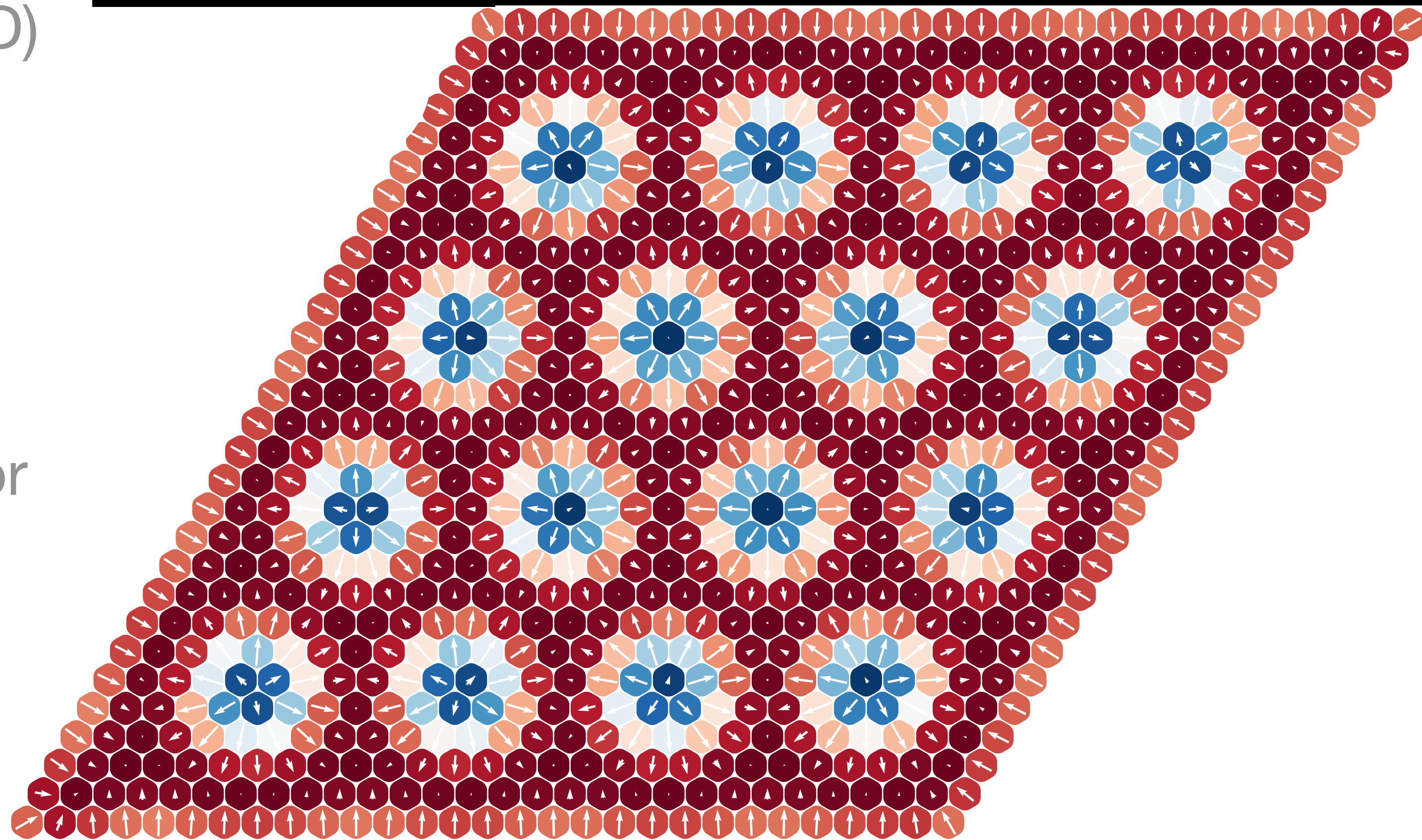
Quantum Skyrmions

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- **Quantum skyrmions are entangled** superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- **Quantum skyrmion lattices** in the thermodynamic limit

$$J = -0.5D \quad B = -B_0 \mathbf{e}_z$$
$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + D_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + B \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$
$$D_{\mathbf{a}_i} = D \mathbf{e}_z \times \mathbf{a}_i$$

- locality of entanglement allows to access large system sizes — skyrmion lattices!



Quantum Spintronics

Future directions

- analytic description of the quantum skyrmion lattice phase
- find and utilize unique dynamics of quantum skyrmion lattices
 - tensor networks
 - state tomography
 - time evolution with tensor networks
- finite temperature simulations (using auxiliary field QMC?)

$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{B} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$

Quantum Spintronics

Future directions

- analytic description of the quantum skyrmion lattice phase
- find and utilize unique dynamics of quantum skyrmion lattices
 - tensor networks
 - state tomography
 - time evolution with tensor networks
 - finite temperature simulations (using auxiliary field QMC?)

$$\hat{H} = J \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} \left(\hat{\mathbf{S}}_{\mathbf{r}} \cdot \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{D}_{\mathbf{r}-\mathbf{r}'} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \times \hat{\mathbf{S}}_{\mathbf{r}'} + \mathbf{B} \cdot \hat{\mathbf{S}}_{\mathbf{r}} \delta_{\mathbf{r}, \mathbf{r}'} \right)$$

- To use Hubbard-Stratonovich trafo's, we must write interactions as perfect squares

$$\epsilon_{\alpha\beta\gamma} \hat{S}_{\mathbf{r}}^{\beta} \hat{S}_{\mathbf{r}'}^{\gamma} = \frac{|\epsilon_{\alpha\beta\gamma}|}{2} \left(\hat{S}_{\mathbf{r}}^{\beta} + \text{sign}(\epsilon_{\alpha\beta\gamma}) \hat{S}_{\mathbf{r}'}^{\gamma} \right)^2 + \dots$$

Quantum Spintronics

Future directions

- analytic description of the quantum skyrmion lattice phase
- find and utilize unique dynamics of quantum skyrmion lattices
 - tensor networks
 - state tomography
 - time evolution with tensor networks
 - finite temperature simulations (using auxiliary field QMC?)

$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$

- To use Hubbard-Stratonovich trafo's, we must write interactions as perfect squares
$$\epsilon_{\alpha\beta\gamma} \hat{S}_r^\beta \hat{S}_{r'}^\gamma = \frac{|\epsilon_{\alpha\beta\gamma}|}{2} \left(\hat{S}_r^\beta + \text{sign}(\epsilon_{\alpha\beta\gamma}) \hat{S}_{r'}^\gamma \right)^2 + \dots$$
- “...” refers to terms that can be made constant by restrictions on the local Hilbert space (on-site Hubbard repulsion)

Quantum Spintronics

Future directions

- analytic description of the quantum skyrmion lattice phase
- find and utilize unique dynamics of quantum skyrmion lattices
 - tensor networks
 - state tomography
 - time evolution with tensor networks
 - finite temperature simulations (using auxiliary field QMC?)

$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{\mathbf{S}}_r \cdot \hat{\mathbf{S}}_{r'} + \mathbf{D}_{r-r'} \cdot \hat{\mathbf{S}}_r \times \hat{\mathbf{S}}_{r'} + \mathbf{B} \cdot \hat{\mathbf{S}}_r \delta_{r, r'} \right)$$

- To use Hubbard-Stratonovich trafo's, we must write interactions as perfect squares
$$\epsilon_{\alpha\beta\gamma} \hat{S}_r^\beta \hat{S}_{r'}^\gamma = \frac{|\epsilon_{\alpha\beta\gamma}|}{2} \left(\hat{S}_r^\beta + \text{sign}(\epsilon_{\alpha\beta\gamma}) \hat{S}_{r'}^\gamma \right)^2 + \dots$$
- “...” refers to terms that can be made constant by restrictions on the local Hilbert space (on-site Hubbard repulsion)
- Auxiliary field QMC: fermionize...
$$\hat{S}_i^\alpha = \frac{1}{2} \hat{f}_i^\dagger \sigma_\alpha \hat{f}_i$$

Quantum Spintronics

Future directions

- analytic description of the quantum skyrmion lattice phase
- find and utilize unique dynamics of quantum skyrmion lattices
 - tensor networks
 - state tomography
 - time evolution with tensor networks
 - finite temperature simulations (using auxiliary field QMC?)

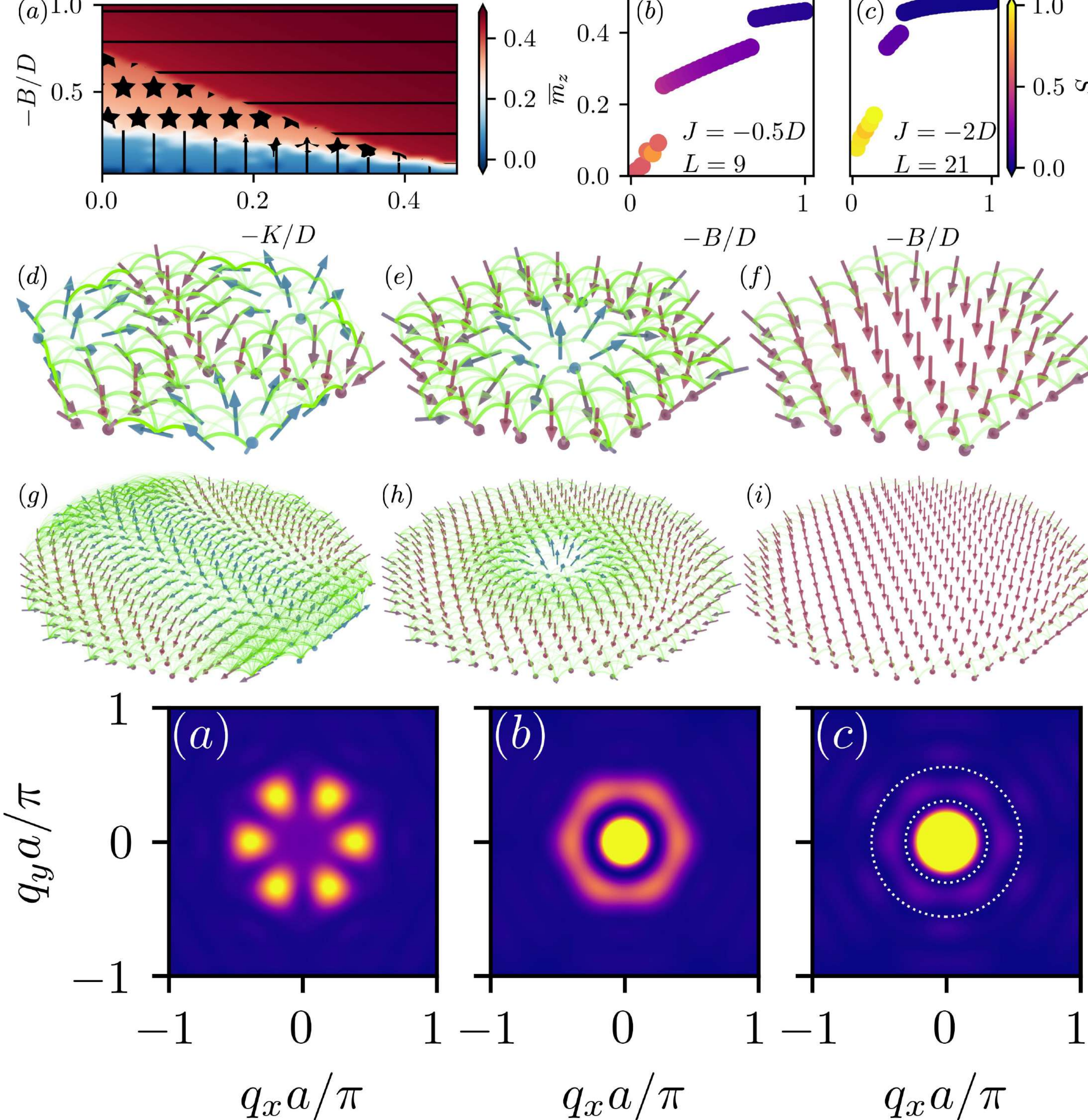
$$\hat{H} = J \sum_{\langle r, r' \rangle} \left(\hat{S}_r \cdot \hat{S}_{r'} + D_{r-r'} \cdot \hat{S}_r \times \hat{S}_{r'} + B \cdot \hat{S}_r \delta_{r, r'} \right)$$

- To use Hubbard-Stratonovich trafo's, we must write interactions as perfect squares
$$\epsilon_{\alpha\beta\gamma} \hat{S}_r^\beta \hat{S}_{r'}^\gamma = \frac{|\epsilon_{\alpha\beta\gamma}|}{2} \left(\hat{S}_r^\beta + \text{sign}(\epsilon_{\alpha\beta\gamma}) \hat{S}_{r'}^\gamma \right)^2 + \dots$$
- “...” refers to terms that can be made constant by restrictions on the local Hilbert space (on-site Hubbard repulsion)
- Auxiliary field QMC: fermionize...
$$\hat{S}_i^\alpha = \frac{1}{2} \hat{f}_i^\dagger \sigma_\alpha \hat{f}_i$$
- Sign problem? If yes, how bad is it...

Summary

arXiv:2112.12475*

- Use tensor networks (matrix product states) to overcome size limits of exact diagonalization (ED)
- **Quantum skyrmions are entangled** superpositions of classical states
- Detect quantum skyrmions by measuring average polarization or neutron scattering cross section
- **Quantum skyrmion lattices** in the thermodynamic limit



Acknowledgements

arXiv:2112.12475*

Quantum skyrmion lattices in Heisenberg Ferromagnets

- Dr. Alireza Habibi (Uni Lu)
- Dr. Solofo Groenendijk (Uni Lu)
- Prof. Andreas Michels
- Prof. Thomas Schmidt

